

# Supersymmetric Particle Searches

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## SUPERSYMMETRIC MODEL ASSUMPTIONS

The exclusion of particle masses within a mass range ( $m_1, m_2$ ) will be denoted with the notation “none  $m_1-m_2$ ” in the VALUE column of the following Listings. The latest unpublished results are described in the “Supersymmetry: Experiment” review.

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## $\tilde{\chi}_1^0$ (Lightest Neutralino) MASS LIMIT

$\tilde{\chi}_1^0$  is often assumed to be the lightest supersymmetric particle (LSP). See also the  $\tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0$  section below.

We have divided the  $\tilde{\chi}_1^0$  listings below into five sections:

- 1) Accelerator limits for stable  $\tilde{\chi}_1^0$ ,
- 2) Bounds on  $\tilde{\chi}_1^0$  from dark matter searches,
- 3)  $\tilde{\chi}_1^0 - p$  elastic cross section (spin-dependent, spin-independent interactions),
- 4) Other bounds on  $\tilde{\chi}_1^0$  from astrophysics and cosmology, and
- 5) Unstable  $\tilde{\chi}_1^0$  (Lightest Neutralino) mass limit.

### ———— Accelerator limits for stable $\tilde{\chi}_1^0$ ————

Unless otherwise stated, results in this section assume spectra, production rates, decay modes, and branching ratios as evaluated in the MSSM, with gaugino and sfermion mass unification at the GUT scale. These papers generally study production of  $\tilde{\chi}_i^0 \tilde{\chi}_j^0$  ( $i \geq 1, j \geq 2$ ),  $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ , and (in the case of hadronic collisions)  $\tilde{\chi}_1^+ \tilde{\chi}_2^0$  pairs. The mass limits on  $\tilde{\chi}_1^0$  are either direct, or follow indirectly from the constraints set by the non-observation of  $\tilde{\chi}_1^\pm$  and  $\tilde{\chi}_2^0$  states on the gaugino and higgsino MSSM parameters  $M_2$  and  $\mu$ . In some cases, information is used from the nonobservation of slepton decays.

Obsolete limits obtained from  $e^+ e^-$  collisions up to  $\sqrt{s}=184$  GeV have been removed from this compilation and can be found in the 2000 Edition (The European Physical Journal **C15** 1 (2000)) of this Review.

$$\Delta m = m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}.$$

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
>40	95	<sup>1</sup> ABBIENDI	04H OPAL	all $\tan\beta$ , $\Delta m > 5$ GeV, $m_0 > 500$ GeV, $A_0 = 0$
>42.4	95	<sup>2</sup> HEISTER	04 ALEP	all $\tan\beta$ , all $\Delta m$ , all $m_0$
>39.2	95	<sup>3</sup> ABDALLAH	03M DLPH	all $\tan\beta$ , $m_{\tilde{\nu}} > 500$ GeV
<b>&gt;46</b>	95	<sup>4</sup> ABDALLAH	03M DLPH	all $\tan\beta$ , all $\Delta m$ , all $m_0$
>32.5	95	<sup>5</sup> ACCIARRI	00D L3	$\tan\beta > 0.7$ , $\Delta m > 3$ GeV, all $m_0$

• • • We do not use the following data for averages, fits, limits, etc. • • •

<sup>6</sup> DREINER 09 THEO

<sup>1</sup> ABBIENDI 04H search for charginos and neutralinos in events with acoplanar leptons+jets and multi-jet final states in the 192–209 GeV data, combined with the results on leptonic final states from ABBIENDI 04. The results hold for a scan over the parameter space covering the region  $0 < M_2 < 5000$  GeV,  $-1000 < \mu < 1000$  GeV and  $\tan\beta$  from 1 to 40. This limit supersedes ABBIENDI 00H.

<sup>2</sup> HEISTER 04 data collected up to 209 GeV. Updates earlier analysis of selectrons from HEISTER 02E, includes a new analysis of charginos and neutralinos decaying into stau and uses results on charginos with initial state radiation from HEISTER 02J. The limit is based on the direct search for charginos and neutralinos, the constraints from the slepton search and the Higgs mass limits from HEISTER 02 using a top mass of 175 GeV, interpreted in a framework with universal gaugino and sfermion masses. Assuming the mixing in the stau sector to be negligible, the limit improves to 43.1 GeV. Under the

assumption of MSUGRA with unification of the Higgs and sfermion masses, the limit improves to 50 GeV, and reaches 53 GeV for  $A_0 = 0$ . These limits include and update the results of BARATE 01.

<sup>3</sup> ABDALLAH 03M uses data from  $\sqrt{s} = 192\text{--}208$  GeV. A limit on the mass of  $\tilde{\chi}_1^0$  is derived from direct searches for neutralinos combined with the chargino search. Neutralinos are searched in the production of  $\tilde{\chi}_1^0\tilde{\chi}_2^0$ ,  $\tilde{\chi}_1^0\tilde{\chi}_3^0$ , as well as  $\tilde{\chi}_2^0\tilde{\chi}_3^0$  and  $\tilde{\chi}_2^0\tilde{\chi}_4^0$  giving rise to cascade decays, and  $\tilde{\chi}_1^0\tilde{\chi}_2^0$  and  $\tilde{\chi}_1^0\tilde{\chi}_2^0$ , followed by the decay  $\tilde{\chi}_2^0 \rightarrow \tilde{\tau}\tau$ . The results hold for the parameter space defined by values of  $M_2 < 1$  TeV,  $|\mu| \leq 2$  TeV with the  $\tilde{\chi}_1^0$  as LSP. The limit is obtained for  $\tan\beta = 1$  and large  $m_0$ , where  $\tilde{\chi}_2^0\tilde{\chi}_4^0$  and chargino pair production are important. If the constraint from Higgs searches is also imposed, the limit improves to 49.0 GeV in the  $m_h^{\text{max}}$  scenario with  $m_t=174.3$  GeV. These limits update the results of ABREU 00J.

<sup>4</sup> ABDALLAH 03M uses data from  $\sqrt{s} = 192\text{--}208$  GeV. An indirect limit on the mass of  $\tilde{\chi}_1^0$  is derived by constraining the MSSM parameter space by the results from direct searches for neutralinos (including cascade decays and  $\tilde{\tau}\tau$  final states), for charginos (for all  $\Delta m_+$ ) and for sleptons, stop and sbottom. The results hold for the full parameter space defined by values of  $M_2 < 1$  TeV,  $|\mu| \leq 2$  TeV with the  $\tilde{\chi}_1^0$  as LSP. Constraints from the Higgs search in the  $m_h^{\text{max}}$  scenario assuming  $m_t=174.3$  GeV are included. The limit is obtained for  $\tan\beta \geq 5$  when stau mixing leads to mass degeneracy between  $\tilde{\tau}_1$  and  $\tilde{\chi}_1^0$  and the limit is based on  $\tilde{\chi}_2^0$  production followed by its decay to  $\tilde{\tau}_1\tau$ . In the pathological scenario where  $m_0$  and  $|\mu|$  are large, so that the  $\tilde{\chi}_2^0$  production cross section is negligible, and where there is mixing in the stau sector but not in stop nor sbottom, the limit is based on charginos with soft decay products and an ISR photon. The limit then degrades to 39 GeV. See Figs 40–42 for the dependence of the limit on  $\tan\beta$  and  $m_{\tilde{\nu}}$ . These limits update the results of ABREU 00W.

<sup>5</sup> ACCIARRI 00D data collected at  $\sqrt{s}=189$  GeV. The results hold over the full parameter space defined by  $0.7 \leq \tan\beta \leq 60$ ,  $0 \leq M_2 \leq 2$  TeV,  $m_0 \leq 500$  GeV,  $|\mu| \leq 2$  TeV. The minimum mass limit is reached for  $\tan\beta=1$  and large  $m_0$ . The results of slepton searches from ACCIARRI 99W are used to help set constraints in the region of small  $m_0$ . The limit improves to 48 GeV for  $m_0 \gtrsim 200$  GeV and  $\tan\beta \gtrsim 10$ . See their Figs. 6–8 for the  $\tan\beta$  and  $m_0$  dependence of the limits. Updates ACCIARRI 98F.

<sup>6</sup> DREINER 09 show that in the general MSSM with non-universal gaugino masses there exists no model-independent laboratory bound on the mass of the lightest neutralino. An essentially massless  $\tilde{\chi}_1^0$  is allowed by the experimental and observational data, imposing some constraints on other MSSM parameters, including  $M_2$ ,  $\mu$  and the slepton and squark masses.

### ————— Bounds on $\tilde{\chi}_1^0$ from dark matter searches —————

These papers generally exclude regions in the  $M_2 - \mu$  parameter plane assuming that  $\tilde{\chi}_1^0$  is the dominant form of dark matter in the galactic halo. These limits are based on the lack of detection in laboratory experiments, telescopes, or by the absence of a signal in underground neutrino detectors. The latter signal is expected if  $\tilde{\chi}_1^0$  accumulates in the Sun or the Earth and annihilates into high-energy  $\nu$ 's.

VALUE	DOCUMENT ID	TECN
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• • • We do not use the following data for averages, fits, limits, etc. • • •

1	AARTSEN	13	ICCB
2	AARTSEN	13C	ICCB
3	ABRAMOWSKI13		HESS

	4	ACKERMANN 13A	FRMI
	5	ADRIAN-MAR.13	ANTR
	6	BERGSTROM 13	COSM
	7	BOLIEV 13	BAKS
	6	JIN 13	ASTR
	6	KOPP 13	COSM
	8	ABBASI 12	ICCB
	9	ABRAMOWSKI11	HESS
	10	ABDO 10	FRMI
	11	ACKERMANN 10	FRMI
	12	ABBASI 09B	ICCB
	13	ACHTERBERG 06	AMND
	14	ACKERMANN 06	AMND
	15	DEBOER 06	RVUE
	16	DESAI 04	SKAM
	16	AMBROSIO 99	MCRO
	17	LOSECCO 95	RVUE
	18	MORI 93	KAMI
	19	BOTTINO 92	COSM
	20	BOTTINO 91	RVUE
	21	GELMINI 91	COSM
	22	KAMIONKOW.91	RVUE
	23	MORI 91B	KAMI
none 4–15 GeV	24	OLIVE 88	COSM

<sup>1</sup> AARTSEN 13 is based on data collected during 317 effective days with the IceCube 79-string detector including the DeepCore sub-array. They looked for interactions of  $\nu_{\mu}$ 's from neutralino annihilations in the Sun over a background of atmospheric neutrinos and set 90% CL limits on the muon flux. They also obtain limits on the spin dependent and spin independent neutralino-proton cross section for neutralino masses in the range 20–5000 GeV.

<sup>2</sup> AARTSEN 13C is based on data collected during 339.8 effective days with the IceCube 59-string detector. They looked for interactions of  $\nu_{\mu}$ 's from neutralino annihilations in nearby galaxies and galaxy clusters. They obtain limits on the neutralino annihilation cross section for neutralino masses in the range 30–100,000 GeV.

<sup>3</sup> ABRAMOWSKI 13 place upper limits on the annihilation cross section with  $\gamma\gamma$  final states in the energy range of 0.5–25 TeV.

<sup>4</sup> ACKERMANN 13A is based on 3.7 years of data with Fermi-LAT and search for monochromatic gamma-rays in the energy range of 5–300 GeV from dark matter annihilations. No globally significant lines are reported.

<sup>5</sup> ADRIAN-MARTINEZ 13 is based on data from the ANTARES neutrino telescope. They looked for interactions of  $\nu_{\mu}$ 's from neutralino annihilations in the Sun over a background of atmospheric neutrinos and set 90% CL limits on the muon flux. They also obtain limits on the spin dependent and spin independent neutralino-proton cross section for neutralino masses in the range 50–10,000 GeV.

<sup>6</sup> BERGSTROM 13, JIN 13, and KOPP 13 derive limits on the mass and annihilation cross section using AMS-02 data. JIN 13 also sets a limit on the lifetime of the dark matter particle.

<sup>7</sup> BOLIEV 13 is based on data collected during 24.12 years of live time with the Bakson Underground Scintillator Telescope. They looked for interactions of  $\nu_{\mu}$ 's from neutralino annihilations in the Sun over a background of atmospheric neutrinos and set 90% CL limits on the muon flux. They also obtain limits on the spin dependent and spin independent neutralino-proton cross section for neutralino masses in the range 10–1000 GeV.

- <sup>8</sup> ABBASI 12 is based on data collected during 812 effective days with AMANDA II and 149 days of the IceCube 40-string detector combined with the data of ABBASI 09B. They looked for interactions of  $\nu_\mu$ 's from neutralino annihilations in the Sun over a background of atmospheric neutrinos and set 90% CL limits on the muon flux. No excess is observed. They also obtain limits on the spin dependent neutralino-proton cross section for neutralino masses in the range 50–5000 GeV.
- <sup>9</sup> ABRAMOWSKI 11 place upper limits on the annihilation cross section with  $\gamma\gamma$  final states.
- <sup>10</sup> ABDO 10 place upper limits on the annihilation cross section with  $\gamma\gamma$  or  $\mu^+\mu^-$  final states.
- <sup>11</sup> ACKERMANN 10 place upper limits on the annihilation cross section with  $b\bar{b}$  or  $\mu^+\mu^-$  final states.
- <sup>12</sup> ABBASI 09B is based on data collected during 104.3 effective days with the IceCube 22-string detector. They looked for interactions of  $\nu_\mu$ 's from neutralino annihilations in the Sun over a background of atmospheric neutrinos and set 90% CL limits on the muon flux. They also obtain limits on the spin dependent neutralino-proton cross section for neutralino masses in the range 250–5000 GeV.
- <sup>13</sup> ACHTERBERG 06 is based on data collected during 421.9 effective days with the AMANDA detector. They looked for interactions of  $\nu_\mu$ 's from the centre of the Earth over a background of atmospheric neutrinos and set 90% CL limits on the muon flux. Their limit is compared with the muon flux expected from neutralino annihilations into  $W^+W^-$  and  $b\bar{b}$  at the centre of the Earth for MSSM parameters compatible with the relic dark matter density, see their Fig. 7.
- <sup>14</sup> ACKERMANN 06 is based on data collected during 143.7 days with the AMANDA-II detector. They looked for interactions of  $\nu_\mu$ 's from the Sun over a background of atmospheric neutrinos and set 90% CL limits on the muon flux. Their limit is compared with the muon flux expected from neutralino annihilations into  $W^+W^-$  in the Sun for SUSY model parameters compatible with the relic dark matter density, see their Fig. 3.
- <sup>15</sup> DEBOER 06 interpret an excess of diffuse Galactic gamma rays observed with the EGRET satellite as originating from  $\pi^0$  decays from the annihilation of neutralinos into quark jets. They analyze the corresponding parameter space in a supergravity inspired MSSM model with radiative electroweak symmetry breaking, see their Fig. 3 for the preferred region in the  $(m_0, m_{1/2})$  plane of a scenario with large  $\tan\beta$ .
- <sup>16</sup> AMBROSIO 99 and DESAI 04 set new neutrino flux limits which can be used to limit the parameter space in supersymmetric models based on neutralino annihilation in the Sun and the Earth.
- <sup>17</sup> LOSECCO 95 reanalyzed the IMB data and places lower limit on  $m_{\tilde{\chi}_1^0}$  of 18 GeV if the LSP is a photino and 10 GeV if the LSP is a higgsino based on LSP annihilation in the sun producing high-energy neutrinos and the limits on neutrino fluxes from the IMB detector.
- <sup>18</sup> MORI 93 excludes some region in  $M_2-\mu$  parameter space depending on  $\tan\beta$  and lightest scalar Higgs mass for neutralino dark matter  $m_{\tilde{\chi}_1^0} > m_W$ , using limits on upgoing muons produced by energetic neutrinos from neutralino annihilation in the Sun and the Earth.
- <sup>19</sup> BOTTINO 92 excludes some region  $M_2-\mu$  parameter space assuming that the lightest neutralino is the dark matter, using upgoing muons at Kamiokande, direct searches by Ge detectors, and by LEP experiments. The analysis includes top radiative corrections on Higgs parameters and employs two different hypotheses for nucleon-Higgs coupling. Effects of rescaling in the local neutralino density according to the neutralino relic abundance are taken into account.
- <sup>20</sup> BOTTINO 91 excluded a region in  $M_2-\mu$  plane using upgoing muon data from Kamioka experiment, assuming that the dark matter surrounding us is composed of neutralinos and that the Higgs boson is not too heavy.
- <sup>21</sup> GELMINI 91 exclude a region in  $M_2-\mu$  plane using dark matter searches.
- <sup>22</sup> KAMIONKOWSKI 91 excludes a region in the  $M_2-\mu$  plane using IMB limit on upgoing muons originated by energetic neutrinos from neutralino annihilation in the sun, assuming

that the dark matter is composed of neutralinos and that  $m_{H_1^0} \lesssim 50$  GeV. See Fig. 8 in the paper.

<sup>23</sup> MORI 91B exclude a part of the region in the  $M_2-\mu$  plane with  $m_{\tilde{\chi}_1^0} \lesssim 80$  GeV using a limit on upgoing muons originated by energetic neutrinos from neutralino annihilation in the earth, assuming that the dark matter surrounding us is composed of neutralinos and that  $m_{H_1^0} \lesssim 80$  GeV.

<sup>24</sup> OLIVE 88 result assumes that photinos make up the dark matter in the galactic halo. Limit is based on annihilations in the sun and is due to an absence of high energy neutrinos detected in underground experiments. The limit is model dependent.

—————  $\tilde{\chi}_1^0$ - $p$  elastic cross section —————

Experimental results on the  $\tilde{\chi}_1^0$ - $p$  elastic cross section are evaluated at  $m_{\tilde{\chi}_1^0}=100$  GeV. The experimental results on the cross section are often mass dependent. Therefore, the mass and cross section results are also given where the limit is strongest, when appropriate. Results are quoted separately for spin-dependent interactions (based on an effective 4-Fermi Lagrangian of the form  $\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$ ) and spin-independent interactions ( $\bar{\chi}\chi\bar{q}q$ ). For calculational details see GRIEST 88B, ELLIS 88D, BARBIERI 89C, DREES 93B, ARNOWITT 96, BERGSTROM 96, and BAER 97 in addition to the theory papers listed in the Tables. For a description of the theoretical assumptions and experimental techniques underlying most of the listed papers, see the review on "Dark matter" in this "Review of Particle Physics," and references therein. Most of the following papers use galactic halo and nuclear interaction assumptions from (LEWIN 96).

**Spin-dependent interactions**

VALUE (pb)	CL%	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
< 0.01	90	1 APRILE 13	X100	Xe
< 0.01	90	2 AKIMOV 12	ZEP3	Xe
< 0.07	90	3 ARCHAMBAU..12	PICA	F
< 7 × 10 <sup>-3</sup>		4 BEHNKE 12	COUP	CF <sub>3</sub> I
< 1.8	90	5 DAW 12	DRFT	CS <sub>2</sub> ; CF <sub>4</sub>
< 8.5 × 10 <sup>-3</sup>		6 FELIZARDO 12	SMPL	C <sub>2</sub> ClF <sub>5</sub>
< 0.016	90	7 KIM 12	KIMS	CsI
5 × 10 <sup>-10</sup> to 10 <sup>-5</sup>	95	8 BUCHMUEL... 11B	THEO	
< 1	90	9 ANGLE 08A	XE10	Xe
< 0.055		10 BEDNYAKOV 08	HDMS	Ge
< 0.33	90	11 BEHNKE 08	COUP	CF <sub>3</sub> I
< 5		12 AKERIB 06	CDMS	Ge
< 2		13 SHIMIZU 06A	CNTR	CaF <sub>2</sub>
< 0.4		14 ALNER 05	NAIA	NaI Spin Dep.
< 2		15 BARNABE-HE..05	PICA	C
2 × 10 <sup>-11</sup> to 1 × 10 <sup>-4</sup>		16 ELLIS 04	THEO	μ > 0
< 0.8		17 AHMED 03	NAIA	NaI Spin Dep.
< 40		18 TAKEDA 03	BOLO	NaF Spin Dep.

< 10	19	ANGLOHER	02	CRES	Sapphire
$8 \times 10^{-7}$ to $2 \times 10^{-5}$	20	ELLIS	01C	THEO	$\tan\beta \leq 10$
< 3.8	21	BERNABEI	00D	DAMA	Xe
< 0.8		SPOONER	00	UKDM	Nal
< 4.8	22	BELLI	99C	DAMA	F
<100	23	OOTANI	99	BOLO	LiF
< 0.6		BERNABEI	98C	DAMA	Xe
< 5	22	BERNABEI	97	DAMA	F

<sup>1</sup> The strongest limit is 0.006 pb and occurs at  $m_\chi = 60$  GeV. APRILE 13 also presents limits for the scattering on neutrons. At 100 GeV, the upper limit is  $4 \times 10^{-4}$  pb and the strongest limit is  $3.5 \times 10^{-4}$  pb at 45 GeV.

<sup>2</sup> This result updates LEBEDENKO 09A. The strongest limit is  $8 \times 10^{-3}$  pb at  $m_\chi = 50$  GeV. Limit applies to the neutralino neutron elastic cross section.

<sup>3</sup> This result updates ARCHAMBAULT 09. The strongest limit is 0.032 pb at  $m_\chi = 20$  GeV.

<sup>4</sup> The strongest limit is  $6 \times 10^{-3}$  at  $m_\chi = 60$  GeV.

<sup>5</sup> The strongest limit is 1.8 pb and occurs at  $m_\chi = 100$  GeV.

<sup>6</sup> The strongest limit is  $5.7 \times 10^{-3}$  at  $m_\chi = 35$  GeV.

<sup>7</sup> This result updates LEE 07A. The strongest limit is at  $m_\chi = 80$  GeV.

<sup>8</sup> Predictions for the spin-dependent elastic cross section based on a frequentist approach to electroweak observables in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry.

<sup>9</sup> The strongest limit is 0.6 pb and occurs at  $m_\chi = 30$  GeV. The limit for scattering on neutrons is 0.01 pb at  $m_\chi = 100$  GeV, and the strongest limit is 0.0045 pb at  $m_\chi = 30$  GeV.

<sup>10</sup> Limit applies to neutron elastic cross section.

<sup>11</sup> The strongest upper limit is 0.25 pb and occurs at  $m_\chi \simeq 40$  GeV.

<sup>12</sup> The strongest upper limit is 4 pb and occurs at  $m_\chi \simeq 60$  GeV. The limit on the neutron spin-dependent elastic cross section is 0.07 pb. This latter limit is improved in AHMED 09, where a limit of 0.02 pb is obtained at  $m_\chi = 100$  GeV. The strongest limit in AHMED 09 is 0.018 pb and occurs at  $m_\chi = 60$  GeV.

<sup>13</sup> The strongest upper limit is 1.2 pb and occurs at  $m_\chi \simeq 40$  GeV. The limit on the neutron spin-dependent cross section is 35 pb.

<sup>14</sup> The strongest upper limit is 0.35 pb and occurs at  $m_\chi \simeq 60$  GeV.

<sup>15</sup> The strongest upper limit is 1.2 pb and occurs  $m_\chi \simeq 30$  GeV.

<sup>16</sup> ELLIS 04 calculates the  $\chi p$  elastic scattering cross section in the framework of  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry, but without universal scalar masses. In the case of universal squark and slepton masses, but non-universal Higgs masses, the limit becomes  $2 \times 10^{-4}$ , see ELLIS 03E.

<sup>17</sup> The strongest upper limit is 0.75 pb and occurs at  $m_\chi \approx 70$  GeV.

<sup>18</sup> The strongest upper limit is 30 pb and occurs at  $m_\chi \approx 20$  GeV.

<sup>19</sup> The strongest upper limit is 8 pb and occurs at  $m_\chi \simeq 30$  GeV.

<sup>20</sup> ELLIS 01C calculates the  $\chi$ - $p$  elastic scattering cross section in the framework of  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry. In models with nonuniversal Higgs masses, the upper limit to the cross section is  $6 \times 10^{-4}$ .

<sup>21</sup> The strongest upper limit is 3 pb and occurs at  $m_\chi \simeq 60$  GeV. The limits are for inelastic scattering  $X^0 + {}^{129}\text{Xe} \rightarrow X^0 + {}^{129}\text{Xe}^*$  (39.58 keV).

<sup>22</sup> The strongest upper limit is 4.4 pb and occurs at  $m_\chi \simeq 60$  GeV.

<sup>23</sup> The strongest upper limit is about 35 pb and occurs at  $m_\chi \simeq 15$  GeV.

### Spin-independent interactions

<u>VALUE (pb)</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
$< 2.2 \times 10^{-6}$	90	1 AALSETH	13 CGNT	Ge
		2 AGNESE	13 CDMS	Si
		3 LI	13B TEXO	Ge
$< 5 \times 10^{-8}$	90	4 AKIMOV	12 ZEP3	Xe
$1.6 \times 10^{-6}; 3.7 \times 10^{-5}$		5 ANGLOHER	12 CRES	CaWO <sub>4</sub>
$< 2.6 \times 10^{-9}$	90	6 APRILE	12 X100	Xe
	90	7 ARCHAMBAU..12	PICA	C <sub>4</sub> F <sub>10</sub>
$3 \times 10^{-12}$ to $3 \times 10^{-9}$	95	8 BECHTLE	12 THEO	
$< 1.6 \times 10^{-7}$		9 BEHNKE	12 COUP	CF <sub>3</sub> I
$< 10^{-9}$	95	10 BUCHMUEL...	12A THEO	
$< 6.5 \times 10^{-6}$		11 FELIZARDO	12 SMPL	C <sub>2</sub> ClF <sub>5</sub>
$< 10^{-9}$	95	12 FOWLIE	12A THEO	
$< 2.3 \times 10^{-7}$	90	13 KIM	12 KIMS	CsI
$< 3.3 \times 10^{-8}$	90	14 AHMED	11A	Ge
$< 4.4 \times 10^{-8}$	90	15 ARMENGAUD	11 EDE2	Ge
$< 4 \times 10^{-8}$	90	16 AHMED	10 CDMS	Ge
$< 7 \times 10^{-7}$	90	17 ANGLOHER	09 CRES	CaWO <sub>4</sub>
$< 1 \times 10^{-7}$	90	18 ANGLE	08 XE10	Xe
$< 1 \times 10^{-6}$	90	BENETTI	08 WARP	Ar
$< 7.5 \times 10^{-7}$	90	19 ALNER	07A ZEP2	Xe
$< 2 \times 10^{-7}$		20 AKERIB	06A CDMS	Ge
$< 90 \times 10^{-7}$		ALNER	05 NAIA	Nal Spin Indep.
$< 12 \times 10^{-7}$		21 ALNER	05A ZEPL	
$< 20 \times 10^{-7}$		22 ANGLOHER	05 CRES	CaWO <sub>4</sub>
$< 14 \times 10^{-7}$		SANGLARD	05 EDEL	Ge
$< 4 \times 10^{-7}$		23 AKERIB	04 CDMS	Ge
$2 \times 10^{-11}$ to $1.5 \times 10^{-7}$	95	24 BALTZ	04 THEO	
$2 \times 10^{-11}$ to $8 \times 10^{-6}$		25,26 ELLIS	04 THEO	$\mu > 0$
$< 5 \times 10^{-8}$		27 PIERCE	04A THEO	
$< 2 \times 10^{-5}$		28 AHMED	03 NAIA	Nal Spin Indep.
$< 3 \times 10^{-6}$		29 AKERIB	03 CDMS	Ge
$2 \times 10^{-13}$ to $2 \times 10^{-7}$		30 BAER	03A THEO	
$< 1.4 \times 10^{-5}$		31 KLAPDOR-K...	03 HDMS	Ge
$< 6 \times 10^{-6}$		32 ABRAMS	02 CDMS	Ge
$< 1.4 \times 10^{-6}$		33 BENOIT	02 EDEL	Ge
$1 \times 10^{-12}$ to $7 \times 10^{-6}$		25 KIM	02B THEO	
$< 3 \times 10^{-5}$		34 MORALES	02B CSME	Ge
$< 1 \times 10^{-5}$		35 MORALES	02C IGEX	Ge
$< 1 \times 10^{-6}$		BALTZ	01 THEO	
$< 3 \times 10^{-5}$	36	BAUDIS	01 HDMS	Ge
$< 4.5 \times 10^{-6}$		BENOIT	01 EDEL	Ge
$< 7 \times 10^{-6}$	37	BOTTINO	01 THEO	
$< 1 \times 10^{-8}$	38	CORSETTI	01 THEO	$\tan\beta \leq 25$
$5 \times 10^{-10}$ to $1.5 \times 10^{-8}$	39	ELLIS	01C THEO	$\tan\beta \leq 10$
$< 4 \times 10^{-6}$	38	GOMEZ	01 THEO	
$2 \times 10^{-10}$ to $1 \times 10^{-7}$	38	LAHANAS	01 THEO	

$< 3 \times 10^{-6}$	ABUSAIDI	00	CDMS	Ge, Si
$< 6 \times 10^{-7}$	40 ACCOMANDO	00	THEO	
	41 BERNABEI	00	DAMA	Nal
$2.5 \times 10^{-9}$ to $3.5 \times 10^{-8}$	42 FENG	00	THEO	$\tan\beta=10$
$< 1.5 \times 10^{-5}$	MORALES	00	IGEX	Ge
$< 4 \times 10^{-5}$	SPOONER	00	UKDM	Nal
$< 7 \times 10^{-6}$	BAUDIS	99	HDMO	$^{76}\text{Ge}$
	43 BERNABEI	99	DAMA	Nal
	44 BERNABEI	98	DAMA	Nal
$< 7 \times 10^{-6}$	BERNABEI	98C	DAMA	Xe

<sup>1</sup> AALSETH 13 presents 90% CL limits on the elastic cross section for masses in the range 4–25 GeV in addition to a region of interest at about 8 GeV. The strongest upper limit is  $2 \times 10^{-5}$  pb at  $m_\chi = 14$  GeV.

<sup>2</sup> AGNESE 13 presents 90% CL limits on the elastic cross section for masses in the range 7–100 GeV using the Si based detector. The strongest upper limit is  $1.8 \times 10^{-6}$  pb at  $m_\chi = 50$  GeV. This limit is improved to  $7 \times 10^{-7}$  pb in AGNESE 13A.

<sup>3</sup> LI 13B presents 90% CL limits on the elastic cross section for masses in the range 4–40 GeV. The strongest upper limit is  $4 \times 10^{-5}$  pb at  $m_\chi = 14$  GeV.

<sup>4</sup> This result updates LEBEDENKO 09. The strongest limit is  $3.9 \times 10^{-8}$  pb at  $m_\chi = 52$  GeV.

<sup>5</sup> ANGLOHER 12 presents results of 730 kg days from the CRESST-II dark matter detector. They find two maxima in the likelihood function corresponding to best fit WIMP masses of 25.3 and 11.6 GeV with elastic cross sections of  $1.6 \times 10^{-6}$  and  $3.7 \times 10^{-5}$  pb respectively, see their Table 4. The statistical significance is more than  $4\sigma$ .

<sup>6</sup> APRILE 12 updates the result of APRILE 11B. The strongest upper limit is  $< 2.0 \times 10^{-9}$  pb and occurs at  $m_\chi \simeq 50$  GeV.

<sup>7</sup> The strongest limit is  $6.1 \times 10^{-5}$  pb at  $m_\chi = 20$  GeV.

<sup>8</sup> Predictions for the spin-independent elastic cross section based on a frequentist approach to electroweak observables in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry using the  $5 \text{ fb}^{-1}$  LHC data and XENON100.

<sup>9</sup> The strongest limit is  $1.4 \times 10^{-7}$  at  $m_\chi = 60$  GeV.

<sup>10</sup> Predictions for the spin-independent elastic cross section based on a frequentist approach to electroweak observables in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry using the  $5 \text{ fb}^{-1}$  LHC data and XENON100.

<sup>11</sup> The strongest limit is  $4.7 \times 10^{-6}$  at  $m_\chi = 35$  GeV.

<sup>12</sup> Predictions for the spin-independent elastic cross section based on a Bayesian approach to electroweak observables in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry using the  $5 \text{ fb}^{-1}$  LHC data and XENON100.

<sup>13</sup> This result updates LEE 07A. The strongest limit is  $2.1 \times 10^{-7}$  at  $m_\chi = 70$  GeV.

<sup>14</sup> AHMED 11A gives combined results from CDMS and EDELWEISS. The strongest limit is at  $m_\chi = 90$  GeV.

<sup>15</sup> ARMENGAUD 11 updates result of ARMENGAUD 10. Strongest limit at  $m_\chi = 85$  GeV.

<sup>16</sup> The strongest upper limit is  $< 3.8 \times 10^{-8}$  pb and occurs at  $m_\chi \simeq 70$  GeV. AHMED 10 updates the results of AHMED 09.

<sup>17</sup> The strongest upper limit is  $4.8 \times 10^{-7}$  pb and occurs at  $m_\chi = 50$  GeV.

<sup>18</sup> The strongest upper limit is  $5.1 \times 10^{-8}$  pb and occurs at  $m_\chi \simeq 30$  GeV. The values quoted here are based on the analysis performed in ANGLE 08 with the update from SORENSEN 09.

<sup>19</sup> The strongest upper limit is  $6.6 \times 10^{-7}$  pb and occurs at  $m_\chi \simeq 65$  GeV.

- 20 AKERIB 06A updates the results of AKERIB 05. The strongest upper limit is  $1.6 \times 10^{-7}$  pb and occurs at  $m_\chi \approx 60$  GeV.
- 21 The strongest upper limit is also close to  $1.0 \times 10^{-6}$  pb and occurs at  $m_\chi \approx 70$  GeV. BENOIT 06 claim that the discrimination power of ZEPLIN-I measurement (ALNER 05A) is not reliable enough to obtain a limit better than  $1 \times 10^{-3}$  pb. However, SMITH 06 do not agree with the criticisms of BENOIT 06.
- 22 The strongest upper limit is also close to  $1.4 \times 10^{-6}$  pb and occurs at  $m_\chi \approx 70$  GeV.
- 23 AKERIB 04 is incompatible with BERNABEI 00 most likely value, under the assumption of standard WIMP-halo interactions. The strongest upper limit is  $4 \times 10^{-7}$  pb and occurs at  $m_\chi \approx 60$  GeV.
- 24 Predictions for the spin-independent elastic cross section in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry.
- 25 KIM 02 and ELLIS 04 calculate the  $\chi p$  elastic scattering cross section in the framework of  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry, but without universal scalar masses.
- 26 In the case of universal squark and slepton masses, but non-universal Higgs masses, the limit becomes  $2 \times 10^{-6}$  ( $2 \times 10^{-11}$  when constraint from the BNL  $g-2$  experiment are included), see ELLIS 03E. ELLIS 05 display the sensitivity of the elastic scattering cross section to the  $\pi$ -Nucleon  $\Sigma$  term.
- 27 PIERCE 04A calculates the  $\chi p$  elastic scattering cross section in the framework of models with very heavy scalar masses. See Fig. 2 of the paper.
- 28 The strongest upper limit is  $1.8 \times 10^{-5}$  pb and occurs at  $m_\chi \approx 80$  GeV.
- 29 Under the assumption of standard WIMP-halo interactions, Akerib 03 is incompatible with BERNABEI 00 most likely value at the 99.98% CL. See Fig. 4.
- 30 BAER 03A calculates the  $\chi p$  elastic scattering cross section in several models including the framework of  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry.
- 31 The strongest upper limit is  $7 \times 10^{-6}$  pb and occurs at  $m_\chi \approx 30$  GeV.
- 32 ABRAMS 02 is incompatible with the DAMA most likely value at the 99.9% CL. The strongest upper limit is  $3 \times 10^{-6}$  pb and occurs at  $m_\chi \approx 30$  GeV.
- 33 BENOIT 02 excludes the central result of DAMA at the 99.8%CL.
- 34 The strongest upper limit is  $2 \times 10^{-5}$  pb and occurs at  $m_\chi \approx 40$  GeV.
- 35 The strongest upper limit is  $7 \times 10^{-6}$  pb and occurs at  $m_\chi \approx 46$  GeV.
- 36 The strongest upper limit is  $1.8 \times 10^{-5}$  pb and occurs at  $m_\chi \approx 32$  GeV
- 37 BOTTINO 01 calculates the  $\chi$ - $p$  elastic scattering cross section in the framework of the following supersymmetric models:  $N=1$  supergravity with the radiative breaking of the electroweak gauge symmetry,  $N=1$  supergravity with nonuniversal scalar masses and an effective MSSM model at the electroweak scale.
- 38 Calculates the  $\chi$ - $p$  elastic scattering cross section in the framework of  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry.
- 39 ELLIS 01C calculates the  $\chi$ - $p$  elastic scattering cross section in the framework of  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry. ELLIS 02B find a range  $2 \times 10^{-8}$ – $1.5 \times 10^{-7}$  at  $\tan\beta=50$ . In models with nonuniversal Higgs masses, the upper limit to the cross section is  $4 \times 10^{-7}$ .
- 40 ACCOMANDO 00 calculate the  $\chi$ - $p$  elastic scattering cross section in the framework of minimal  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry. The limit is relaxed by at least an order of magnitude when models with nonuniversal scalar masses are considered. A subset of the authors in ARNOWITT 02 updated the limit to  $< 9 \times 10^{-8}$  ( $\tan\beta < 55$ ).
- 41 BERNABEI 00 search for annual modulation of the WIMP signal. The data favor the hypothesis of annual modulation at  $4\sigma$  and are consistent, for a particular model framework quoted there, with  $m_{\chi_0} = 44^{+12}_{-9}$  GeV and a spin-independent  $\chi^0$ -proton cross section of  $(5.4 \pm 1.0) \times 10^{-6}$  pb. See also BERNABEI 01 and BERNABEI 00C.

- 42 FENG 00 calculate the  $\chi$ - $p$  elastic scattering cross section in the framework of  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry with a particular emphasis on focus point models. At  $\tan\beta=50$ , the range is  $8 \times 10^{-8} - 4 \times 10^{-7}$ .
- 43 BERNABEI 99 search for annual modulation of the WIMP signal. The data favor the hypothesis of annual modulation at 99.6%CL and are consistent, for the particular model framework considered there, with  $m_{\chi_0} = 59_{-14}^{+17}$  GeV and spin-independent  $\chi^0$ -proton cross section of  $(7.0_{-1.2}^{+0.4}) \times 10^{-6}$  pb ( $1\sigma$  errors).
- 44 BERNABEI 98 search for annual modulation of the WIMP signal. The data are consistent, for the particular model framework considered there, with  $m_{\chi_0} = 59_{-19}^{+36}$  GeV and spin-independent  $\chi^0$ -proton cross section of  $(1.0_{-0.4}^{+0.1}) \times 10^{-5}$  pb ( $1\sigma$  errors).

### Other bounds on $\tilde{\chi}_1^0$ from astrophysics and cosmology

Most of these papers generally exclude regions in the  $M_2 - \mu$  parameter plane by requiring that the  $\tilde{\chi}_1^0$  contribution to the overall cosmological density is less than some maximal value to avoid overclosure of the Universe. Those not based on the cosmological density are indicated. Many of these papers also include LEP and/or other bounds.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>&gt;46 GeV</b>	1 ELLIS	00	RVUE
• • • We do not use the following data for averages, fits, limits, etc. • • •			
	2 CABRERA	13	COSM
	3 ELLIS	13B	COSM
	2 STREGE	13	COSM
	4 AKULA	12	COSM
	4 ARBEY	12A	COSM
	4 BAER	12	COSM
	5 BALAZS	12	COSM
	6 BECHTLE	12	COSM
	7 BESKIDT	12	COSM
> 18 GeV	8 BOTTINO	12	COSM
	4 BUCHMUEL...	12	COSM
	9 BUCHMUEL...	12A	COSM
	4 CAO	12A	COSM
	4 ELLIS	12B	COSM
	10 FENG	12B	COSM
	11 FOWLIE	12A	COSM
	4 KADASTIK	12	COSM
	5 STREGE	12	COSM
	12 BUCHMUEL...	11	COSM
	13 ROSZKOWSKI	11	COSM
	14 ELLIS	10	COSM
	15 BUCHMUEL...	09	COSM
	16 DREINER	09	THEO
	17 BUCHMUEL...	08	COSM
	13 ELLIS	08	COSM
	18 CALIBBI	07	COSM
	19 ELLIS	07	COSM
	20 ALLANACH	06	COSM
	21 DE-AUSTRI	06	COSM

	13	BAER	05	COSM	
	22	BALTZ	04	COSM	
> 6 GeV	8,23	BELANGER	04	THEO	
	24	ELLIS	04B	COSM	
	25	PIERCE	04A	COSM	
	26	BAER	03	COSM	
> 6 GeV	8	BOTTINO	03	COSM	
	26	CHATTOPAD...	03	COSM	
	27	ELLIS	03	COSM	
	13	ELLIS	03B	COSM	
	26	ELLIS	03C	COSM	
	26	LAHANAS	03	COSM	
	28	LAHANAS	02	COSM	
	29	BARGER	01C	COSM	
	30	ELLIS	01B	COSM	
	27	BOEHM	00B	COSM	
	31	FENG	00	COSM	
< 600 GeV	32	ELLIS	98B	COSM	
	33	EDSJO	97	COSM	Co-annihilation
	34	BAER	96	COSM	
	13	BEREZINSKY	95	COSM	
	35	FALK	95	COSM	$CP$ -violating phases
	36	DREES	93	COSM	Minimal supergravity
	37	FALK	93	COSM	Sfermion mixing
	36	KELLEY	93	COSM	Minimal supergravity
	38	MIZUTA	93	COSM	Co-annihilation
	39	LOPEZ	92	COSM	Minimal supergravity, $m_0=A=0$
	40	MCDONALD	92	COSM	
	41	GRIEST	91	COSM	
	42	NOJIRI	91	COSM	Minimal supergravity
	43	OLIVE	91	COSM	
	44	ROSZKOWSKI	91	COSM	
	45	GRIEST	90	COSM	
	43	OLIVE	89	COSM	
none 100 eV – 15 GeV		SREDNICKI	88	COSM	$\tilde{\gamma}$ ; $m_{\tilde{f}}=100$ GeV
none 100 eV–5 GeV		ELLIS	84	COSM	$\tilde{\gamma}$ ; for $m_{\tilde{f}}=100$ GeV
		GOLDBERG	83	COSM	$\tilde{\gamma}$
	46	KRAUSS	83	COSM	$\tilde{\gamma}$
		VYSOTSKII	83	COSM	$\tilde{\gamma}$

<sup>1</sup> ELLIS 00 updates ELLIS 98. Uses LEP  $e^+e^-$  data at  $\sqrt{s}=202$  and 204 GeV to improve bound on neutralino mass to 51 GeV when scalar mass universality is assumed and 46 GeV when Higgs mass universality is relaxed. Limits on  $\tan\beta$  improve to  $> 2.7$  ( $\mu > 0$ ),  $> 2.2$  ( $\mu < 0$ ) when scalar mass universality is assumed and  $> 1.9$  (both signs of  $\mu$ ) when Higgs mass universality is relaxed.

<sup>2</sup> CABRERA 13 and STREGE 13 place constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry with and without non-universal Higgs masses using the  $5.8 \text{ fb}^{-1}$ ,  $\sqrt{s} = 7$  TeV ATLAS supersymmetry searches and XENON100 results.

- <sup>3</sup> ELLIS 13B place constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry with and without Higgs mass universality. Models with universality below the GUT scale are also considered.
- <sup>4</sup> Implications of the LHC result on the Higgs mass and on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry.
- <sup>5</sup> BALAZS 12 and STREGE 12 place constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry using the  $1 \text{ fb}^{-1}$  LHC supersymmetry searches, the  $5 \text{ fb}^{-1}$  Higgs mass constraints, both with  $\sqrt{s} = 7 \text{ TeV}$ , and XENON100 results.
- <sup>6</sup> BECHTLE 12 places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry using indirect experimental searches, using the  $5 \text{ fb}^{-1}$  LHC and XENON100 data.
- <sup>7</sup> BESKIDT 12 places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry using indirect experimental searches, the  $5 \text{ fb}^{-1}$  LHC and the XENON100 data.
- <sup>8</sup> BELANGER 04 and BOTTINO 12 (see also BOTTINO 03, BOTTINO 03A and BOTTINO 04) do not assume gaugino or scalar mass unification.
- <sup>9</sup> BUCHMUELLER 12A places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry using indirect experimental searches, using the  $5 \text{ fb}^{-1}$  LHC and XENON100 data.
- <sup>10</sup> FENG 12B places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry and large sfermion masses using the  $1 \text{ fb}^{-1}$  LHC supersymmetry searches, the  $5 \text{ fb}^{-1}$  LHC Higgs mass constraints both with  $\sqrt{s} = 7 \text{ TeV}$ , and XENON100 results.
- <sup>11</sup> FOWLIE 12A places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry using indirect experimental searches, using the  $5 \text{ fb}^{-1}$  LHC and XENON100 data.
- <sup>12</sup> BUCHMUELLER 11 places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry using indirect experimental searches and including supersymmetry breaking relations between A and B parameters.
- <sup>13</sup> Places constraints on the SUSY parameter space in the framework of  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry but non-Universal Higgs masses.
- <sup>14</sup> ELLIS 10 places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry with universality above the GUT scale.
- <sup>15</sup> BUCHMUELLER 09 places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry using indirect experimental searches.
- <sup>16</sup> DREINER 09 show that in the general MSSM with non-universal gaugino masses there exists no model-independent laboratory bound on the mass of the lightest neutralino. An essentially massless  $\chi_1^0$  is allowed by the experimental and observational data, imposing some constraints on other MSSM parameters, including  $M_2$ ,  $\mu$  and the slepton and squark masses.
- <sup>17</sup> BUCHMUELLER 08 places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry using indirect experimental searches.
- <sup>18</sup> CALIBBI 07 places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry with universality above the GUT scale including the effects of right-handed neutrinos.
- <sup>19</sup> ELLIS 07 places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry with universality below the GUT scale.

- 20 ALLANACH 06 places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry.
- 21 DE-AUSTRI 06 places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry.
- 22 BALTZ 04 places constraints on the SUSY parameter space in the framework of  $N = 1$  supergravity models with radiative breaking of the electroweak gauge symmetry.
- 23 Limit assumes a pseudo scalar mass  $< 200$  GeV. For larger pseudo scalar masses,  $m_{\chi} > 18(29)$  GeV for  $\tan\beta = 50(10)$ . Bounds from WMAP,  $(g - 2)_{\mu}$ ,  $b \rightarrow s\gamma$ , LEP.
- 24 ELLIS 04B places constraints on the SUSY parameter space in the framework of  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry including supersymmetry breaking relations between A and B parameters. See also ELLIS 03D.
- 25 PIERCE 04A places constraints on the SUSY parameter space in the framework of models with very heavy scalar masses.
- 26 BAER 03, CHATTOPADHYAY 03, ELLIS 03C and LAHANAS 03 place constraints on the SUSY parameter space in the framework of  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry based on WMAP results for the cold dark matter density.
- 27 BOEHM 00B and ELLIS 03 place constraints on the SUSY parameter space in the framework of minimal  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry. Includes the effect of  $\chi$ - $\tilde{t}$  co-annihilations.
- 28 LAHANAS 02 places constraints on the SUSY parameter space in the framework of minimal  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry. Focuses on the role of pseudo-scalar Higgs exchange.
- 29 BARGER 01C use the cosmic relic density inferred from recent CMB measurements to constrain the parameter space in the framework of minimal  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry.
- 30 ELLIS 01B places constraints on the SUSY parameter space in the framework of minimal  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry. Focuses on models with large  $\tan\beta$ .
- 31 FENG 00 explores cosmologically allowed regions of MSSM parameter space with multi-TeV masses.
- 32 ELLIS 98B assumes a universal scalar mass and radiative supersymmetry breaking with universal gaugino masses. The upper limit to the LSP mass is increased due to the inclusion of  $\chi - \tilde{\tau}_R$  coannihilations.
- 33 EDSJO 97 included all coannihilation processes between neutralinos and charginos for any neutralino mass and composition.
- 34 Notes the location of the neutralino  $Z$  resonance and  $h$  resonance annihilation corridors in minimal supergravity models with radiative electroweak breaking.
- 35 Mass of the bino (=LSP) is limited to  $m_{\tilde{B}} \lesssim 350$  GeV for  $m_t = 174$  GeV.
- 36 DREES 93, KELLEY 93 compute the cosmic relic density of the LSP in the framework of minimal  $N=1$  supergravity models with radiative breaking of the electroweak gauge symmetry.
- 37 FALK 93 relax the upper limit to the LSP mass by considering sfermion mixing in the MSSM.
- 38 MIZUTA 93 include coannihilations to compute the relic density of Higgsino dark matter.
- 39 LOPEZ 92 calculate the relic LSP density in a minimal SUSY GUT model.
- 40 MCDONALD 92 calculate the relic LSP density in the MSSM including exact tree-level annihilation cross sections for all two-body final states.
- 41 GRIEST 91 improve relic density calculations to account for coannihilations, pole effects, and threshold effects.
- 42 NOJIRI 91 uses minimal supergravity mass relations between squarks and sleptons to narrow cosmologically allowed parameter space.
- 43 Mass of the bino (=LSP) is limited to  $m_{\tilde{B}} \lesssim 350$  GeV for  $m_t \leq 200$  GeV. Mass of the higgsino (=LSP) is limited to  $m_{\tilde{H}} \lesssim 1$  TeV for  $m_t \leq 200$  GeV.

- 44 ROSZKOWSKI 91 calculates LSP relic density in mixed gaugino/higgsino region.  
 45 Mass of the bino (=LSP) is limited to  $m_{\tilde{B}} \lesssim 550$  GeV. Mass of the higgsino (=LSP) is limited to  $m_{\tilde{H}} \lesssim 3.2$  TeV.  
 46 KRAUSS 83 finds  $m_{\tilde{\gamma}}$  not 30 eV to 2.5 GeV. KRAUSS 83 takes into account the gravitino decay. Find that limits depend strongly on reheated temperature. For example a new allowed region  $m_{\tilde{\gamma}} = 4\text{--}20$  MeV exists if  $m_{\text{gravitino}} < 40$  TeV. See figure 2.

## ———— Unstable $\tilde{\chi}_1^0$ (Lightest Neutralino) MASS LIMIT ————

Unless otherwise stated, results in this section assume spectra and production rates as evaluated in the MSSM. Unless otherwise stated, the goldstino or gravitino mass  $m_{\tilde{G}}$  is assumed to be negligible relative to all other masses. In the following,  $\tilde{G}$  is assumed to be undetected and to give rise to a missing energy ( $\cancel{E}$ ) signature.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
none	220–380	95	1 AAD	13AP ATLS $2\gamma + \cancel{E}_T$ , GMSB, SPS8
			2 AAD	13Q ATLS $\gamma + b + \cancel{E}_T$ , higgsino-like neutralino, GMSB
>220	95	3 AAD	13R ATLS $\tilde{\chi}_1^0 \rightarrow \mu jj, \cancel{R}, \lambda'_{211} \neq 0$	
		4 AALTONEN	13I CDF $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}, \cancel{E}_T$ , GMSB	
		5 CHATRCHYAN	13AH CMS $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ , GMSB, SPS8, $c\tau < 500$ mm	
		6 AAD	12CP ATLS $2\gamma + \cancel{E}_T$ , GMSB	
>149	95	7 AAD	12CT ATLS $\geq 4\ell^\pm, \cancel{R}$	
		8 AAD	12R ATLS $\tilde{\chi}_1^0 \rightarrow \mu jj, \cancel{R}, \lambda'_{211} \neq 0$	
		9 ABAZOV	12AD D0 $\tilde{\chi}_1^0 \tilde{\chi}_1^0 \rightarrow \gamma Z \tilde{G} \tilde{G}$ , GMSB	
		10 CHATRCHYAN	12BK CMS $2\gamma + \cancel{E}_T$ , GMSB	
		11 CHATRCHYAN	11B CMS $\tilde{W}^0 \rightarrow \gamma \tilde{G}, \tilde{W}^\pm \rightarrow \ell^\pm \tilde{G}$ , GMSB	
		12 AALTONEN	10 CDF $p\bar{p} \rightarrow \tilde{\chi}\tilde{\chi}, \tilde{\chi}=\tilde{\chi}_2^0, \tilde{\chi}_1^\pm, \tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ , GMSB	
		13 ABAZOV	10P D0 $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ , GMSB	
		14 AALTONEN	08U CDF $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ , GMSB	
		15 ABAZOV	08F D0 $p\bar{p} \rightarrow \tilde{\chi}\tilde{\chi}, \tilde{\chi}=\tilde{\chi}_2^0, \tilde{\chi}_1^\pm, \tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ , GMSB	
		16 ABAZOV	08X D0 $\tilde{\chi}_1^0 \rightarrow Z^0 \tilde{G}$ , GMSB	
> 96.8	95	17 ABULENCIA	07H CDF $\cancel{R}, LL\bar{E}$	
		18 ABAZOV	06D D0 $\cancel{R}, LL\bar{E}$	
		19 ABAZOV	06P D0 $\cancel{R}, \lambda_{122}$	
> 96	95	20 ABBIENDI	06B OPAL $e^+e^- \rightarrow \tilde{B}\tilde{B}, (\tilde{B} \rightarrow \tilde{G}\gamma)$	
		21 ABDALLAH	05B DLPH $e^+e^- \rightarrow \tilde{G}\tilde{\chi}_1^0, (\tilde{\chi}_1^0 \rightarrow \tilde{G}\gamma)$	
> 93	95	22 ABDALLAH	05B DLPH $e^+e^- \rightarrow \tilde{B}\tilde{B}, (\tilde{B} \rightarrow \tilde{G}\gamma)$	
> 93	95	23 ACOSTA	05E CDF $p\bar{p} \rightarrow \tilde{\chi}\tilde{\chi}, \tilde{\chi}=\tilde{\chi}_2^0, \tilde{\chi}_1^\pm, \tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ , GMSB	
		24 AKTAS	05 H1 $e^\pm p \rightarrow q \tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ , GMSB+ $\cancel{R} LQ\bar{D}$	

		25	ABBIENDI	04N	OPAL	$e^+e^- \rightarrow \gamma\gamma\cancel{E}$
> 66		95 26,27	ABDALLAH	04H	DLPH	$\text{AMSB}, \mu > 0$
> 38.0		95 28,29	ABDALLAH	04M	DLPH	$\cancel{R}(\bar{U}\bar{D}\bar{D})$
		30	ACHARD	04E	L3	$e^+e^- \rightarrow \tilde{G}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tilde{G}\gamma$
> 99.5		95 31	ACHARD	04E	L3	$e^+e^- \rightarrow \tilde{B}\tilde{B}, (\tilde{B} \rightarrow \tilde{G}\gamma)$
> 89		32	ABDALLAH	03D	DLPH	$e^+e^- \rightarrow \tilde{\chi}_1^0\tilde{\chi}_1^0, \text{GMSB},$ $m(\tilde{G}) < 1\text{eV}$
		33	HEISTER	03C	ALEP	$e^+e^- \rightarrow \tilde{B}\tilde{B}, (\tilde{B} \rightarrow \gamma\tilde{G})$
		34	HEISTER	03C	ALEP	$e^+e^- \rightarrow \tilde{G}\tilde{\chi}_1^0, (\tilde{\chi}_1^0 \rightarrow \tilde{G}\gamma)$
> 39.9		95 35	ACHARD	02	L3	$\cancel{R}, \text{MSUGRA}$
> 92		95 36	HEISTER	02R	ALEP	short lifetime
> 54		95 36	HEISTER	02R	ALEP	any lifetime
> 85		95 37	ABBIENDI	01	OPAL	$e^+e^- \rightarrow \tilde{\chi}_1^0\tilde{\chi}_1^0, \text{GMSB}, \tan\beta=2$
> 76		95 37	ABBIENDI	01	OPAL	$e^+e^- \rightarrow \tilde{\chi}_1^0\tilde{\chi}_1^0, \text{GMSB}, \tan\beta=20$
> 32.5		95 38	ACCIARRI	01	L3	$\cancel{R}, \text{all } m_0, 0.7 \leq \tan\beta \leq 40$
		39	ADAMS	01	NTEV	$\tilde{\chi}^0 \rightarrow \mu\mu\nu, \cancel{R}, \text{LL}\bar{E}$
> 29		95 40	ABBIENDI	99T	OPAL	$e^+e^- \rightarrow \tilde{\chi}_1^0\tilde{\chi}_1^0, \cancel{R}, m_0=500 \text{ GeV},$ $\tan\beta > 1.2$
> 29		95 41	BARATE	99E	ALEP	$\cancel{R}, \text{LQ}\bar{D}, \tan\beta=1.41, m_0=500 \text{ GeV}$
		42	ABREU	98	DLPH	$e^+e^- \rightarrow \tilde{\chi}_1^0\tilde{\chi}_1^0 (\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G})$
> 23		95 43	BARATE	98S	ALEP	$\cancel{R}, \text{LL}\bar{E}$
		44	ELLIS	97	THEO	$e^+e^- \rightarrow \tilde{\chi}_1^0\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$
		45	CABIBBO	81	COSM	

<sup>1</sup> AAD 13AP searched in  $4.8 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing non-pointing photons in a diphoton plus missing transverse energy final state. No excess is observed above the background expected from Standard Model processes. The results are used to set 95% C.L. exclusion limits in the context of gauge-mediated supersymmetric breaking models, with the lightest neutralino being the next-to-lightest supersymmetric particle and decaying with a lifetime in excess of 0.25 ns into a photon and a gravitino. For limits in the NLSP lifetime versus  $\Lambda$  plane, for the SPS8 model, see their Fig. 8.

<sup>2</sup> AAD 13Q searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing a high- $p_T$  isolated photon, at least one jet identified as originating from a bottom quark, and high missing transverse momentum. Such signatures may originate from supersymmetric models with gauge-mediated supersymmetry breaking in events in which one of a pair of higgsino-like neutralinos decays into a photon and a gravitino while the other decays into a Higgs boson and a gravitino. No significant excess above the expected background was found and limits were set on the neutralino mass in a generalized GMSB model (GGM) with a higgsino-like neutralino NLSP, see their Fig. 4. Intermediate neutralino masses between 220 and 380 GeV are excluded at 95% C.L. regardless of the squark and gluino masses, purely on the basis of the expected weak production.

<sup>3</sup> AAD 13R looked in  $4.4 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing new, heavy particles that decay at a significant distance from their production point into a final state containing a high-momentum muon and charged hadrons. No excess over the expected background is observed and limits are placed on the production cross-section of neutralinos via squarks for various  $m_{\tilde{q}}, m_{\tilde{\chi}_1^0}$  in an R-parity violating scenario with

$\lambda'_{211} \neq 0$ , as a function of the neutralino lifetime, see their Fig. 6.

<sup>4</sup> AALTONEN 13I searched in  $6.3 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events containing  $\cancel{E}_T$  and a delayed photon that arrives late in the detector relative to the time expected from prompt production. No evidence of delayed photon production is observed.

- <sup>5</sup> CHATRCHYAN 13AH searched in  $4.9 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing  $\cancel{E}_T$  and a delayed photon that arrives late in the detector relative to the time expected from prompt production. No significant excess above the expected background was found and limits were set on the pair production of  $\tilde{\chi}_1^0$  depending on the neutralino proper decay length, see Fig. 8. Supersedes CHATRCHYAN 12BK.
- <sup>6</sup> AAD 12CP searched in  $4.8 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with two photons and large  $\cancel{E}_T$  due to  $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$  decays in a GMSB framework. No significant excess above the expected background was found and limits were set on the neutralino mass in a generalized GMSB model (GGM) with a bino-like neutralino NLSP, see Figs. 6 and 7. The other sparticle masses were decoupled,  $\tan\beta = 2$  and  $c\tau_{NLSP} < 0.1 \text{ mm}$ . Also, in the framework of the SPS8 model, limits are presented in Fig. 8.
- <sup>7</sup> AAD 12CT searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing four or more leptons (electrons or muons) and either moderate values of missing transverse momentum or large effective mass. No significant excess is found in the data. Limits are presented in a simplified model of  $R$ -parity violating supersymmetry in which charginos are pair-produced and then decay into a  $W$ -boson and a  $\tilde{\chi}_1^0$ , which in turn decays through an RPV coupling into two charged leptons ( $e^\pm e^\mp$  or  $\mu^\pm \mu^\mp$ ) and a neutrino. In this model, limits are set on the neutralino mass as a function of the chargino mass, see Fig. 3a. Limits are also set in an  $R$ -parity violating mSUGRA model, see Fig. 3b.
- <sup>8</sup> AAD 12R looked in  $33 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing new, heavy particles that decay at a significant distance from their production point into a final state containing a high-momentum muon and charged hadrons. No excess over the expected background is observed and limits are placed on the production cross-section of neutralinos via squarks for various  $(m_{\tilde{q}}, m_{\tilde{\chi}_1^0})$  in an  $R$ -parity violating scenario with  $\lambda'_{211} \neq 0$ , as a function of the neutralino lifetime, see their Fig. 8. Superseded by AAD 13R.
- <sup>9</sup> ABAZOV 12AD looked in  $6.2 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with a photon, a  $Z$ -boson, and large  $\cancel{E}_T$  in the final state. This topology corresponds to a GMSB model where pairs of neutralino NLSPs are either pair produced promptly or from decays of other supersymmetric particles and then decay to either  $Z \tilde{G}$  or  $\gamma \tilde{G}$ . No significant excess over the SM expectation is observed and a limit at 95% C.L. on the cross section is derived as a function of the effective SUSY breaking scale  $\Lambda$ , see Fig. 3. Assuming  $N_{mes} = 2$ ,  $M_{mes} = 3 \Lambda$ ,  $\tan\beta = 3$ ,  $\mu = 0.75 M_1$ , and  $C_{grav} = 1$ , the model is excluded at 95% C.L. for values of  $\Lambda < 87 \text{ TeV}$ .
- <sup>10</sup> CHATRCHYAN 12BK searched in  $2.23 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with two photons and large  $\cancel{E}_T$  due to  $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$  decays in a GMSB framework. No significant excess above the expected background was found and limits were set on the pair production of  $\tilde{\chi}_1^0$  depending on the neutralino lifetime, see Fig. 6.
- <sup>11</sup> CHATRCHYAN 11B looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s}=7 \text{ TeV}$  for events with an isolated lepton ( $e$  or  $\mu$ ), a photon and  $\cancel{E}_T$  which may arise in a generalized gauge mediated model from the decay of Wino-like NLSPs. No evidence for an excess over the expected background is observed. Limits are derived in the plane of squark/gluino mass versus Wino mass (see Fig. 4). Mass degeneracy of the produced squarks and gluinos is assumed.
- <sup>12</sup> AALTONEN 10 searched in  $2.6 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for diphoton events with large  $\cancel{E}_T$ . They may originate from the production of  $\tilde{\chi}^\pm$  in pairs or associated to a  $\tilde{\chi}_2^0$ , decaying into  $\tilde{\chi}_1^0$  which itself decays in GMSB to  $\gamma \tilde{G}$ . There is no excess of events beyond expectation. An upper limit on the cross section is calculated in the GMSB model as a function of the  $\tilde{\chi}_1^0$  mass and lifetime, see their Fig. 2. A limit is derived on the  $\tilde{\chi}_1^0$  mass of  $149 \text{ GeV}$  for  $\tau_{\tilde{\chi}_1^0} \ll 1 \text{ ns}$ , which improves the results of previous searches.

- 13 ABAZOV 10P looked in  $6.3 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with at least two isolated  $\gamma$ s and large  $\cancel{E}_T$ . These could be the signature of  $\tilde{\chi}_2^0$  and  $\tilde{\chi}_1^\pm$  production, decaying to  $\tilde{\chi}_1^0$  and finally  $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$  in a GMSB framework. No significant excess over the SM expectation is observed, and a limit at 95% C.L. on the cross section is derived for  $N_{mes} = 1$ ,  $\tan\beta = 15$  and  $\mu > 0$ , see their Fig. 2. This allows them to set a limit on the effective SUSY breaking scale  $\Lambda > 124 \text{ TeV}$ , from which the excluded  $\tilde{\chi}_1^0$  mass range is obtained.
- 14 AALTONEN 08U searched in  $570 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events that contain a time-delayed photon, at least one jet, and large  $\cancel{E}_T$ . The time-of-arrival is measured for each electromagnetic tower with a resolution of 0.50 ns. The number of observed events in the signal region is consistent with the background estimation. An upper limit on the cross section is derived as a function of the  $\tilde{\chi}_1^0$  mass and lifetime, shown in their Fig. 24. The comparison with the NLO cross section for GMSB yields an exclusion of the  $\tilde{\chi}_1^0$  mass as a function of its lifetime, see Fig. 25. See ABULENCIA 07P for a previous analysis of the same data set.
- 15 ABAZOV 08F looked in  $1.1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for diphoton events with large  $\cancel{E}_T$ . They may originate from the production of  $\tilde{\chi}^\pm$  in pairs or associated to a  $\tilde{\chi}_2^0$ , decaying to a  $\tilde{\chi}_1^0$  which itself decays promptly in GMSB to  $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ . No significant excess was found compared to the background expectation. A limit is derived on the masses of SUSY particles in the GMSB framework for  $M = 2\Lambda$ ,  $N = 1$ ,  $\tan\beta = 15$  and  $\mu > 0$ , see Figure 2. It also excludes  $\Lambda < 91.5 \text{ TeV}$ . Supersedes the results of ABAZOV 05A. Superseded by ABAZOV 10P.
- 16 ABAZOV 08X searched in  $1.1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for an excess of events with electron pairs. Their vertex, reconstructed from the directions measured in the segmented electromagnetic calorimeter, is required to be away from the primary interaction point. Such delayed decays might be expected for a Higgsino-like  $\tilde{\chi}_1^0$  in GMSB. No significant excess was found compared to the background expectation. Upper limits on the cross-section times branching ratio are extracted as a function of the lifetime for several ranges of dielectron invariant masses, see their Fig. 3.
- 17 ABULENCIA 07H searched in  $346 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with at least three leptons ( $e$  or  $\mu$ ) from the decay of  $\tilde{\chi}_1^0$  via  $LL\bar{E}$  couplings. The results are consistent with the hypothesis of no signal. Upper limits on the cross-section are extracted and a limit is derived in the framework of mSUGRA on the masses of  $\tilde{\chi}_1^0$  and  $\tilde{\chi}_1^\pm$ , see e.g. their Fig. 3 and Tab. II.
- 18 ABAZOV 06D looked in  $360 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with three leptons originating from the pair production of charginos and neutralinos, followed by  $\tilde{R}$  decays mediated by  $LL\bar{E}$  couplings. One coupling is assumed to be dominant at a time. No significant excess was found compared to the background expectation in the  $e\ell\ell$ ,  $\mu\mu\ell$  nor  $e\ell\tau$  ( $\ell = e, \mu$ ) final states. Upper limits on the cross-section are extracted in a specific MSUGRA model and a MSSM model without unification of  $M_1$  and  $M_2$  at the GUT scale. A limit is derived on the masses of charginos and neutralinos for both scenarios assuming  $\lambda_{ijk}$  couplings such that the decay length is less than 1 cm, see their Table III and Fig. 4.
- 19 ABAZOV 06P looked in  $380 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with at least 2 opposite sign isolated muons which might arise from the decays of neutralinos into  $\mu\mu\nu$  via  $\tilde{R}$  couplings  $LL\bar{E}$ . No events are observed in the decay region defined by a radius between 5 and 20 cm, in agreement with the SM expectation. Limits are set on the cross-section times branching ratio as a function of lifetime, shown in their Fig. 3. This limit excludes the SUSY interpretation of the NuTeV excess of dimuon events reported in ADAMS 01.
- 20 ABBIENDI 06B use  $600 \text{ pb}^{-1}$  of data from  $\sqrt{s} = 189\text{--}209 \text{ GeV}$ . They look for events with diphotons +  $\cancel{E}$  final states originating from prompt decays of pair-produced neutralinos in a GMSB scenario with  $\tilde{\chi}_1^0$  NLSP. Limits on the cross-section are computed as a function

- of  $m(\tilde{\chi}_1^0)$ , see their Fig. 14. The limit on the  $\tilde{\chi}_1^0$  mass is for a pure Bino state assuming a prompt decay, with lifetimes up to  $10^{-9}$ s. Supersedes the results of ABBIENDI 04N.
- 21 ABDALLAH 05B use data from  $\sqrt{s} = 180\text{--}209$  GeV. They look for events with single photons +  $\cancel{E}$  final states. Limits are computed in the plane  $(m(\tilde{G}), m(\tilde{\chi}_1^0))$ , shown in their Fig. 9b for a pure Bino state in the GMSB framework and in Fig. 9c for a no-scale supergravity model. Supersedes the results of ABREU 00Z.
- 22 ABDALLAH 05B use data from  $\sqrt{s} = 130\text{--}209$  GeV. They look for events with diphotons +  $\cancel{E}$  final states and single photons not pointing to the vertex, expected in GMSB when the  $\tilde{\chi}_1^0$  is the NLSP. Limits are computed in the plane  $(m(\tilde{G}), m(\tilde{\chi}_1^0))$ , see their Fig. 10. The lower limit is derived on the  $\tilde{\chi}_1^0$  mass for a pure Bino state assuming a prompt decay and  $m_{\tilde{e}_R} = m_{\tilde{e}_L} = 2 m_{\tilde{\chi}_1^0}$ . It improves to 100 GeV for  $m_{\tilde{e}_R} = m_{\tilde{e}_L} = 1.1 m_{\tilde{\chi}_1^0}$  and the limit in the plane  $(m(\tilde{\chi}_1^0), m(\tilde{e}_R))$  is shown in Fig. 10b. For long-lived neutralinos, cross-section limits are displayed in their Fig 11. Supersedes the results of ABREU 00Z.
- 23 ACOSTA 05E looked in  $202 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.96$  TeV for diphoton events with large  $\cancel{E}_T$ . They may originate from the production of  $\tilde{\chi}^\pm$  in pairs or associated to a  $\tilde{\chi}_2^0$ , decaying to a  $\tilde{\chi}_1^0$  which itself decays promptly in GMSB to  $\gamma\tilde{G}$ . No events are selected at large  $\cancel{E}_T$  compared to the background expectation. A limit is derived on the masses of SUSY particles in the GMSB framework for  $M = 2 \Lambda$ ,  $N = 1$ ,  $\tan\beta = 15$  and  $\mu > 0$ , see Figure 2. It also excludes  $\Lambda < 69$  TeV. Supersedes the results of ABE 99I.
- 24 AKTAS 05 data collected at 319 GeV with  $64.3 \text{ pb}^{-1}$  of  $e^+p$  and  $13.5 \text{ pb}^{-1}$  of  $e^-p$ . They look for  $\cancel{E}$  resonant  $\tilde{\chi}_1^0$  production via  $t$ -channel exchange of a  $\tilde{e}$ , followed by prompt GMSB decay of the  $\tilde{\chi}_1^0$  to  $\gamma\tilde{G}$ . Upper limits at 95% on the cross section are derived, see their Figure 4, and compared to two example scenarios. In Figure 5, they display 95% exclusion limits in the plane of  $M(\tilde{\chi}_1^0)$  versus  $M(\tilde{e}_L) - M(\tilde{\chi}_1^0)$  for the two scenarios and several values of the  $\lambda'$  Yukawa coupling.
- 25 ABBIENDI 04N use data from  $\sqrt{s} = 189\text{--}209$  GeV, setting limits on  $\sigma(e^+e^- \rightarrow XX) \times B^2(X \rightarrow Y\gamma)$ , with  $Y$  invisible (see their Fig. 4). Limits on  $\tilde{\chi}_1^0$  masses for a specific model are given. Supersedes the results of ABBIENDI,G 00D.
- 26 ABDALLAH 04H use data from LEP 1 and  $\sqrt{s} = 192\text{--}208$  GeV. They re-use results or re-analyze the data from ABDALLAH 03M to put limits on the parameter space of anomaly-mediated supersymmetry breaking (AMSB), which is scanned in the region  $1 < m_{3/2} < 50$  TeV,  $0 < m_0 < 1000$  GeV,  $1.5 < \tan\beta < 35$ , both signs of  $\mu$ . The constraints are obtained from the searches for mass degenerate chargino and neutralino, for SM-like and invisible Higgs, for leptonically decaying charginos and from the limit on non-SM  $Z$  width of 3.2 MeV. The limit is for  $m_t = 174.3$  GeV (see Table 2 for other  $m_t$  values).
- 27 The limit improves to 73 GeV for  $\mu < 0$ .
- 28 ABDALLAH 04M use data from  $\sqrt{s} = 192\text{--}208$  GeV to derive limits on sparticle masses under the assumption of  $\cancel{E}$  with  $LL\bar{E}$  or  $\overline{UD}\bar{D}$  couplings. The results are valid in the ranges  $90 < m_0 < 500$  GeV,  $0.7 < \tan\beta < 30$ ,  $-200 < \mu < 200$  GeV,  $0 < M_2 < 400$  GeV. Supersedes the result of ABREU 01D and ABREU 00U.
- 29 The limit improves to 39.5 GeV for  $LL\bar{E}$  couplings.
- 30 ACHARD 04E use data from  $\sqrt{s} = 189\text{--}209$  GeV. They look for events with single photons +  $\cancel{E}$  final states. Limits are computed in the plane  $(m(\tilde{G}), m(\tilde{\chi}_1^0))$ , shown in their Fig. 8c for a no-scale supergravity model, excluding, e.g., Gravitino masses below  $10^{-5}$  eV for neutralino masses below 172 GeV. Supersedes the results of ACCIARRI 99R.
- 31 ACHARD 04E use data from  $\sqrt{s} = 189\text{--}209$  GeV. They look for events with diphotons +  $\cancel{E}$  final states. Limits are computed in the plane  $(m(\tilde{\chi}_1^0), m(\tilde{e}_R))$ , see their Fig. 8d. The limit on the  $\tilde{\chi}_1^0$  mass is for a pure Bino state assuming a prompt decay, with  $m_{\tilde{e}_L} = 1.1 m_{\tilde{\chi}_1^0}$  and  $m_{\tilde{e}_R} = 2.5 m_{\tilde{\chi}_1^0}$ . Supersedes the results of ACCIARRI 99R.

- <sup>32</sup> ABDALLAH 03D use data from  $\sqrt{s} = 161\text{--}208$  GeV. They look for  $4\text{-tau} + \cancel{E}$  final states, expected in GMSB when the  $\tilde{\tau}_1$  is the NLSP, and  $4\text{-lepton} + \cancel{E}$  final states, expected in the co-NLSP scenario, and assuming a short-lived  $\tilde{\chi}_1^0$  ( $m(\tilde{G}) < 1$  eV). Limits are computed in the plane  $(m(\tilde{\tau}_1), m(\tilde{\chi}_1^0))$  from a scan of the GMSB parameters space, after combining these results with the search for slepton pair production from the same paper to cover prompt decays and for the case of  $\tilde{\chi}_1^0$  NLSP from ABREU 00Z. The limit above is reached for a single generation of messengers and when the  $\tilde{\tau}_1$  is the NLSP. Stronger limits are obtained when more messenger generations are assumed or when the other sleptons are co-NLSP, see their Fig. 10. Supersedes the results of ABREU 01G.
- <sup>33</sup> HEISTER 03C use the data from  $\sqrt{s} = 189\text{--}209$  GeV to search for  $\gamma \cancel{E}_T$  final states with non-pointing photons and  $\gamma\gamma \cancel{E}_T$  events. Interpreted in the framework of Minimal GMSB, a lower bound on the  $\tilde{\chi}_1^0$  mass is obtained as function of its lifetime. For a laboratory lifetime of less than 3 ns, the limit at 95% CL is 98.8 GeV. For other lifetimes, see their Fig. 5. These results are interpreted in a more general GMSB framework in HEISTER 02R.
- <sup>34</sup> HEISTER 03C use the data from  $\sqrt{s} = 189\text{--}209$  GeV to search for  $\gamma \cancel{E}_T$  final states. They obtained an upper bound on the cross section for the process  $e^+e^- \rightarrow \tilde{G}\tilde{\chi}_1^0$ , followed by the prompt decay  $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$ , shown in their Fig. 4. These results supersede BARATE 98H.
- <sup>35</sup> ACHARD 02 searches for the production of sparticles in the case of  $\cancel{R}$  prompt decays with  $LL\bar{E}$  or  $UDD$  couplings at  $\sqrt{s}=189\text{--}208$  GeV. The search is performed for direct and indirect decays, assuming one coupling at the time to be nonzero. The MSUGRA limit results from a scan over the MSSM parameter space with the assumption of gaugino and scalar mass unification at the GUT scale, imposing simultaneously the exclusions from neutralino, chargino, sleptons, and squarks analyses. The limit holds for  $UDD$  couplings and increases to 40.2 GeV for  $LL\bar{E}$  couplings. For L3 limits from  $LQ\bar{D}$  couplings, see ACCIARRI 01.
- <sup>36</sup> HEISTER 02R search for signals of GMSB in the 189–209 GeV data. For the  $\tilde{\chi}_1^0$  NLSP scenario, they looked for topologies consisting of  $\gamma\gamma\cancel{E}$  or a single  $\gamma$  not pointing to the interaction vertex. For the  $\tilde{\ell}$  NLSP case, the topologies consist of  $\ell\ell\cancel{E}$  or  $4\ell\cancel{E}$  (from  $\tilde{\chi}_1^0\tilde{\chi}_1^0$  production), including leptons with large impact parameters, kinks, or stable particles. Limits are derived from a scan over the GMSB parameters (see their Table 5 for the ranges). The limits are valid whichever is the NLSP. The absolute mass bound on the  $\tilde{\chi}_1^0$  for any lifetime includes indirect limits from the chargino search, and from the slepton search HEISTER 02E performed within the MSUGRA framework. A bound for any NLSP and any lifetime of 77 GeV has also been derived by using the constraints from the neutral Higgs search in HEISTER 02. Limits on the universal SUSY mass scale  $\Lambda$  are also derived in the paper. Supersedes the results from BARATE 00G.
- <sup>37</sup> ABBIENDI 01 looked for final states with  $\gamma\gamma\cancel{E}$ ,  $\ell\ell\cancel{E}$ , with possibly additional activity and four leptons +  $\cancel{E}$  to search for prompt decays of  $\tilde{\chi}_1^0$  or  $\tilde{\ell}_1$  in GMSB. They derive limits in the plane  $(m_{\tilde{\chi}_1^0}, m_{\tilde{\tau}_1})$ , see Fig. 6, allowing either the  $\tilde{\chi}_1^0$  or a  $\tilde{\ell}_1$  to be the NLSP. Two scenarios are considered:  $\tan\beta=2$  with the 3 sleptons degenerate in mass and  $\tan\beta=20$  where the  $\tilde{\tau}_1$  is lighter than the other sleptons. Data taken at  $\sqrt{s}=189$  GeV.
- <sup>38</sup> ACCIARRI 01 searches for multi-lepton and/or multi-jet final states from  $\cancel{R}$  prompt decays with  $LL\bar{E}$ ,  $LQ\bar{D}$ , or  $UDD$  couplings at  $\sqrt{s}=189$  GeV. The search is performed for direct and indirect decays of neutralinos, charginos, and scalar leptons, with the  $\tilde{\chi}_1^0$  or a  $\tilde{\ell}$  as LSP and assuming one coupling to be nonzero at a time. Mass limits are derived using simultaneously the constraints from the neutralino, chargino, and slepton analyses; and the  $Z^0$  width measurements from ACCIARRI 00C in a scan of the parameter space assuming MSUGRA with gaugino and scalar mass universality. Updates and supersedes the results from ACCIARRI 99I.

- <sup>39</sup> ADAMS 01 looked for neutral particles with mass  $> 2.2$  GeV, produced by 900 GeV protons incident on a Beryllium oxide target and decaying through weak interactions into  $\mu\mu$ ,  $\mu e$ , or  $\mu\pi$  final states in the decay channel of the NuTeV detector (E815) at Fermilab. The number of observed events is  $3\mu\mu$ ,  $0\mu e$ , and  $0\mu\pi$  with an expected background of  $0.069 \pm 0.010$ ,  $0.13 \pm 0.02$ , and  $0.14 \pm 0.02$ , respectively. The  $\mu\mu$  events are consistent with the  $\tilde{R}$  decay of a neutralino with mass around 5 GeV. However, they share several aspects with  $\nu$ -interaction backgrounds. An upper limit on the differential production cross section of neutralinos in  $p\bar{p}$  interactions as function of the decay length is given in Fig. 3.
- <sup>40</sup> ABBIENDI 99T searches for the production of neutralinos in the case of  $R$ -parity violation with  $LL\bar{E}$ ,  $LQ\bar{D}$ , or  $U\bar{D}\bar{D}$  couplings using data from  $\sqrt{s}=183$  GeV. They investigate topologies with multiple leptons, jets plus leptons, or multiple jets, assuming one coupling at the time to be non-zero and giving rise to direct or indirect decays. Mixed decays (where one particle has a direct, the other an indirect decay) are also considered for the  $U\bar{D}\bar{D}$  couplings. Upper limits on the cross section are derived which, combined with the constraint from the  $Z^0$  width, allow to exclude regions in the  $M_2$  versus  $\mu$  plane for any coupling. Limits on the neutralino mass are obtained for non-zero  $LL\bar{E}$  couplings  $> 10^{-5}$ . The limit disappears for  $\tan\beta < 1.2$  and it improves to 50 GeV for  $\tan\beta > 20$ .
- <sup>41</sup> BARATE 99E looked for the decay of gauginos via  $R$ -violating couplings  $LQ\bar{D}$ . The bound is significantly reduced for smaller values of  $m_0$ . Data collected at  $\sqrt{s}=130$ –172 GeV.
- <sup>42</sup> ABREU 98 uses data at  $\sqrt{s}=161$  and 172 GeV. Upper bounds on  $\gamma\gamma\tilde{E}$  cross section are obtained. Similar limits on  $\gamma\tilde{E}$  are also given, relevant for  $e^+e^- \rightarrow \tilde{\chi}_1^0\tilde{G}$  production.
- <sup>43</sup> BARATE 98S looked for the decay of gauginos via  $R$ -violating coupling  $LL\bar{E}$ . The bound improves to 25 GeV if the chargino decays into neutralino which further decays into lepton pairs. Data collected at  $\sqrt{s}=130$ –172 GeV.
- <sup>44</sup> ELLIS 97 reanalyzed the LEP2 ( $\sqrt{s}=161$  GeV) limits of  $\sigma(\gamma\gamma+E_{\text{miss}}) < 0.2$  pb to exclude  $m_{\tilde{\chi}_1^0} < 63$  GeV if  $m_{\tilde{e}_L} = m_{\tilde{e}_R} < 150$  GeV and  $\tilde{\chi}_1^0$  decays to  $\gamma\tilde{G}$  inside detector.
- <sup>45</sup> CABIBBO 81 consider  $\tilde{\gamma} \rightarrow \gamma + \text{goldstino}$ . Photino must be either light enough ( $< 30$  eV) to satisfy cosmology bound, or heavy enough ( $> 0.3$  MeV) to have disappeared at early universe.

## $\tilde{\chi}_2^0$ , $\tilde{\chi}_3^0$ , $\tilde{\chi}_4^0$ (Neutralinos) MASS LIMITS

Neutralinos are unknown mixtures of photinos, z-inos, and neutral higgsinos (the supersymmetric partners of photons and of  $Z$  and Higgs bosons). The limits here apply only to  $\tilde{\chi}_2^0$ ,  $\tilde{\chi}_3^0$ , and  $\tilde{\chi}_4^0$ .  $\tilde{\chi}_1^0$  is the lightest supersymmetric particle (LSP); see  $\tilde{\chi}_1^0$  Mass Limits. It is not possible to quote rigorous mass limits because they are extremely model dependent; i.e. they depend on branching ratios of various  $\tilde{\chi}^0$  decay modes, on the masses of decay products ( $\tilde{e}$ ,  $\tilde{\gamma}$ ,  $\tilde{q}$ ,  $\tilde{g}$ ), and on the  $\tilde{e}$  mass exchanged in  $e^+e^- \rightarrow \tilde{\chi}_i^0\tilde{\chi}_j^0$ . Limits arise either from direct searches, or from the MSSM constraints set on the gaugino and higgsino mass parameters  $M_2$  and  $\mu$  through searches for lighter charginos and neutralinos. Often limits are given as contour plots in the  $m_{\tilde{\chi}^0} - m_{\tilde{e}}$  plane vs other parameters. When specific assumptions are made, e.g. the neutralino is a pure photino ( $\tilde{\gamma}$ ), pure z-ino ( $\tilde{Z}$ ), or pure neutral higgsino ( $\tilde{H}^0$ ), the neutralinos will be labelled as such.

Limits obtained from  $e^+e^-$  collisions at energies up to 136 GeV, as well as other limits from different techniques, are now superseded and have not been included in this compilation. They can be found in the 1998 Edition (The European Physical Journal **C3** 1 (1998)) of this Review.  $\Delta m = m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$ .

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
		1 AAD	13 ATLS	$3\ell^\pm + \cancel{E}_T$ , pMSSM, SMS
		2 CHATRCHYAN 12BJ	CMS	$\geq 2 \ell$ , jets + $\cancel{E}_T$ , $pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$
> 78	95	3 ABBIENDI	04H OPAL	$\tilde{\chi}_2^0$ , all $\tan\beta$ , $\Delta m > 5$ GeV, $m_0 > 500$ GeV, $A_0 = 0$
> 62.4	95	4 ABREU	00W DLPH	$\tilde{\chi}_2^0$ , $1 \leq \tan\beta \leq 40$ , all $\Delta m$ , all $m_0$
> 99.9	95	4 ABREU	00W DLPH	$\tilde{\chi}_3^0$ , $1 \leq \tan\beta \leq 40$ , all $\Delta m$ , all $m_0$
> 116.0	95	4 ABREU	00W DLPH	$\tilde{\chi}_4^0$ , $1 \leq \tan\beta \leq 40$ , all $\Delta m$ , all $m_0$
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
		5 AAD	12AS ATLS	$3\ell^\pm + \cancel{E}_T$ , pMSSM
		6 AAD	12T ATLS	$\ell^\pm \ell^\pm + \cancel{E}_T$ , $pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$
		7 ABULENCIA	07N CDF	$p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$
		8 ABDALLAH	05B DLPH	$e^+e^- \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_2^0$ , ( $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \gamma$ )
		9 ACHARD	04E L3	$e^+e^- \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_2^0$ , ( $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \gamma$ )
> 80.0	95	10 ACHARD	02 L3	$\tilde{\chi}_2^0$ , $\mathcal{R}$ , MSUGRA
> 107.2	95	10 ACHARD	02 L3	$\tilde{\chi}_3^0$ , $\mathcal{R}$ , MSUGRA
		11 ABREU	01B DLPH	$e^+e^- \rightarrow \tilde{\chi}_i^0 \tilde{\chi}_j^0$
> 68.0	95	12 ACCIARRI	01 L3	$\tilde{\chi}_2^0$ , $\mathcal{R}$ , all $m_0$ , $0.7 \leq \tan\beta \leq 40$
> 99.0	95	12 ACCIARRI	01 L3	$\tilde{\chi}_3^0$ , $\mathcal{R}$ , all $m_0$ , $0.7 \leq \tan\beta \leq 40$
> 50	95	13 ABREU	00U DLPH	$\tilde{\chi}_2^0$ , $\mathcal{R}$ ( $LL\bar{E}$ ), all $\Delta m$ , $1 \leq \tan\beta \leq 30$
		14 ABBIENDI	99F OPAL	$e^+e^- \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^0$ ( $\tilde{\chi}_2^0 \rightarrow \gamma \tilde{\chi}_1^0$ )
		15 ABBIENDI	99F OPAL	$e^+e^- \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_2^0$ ( $\tilde{\chi}_2^0 \rightarrow \gamma \tilde{\chi}_1^0$ )
		16 ABBOTT	98C D0	$p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$
> 82.2	95	17 ABE	98J CDF	$p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$
> 92	95	18 ACCIARRI	98F L3	$\tilde{H}_2^0$ , $\tan\beta=1.41$ , $M_2 < 500$ GeV
		19 ACCIARRI	98V L3	$e^+e^- \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_{1,2}^0$ ( $\tilde{\chi}_2^0 \rightarrow \gamma \tilde{\chi}_1^0$ )
> 53	95	20 BARATE	98H ALEP	$e^+e^- \rightarrow \tilde{\gamma} \tilde{\gamma}$ ( $\tilde{\gamma} \rightarrow \gamma \tilde{H}^0$ )
> 74	95	21 BARATE	98J ALEP	$e^+e^- \rightarrow \tilde{\gamma} \tilde{\gamma}$ ( $\tilde{\gamma} \rightarrow \gamma \tilde{H}^0$ )
		22 ABACHI	96 D0	$p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$
		23 ABE	96K CDF	$p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$

<sup>1</sup> AAD 13 searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for charginos and neutralinos decaying to a final state with three leptons ( $e$  and  $\mu$ ) and missing transverse energy. No excess beyond the Standard Model expectation is observed. Exclusion limits are derived in the phenomenological MSSM, see Fig. 2 and 3, and in simplified models, see Fig. 4. For the simplified models with intermediate slepton decays, degenerate  $\tilde{\chi}_1^\pm$  and  $\tilde{\chi}_2^0$  masses up to 500 GeV are excluded at 95% C.L. for very large mass differences with the  $\tilde{\chi}_1^0$ . Supersedes AAD 12AS.

- <sup>2</sup> CHATRCHYAN 12BJ searched in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for direct electroweak production of charginos and neutralinos in events with at least two leptons, jets and missing transverse momentum. No significant excesses over the expected SM backgrounds are observed and 95% C.L. limits on the production cross section of  $\tilde{\chi}_1^\pm \tilde{\chi}_2^0$  pair production were set in a number of simplified models, see Figs. 7 to 12. Most limits are for exactly 3 jets.
- <sup>3</sup> ABBIENDI 04H search for charginos and neutralinos in events with acoplanar leptons+jets and multi-jet final states in the 192–209 GeV data, combined with the results on leptonic final states from ABBIENDI 04. The results hold for a scan over the parameter space covering the region  $0 < M_2 < 5000 \text{ GeV}$ ,  $-1000 < \mu < 1000 \text{ GeV}$  and  $\tan\beta$  from 1 to 40. This limit supersedes ABBIENDI 00H.
- <sup>4</sup> ABREU 00W combines data collected at  $\sqrt{s}=189 \text{ GeV}$  with results from lower energies. The mass limit is obtained by constraining the MSSM parameter space with gaugino and sfermion mass universality at the GUT scale, using the results of negative direct searches for neutralinos (including cascade decays and  $\tilde{\tau}\tau$  final states) from ABREU 01, for charginos from ABREU 00J and ABREU 00T (for all  $\Delta m_+$ ), and for charged sleptons from ABREU 01B. The results hold for the full parameter space defined by all values of  $M_2$  and  $|\mu| \leq 2 \text{ TeV}$  with the  $\tilde{\chi}_1^0$  as LSP.
- <sup>5</sup> AAD 12AS searched in  $2.06 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for charginos and neutralinos decaying to a final state with three leptons ( $e$  and  $\mu$ ) and missing transverse energy. No excess beyond the Standard Model expectation is observed. Exclusion limits are derived in the phenomenological MSSM, see Fig. 2 (top), and in simplified models, see Fig. 2 (bottom).
- <sup>6</sup> AAD 12T looked in  $1 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for the production of supersymmetric particles decaying into final states with missing transverse momentum and exactly two isolated leptons ( $e$  or  $\mu$ ). Same-sign dilepton events were separately studied. Additionally, in opposite-sign events, a search was made for an excess of same-flavor over different-flavor lepton pairs. No excess over the expected background is observed and limits are placed on the effective production cross section of opposite-sign dilepton events with  $\cancel{E}_T > 250 \text{ GeV}$  and on same-sign dilepton events with  $\cancel{E}_T > 100 \text{ GeV}$ . The latter limit is interpreted in a simplified electroweak gaugino production model.
- <sup>7</sup> ABULENCIA 07N searched in  $1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with two same sign leptons ( $e$  or  $\mu$ ) from the decay of  $\tilde{\chi}_1^\pm \tilde{\chi}_2^0 X$  and large  $\cancel{E}_T$ . A slight excess of 13 events is observed over a SM background expectation of  $7.8 \pm 1.1$ . However, the kinematic distributions do not show any anomalous deviation from expectations in any particular region of parameter space.
- <sup>8</sup> ABDALLAH 05B use data from  $\sqrt{s} = 130\text{--}209 \text{ GeV}$ , looking for events with diphotons +  $\cancel{E}$ . Limits on the cross-section are computed in the plane  $(m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0))$ , see Fig. 12. Supersedes the results of ABREU 00Z.
- <sup>9</sup> ACHARD 04E use data from  $\sqrt{s} = 189\text{--}209 \text{ GeV}$ , looking for events with diphotons +  $\cancel{E}$ . Limits are computed in the plane  $(m(\tilde{\chi}_2^0), m(\tilde{e}_R))$ , for  $\Delta m > 10 \text{ GeV}$ , see Fig. 7. Supersedes the results of ACCIARRI 99R.
- <sup>10</sup> ACHARD 02 searches for the production of sparticles in the case of  $\mathcal{R}$  prompt decays with  $LL\bar{E}$  or  $\overline{UDD}$  couplings at  $\sqrt{s}=189\text{--}208 \text{ GeV}$ . The search is performed for direct and indirect decays, assuming one coupling at the time to be nonzero. The MSUGRA limit results from a scan over the MSSM parameter space with the assumption of gaugino and scalar mass unification at the GUT scale, imposing simultaneously the exclusions from neutralino, chargino, sleptons, and squarks analyses. The limit of  $\tilde{\chi}_2^0$  holds for  $\overline{UDD}$  couplings and increases to  $84.0 \text{ GeV}$  for  $LL\bar{E}$  couplings. The same  $\tilde{\chi}_3^0$  limit holds for both  $LL\bar{E}$  and  $\overline{UDD}$  couplings. For L3 limits from  $LQ\bar{D}$  couplings, see ACCIARRI 01.
- <sup>11</sup> ABREU 01B used data from  $\sqrt{s}=189 \text{ GeV}$  to search for the production of  $\tilde{\chi}_i^0 \tilde{\chi}_j^0$ . They looked for di-jet and di-lepton pairs with  $\cancel{E}$  for events from  $\tilde{\chi}_i^0 \tilde{\chi}_j^0$  with the decay  $\tilde{\chi}_j^0 \rightarrow$

- $f\bar{f}\tilde{\chi}_1^0$ ; multi-jet and multi-lepton pairs with or without additional photons to cover the cascade decays  $\tilde{\chi}_j^0 \rightarrow f\bar{f}\tilde{\chi}_2^0$ , followed by  $\tilde{\chi}_j^0 \rightarrow f\bar{f}\tilde{\chi}_1^0$  or  $\tilde{\chi}_j^0 \rightarrow \gamma\tilde{\chi}_1^0$ ; multi-tau final states from  $\tilde{\chi}_2^0 \rightarrow \tilde{\tau}\tau$  with  $\tilde{\tau} \rightarrow \tau\tilde{\chi}_1^0$ . See Figs. 9 and 10 for limits on the  $(\mu, M_2)$  plane for  $\tan\beta=1.0$  and different values of  $m_0$ .
- 12 ACCIARRI 01 searches for multi-lepton and/or multi-jet final states from  $\mathcal{R}$  prompt decays with  $L\bar{L}\bar{E}$ ,  $LQ\bar{D}$ , or  $\bar{U}\bar{D}\bar{D}$  couplings at  $\sqrt{s}=189$  GeV. The search is performed for direct and indirect decays of neutralinos, charginos, and scalar leptons, with the  $\tilde{\chi}_1^0$  or a  $\tilde{\ell}$  as LSP and assuming one coupling to be nonzero at a time. Mass limits are derived using simultaneously the constraints from the neutralino, chargino, and slepton analyses; and the  $Z^0$  width measurements from ACCIARRI 00C in a scan of the parameter space assuming MSUGRA with gaugino and scalar mass universality. Updates and supersedes the results from ACCIARRI 99I.
- 13 ABREU 00U searches for the production of charginos and neutralinos in the case of  $R$ -parity violation with  $L\bar{L}\bar{E}$  couplings, using data from  $\sqrt{s}=189$  GeV. They investigate topologies with multiple leptons or jets plus leptons, assuming one coupling to be nonzero at the time and giving rise to direct or indirect decays. Limits are obtained in the  $M_2$  versus  $\mu$  plane and a limit on the neutralino mass is derived from a scan over the parameters  $m_0$  and  $\tan\beta$ .
- 14 ABBIENDI 99F looked for  $\gamma\cancel{E}$  final states at  $\sqrt{s}=183$  GeV. They obtained an upper bound on the cross section for the production  $e^+e^- \rightarrow \tilde{\chi}_2^0\tilde{\chi}_1^0$  followed by the prompt decay  $\tilde{\chi}_2^0 \rightarrow \gamma\tilde{\chi}_1^0$  of 0.075–0.80 pb in the region  $m_{\tilde{\chi}_2^0} + m_{\tilde{\chi}_1^0} > m_Z$ ,  $m_{\tilde{\chi}_2^0}=91$ –183 GeV, and  $\Delta m > 5$  GeV. See Fig. 7 for explicit limits in the  $(m_{\tilde{\chi}_2^0}, m_{\tilde{\chi}_1^0})$  plane.
- 15 ABBIENDI 99F looked for  $\gamma\gamma\cancel{E}$  final states at  $\sqrt{s}=183$  GeV. They obtained an upper bound on the cross section for the production  $e^+e^- \rightarrow \tilde{\chi}_2^0\tilde{\chi}_2^0$  followed by the prompt decay  $\tilde{\chi}_2^0 \rightarrow \gamma\tilde{\chi}_1^0$  of 0.08–0.37 pb for  $m_{\tilde{\chi}_2^0}=45$ –81.5 GeV, and  $\Delta m > 5$  GeV. See Fig. 11 for explicit limits in the  $(m_{\tilde{\chi}_2^0}, m_{\tilde{\chi}_1^0})$  plane.
- 16 ABBOTT 98C searches for trilepton final states ( $\ell=e, \mu$ ). See footnote to ABBOTT 98C in the Chargino Section for details on the assumptions. Assuming a negligible decay rate of  $\tilde{\chi}_1^\pm$  and  $\tilde{\chi}_2^0$  to quarks, they obtain  $m_{\tilde{\chi}_2^0} \gtrsim 103$  GeV.
- 17 ABE 98J searches for trilepton final states ( $\ell=e, \mu$ ). See footnote to ABE 98J in the Chargino Section for details on the assumptions. The quoted result for  $m_{\tilde{\chi}_2^0}$  corresponds to the best limit within the selected range of parameters, obtained for  $m_{\tilde{q}} > m_{\tilde{g}}$ ,  $\tan\beta=2$ , and  $\mu=-600$  GeV.
- 18 ACCIARRI 98F is obtained from direct searches in the  $e^+e^- \rightarrow \tilde{\chi}_{1,2}^0\tilde{\chi}_2^0$  production channels, and indirectly from  $\tilde{\chi}_1^\pm$  and  $\tilde{\chi}_1^0$  searches within the MSSM. See footnote to ACCIARRI 98F in the chargino Section for further details on the assumptions. Data taken at  $\sqrt{s} = 130$ –172 GeV.
- 19 ACCIARRI 98V looked for  $\gamma(\gamma)\cancel{E}$  final states at  $\sqrt{s}=183$  GeV. They obtained an upper bound on the cross section for the production  $e^+e^- \rightarrow \tilde{\chi}_2^0\tilde{\chi}_{1,2}^0$  followed by the prompt decay  $\tilde{\chi}_2^0 \rightarrow \gamma\tilde{\chi}_1^0$ . See Figs. 4a and 6a for explicit limits in the  $(m_{\tilde{\chi}_2^0}, m_{\tilde{\chi}_1^0})$  plane.
- 20 BARATE 98H looked for  $\gamma\gamma\cancel{E}$  final states at  $\sqrt{s} = 161, 172$  GeV. They obtained an upper bound on the cross section for the production  $e^+e^- \rightarrow \tilde{\chi}_2^0\tilde{\chi}_2^0$  followed by the prompt decay  $\tilde{\chi}_2^0 \rightarrow \gamma\tilde{\chi}_1^0$  of 0.4–0.8 pb for  $m_{\tilde{\chi}_2^0} = 10$ –80 GeV. The bound above is for the specific case of  $\tilde{\chi}_1^0 = \tilde{H}^0$  and  $\tilde{\chi}_2^0 = \tilde{\gamma}$  and  $m_{\tilde{e}_R} = 100$  GeV. See Fig. 6 and 7 for explicit limits in the  $(\tilde{\chi}_2^0, \tilde{\chi}_1^0)$  plane and in the  $(\tilde{\chi}_2^0, \tilde{e}_R)$  plane.

- <sup>21</sup> BARATE 98J looked for  $\gamma\gamma \cancel{E}$  final states at  $\sqrt{s} = 161\text{--}183$  GeV. They obtained an upper bound on the cross section for the production  $e^+e^- \rightarrow \tilde{\chi}_2^0\tilde{\chi}_2^0$  followed by the prompt decay  $\tilde{\chi}_2^0 \rightarrow \gamma\tilde{\chi}_1^0$  of  $0.08\text{--}0.24$  pb for  $m_{\tilde{\chi}_2^0} < 91$  GeV. The bound above is for the specific case of  $\tilde{\chi}_1^0 = \tilde{H}^0$  and  $\tilde{\chi}_2^0 = \tilde{\gamma}$  and  $m_{\tilde{e}_R} = 100$  GeV.
- <sup>22</sup> ABACHI 96 searches for 3-lepton final states. Efficiencies are calculated using mass relations and branching ratios in the Minimal Supergravity scenario. Results are presented as lower bounds on  $\sigma(\tilde{\chi}_1^\pm\tilde{\chi}_2^0) \times B(\tilde{\chi}_1^\pm \rightarrow \ell\nu_\ell\tilde{\chi}_1^0) \times B(\tilde{\chi}_2^0 \rightarrow \ell^+\ell^-\tilde{\chi}_1^0)$  as a function of  $m_{\tilde{\chi}_1^0}$ . Limits range from 3.1 pb ( $m_{\tilde{\chi}_1^0} = 45$  GeV) to 0.6 pb ( $m_{\tilde{\chi}_1^0} = 100$  GeV).
- <sup>23</sup> ABE 96K looked for trilepton events from chargino-neutralino production. They obtained lower bounds on  $m_{\tilde{\chi}_2^0}$  as a function of  $\mu$ . The lower bounds are in the 45–50 GeV range for gaugino-dominant  $\tilde{\chi}_2^0$  with negative  $\mu$ , if  $\tan\beta < 10$ . See paper for more details of the assumptions.

## $\tilde{\chi}_1^\pm, \tilde{\chi}_2^\pm$ (Charginos) MASS LIMITS

Charginos are unknown mixtures of  $w$ -inos and charged higgsinos (the supersymmetric partners of  $W$  and Higgs bosons). A lower mass limit for the lightest chargino ( $\tilde{\chi}_1^\pm$ ) of approximately 45 GeV, independent of the field composition and of the decay mode, has been obtained by the LEP experiments from the analysis of the  $Z$  width and decays. These results, as well as other now superseded limits from  $e^+e^-$  collisions at energies below 136 GeV, and from hadronic collisions, can be found in the 1998 Edition (The European Physical Journal **C3** 1 (1998)) of this Review.

Unless otherwise stated, results in this section assume spectra, production rates, decay modes and branching ratios as evaluated in the MSSM, with gaugino and sfermion mass unification at the GUT scale. These papers generally study production of  $\tilde{\chi}_1^0\tilde{\chi}_2^0$ ,  $\tilde{\chi}_1^+\tilde{\chi}_1^-$  and (in the case of hadronic collisions)  $\tilde{\chi}_1^+\tilde{\chi}_2^0$  pairs, including the effects of cascade decays. The mass limits on  $\tilde{\chi}_1^\pm$  are either direct, or follow indirectly from the constraints set by the non-observation of  $\tilde{\chi}_2^0$  states on the gaugino and higgsino MSSM parameters  $M_2$  and  $\mu$ . For generic values of the MSSM parameters, limits from high-energy  $e^+e^-$  collisions coincide with the highest value of the mass allowed by phase-space, namely  $m_{\tilde{\chi}_1^\pm} \lesssim \sqrt{s}/2$ . The still unpublished combination of the results of

the four LEP collaborations from the 2000 run of LEP2 at  $\sqrt{s}$  up to  $\simeq 209$  GeV yields a lower mass limit of 103.5 GeV valid for general MSSM models. The limits become however weaker in certain regions of the MSSM parameter space where the detection efficiencies or production cross sections are suppressed. For example, this may happen when: (i) the mass differences  $\Delta m_{+} = m_{\tilde{\chi}_1^\pm} - m_{\tilde{\chi}_1^0}$  or  $\Delta m_{\nu} = m_{\tilde{\chi}_1^\pm} - m_{\tilde{\nu}}$  are very small, and the detection efficiency is reduced; (ii) the electron sneutrino mass is small, and the  $\tilde{\chi}_1^\pm$  production rate is suppressed due to a destructive interference between  $s$  and  $t$  channel exchange diagrams. The regions of MSSM parameter space where the following limits are valid are indicated in the comment lines or in the footnotes.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
		1 AAD	13 ATLS	$3\ell^\pm + \cancel{E}_T$ , pMSSM, SMS
		2 AAD	13B ATLS	$2\ell^\pm + \cancel{E}_T$ , pMSSM, SMS
>540	95	3 AAD	12CT ATLS	$\geq 4\ell^\pm, \cancel{E}, m_{\tilde{\chi}_1^0} > 300$ GeV

		4	CHATRCHYAN 12BJ	CMS	$\geq 2 \ell, \text{jets} + \cancel{E}_T, pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$	
>101	95	5	ABBIENDI	04H OPAL	all $\tan\beta, \Delta m_+ > 5 \text{ GeV}, m_0 > 500 \text{ GeV}, A_0 = 0$	
> 89		6	ABBIENDI	03H OPAL	$0.5 \leq \Delta m_+ \leq 5 \text{ GeV}$ , higgsino-like, $\tan\beta=1.5$	
> 97.1	95	7	ABDALLAH	03M DLPH	$\tilde{\chi}_1^\pm, \Delta m_+ \geq 3 \text{ GeV}, m_{\tilde{\nu}} > m_{\tilde{\chi}^\pm}$	
> 75	95	7	ABDALLAH	03M DLPH	$\tilde{\chi}_1^\pm$ , higgsino, all $\Delta m_+, m_{\tilde{f}} > m_{\tilde{\chi}^\pm}$	
> 70	95	7	ABDALLAH	03M DLPH	$\tilde{\chi}_1^\pm$ , all $\Delta m_+, m_{\tilde{\nu}} > 500 \text{ GeV}, M_2 \leq 2M_1 \leq 10M_2$	
> <b>94</b>	95	8	ABDALLAH	03M DLPH	$\tilde{\chi}_1^\pm, \tan\beta \leq 40, \Delta m_+ > 3 \text{ GeV}$ , all $m_0$	
> 88	95	9	HEISTER	02J ALEP	$\tilde{\chi}_1^\pm$ , all $\Delta m_+$ , large $m_0$	
> 67.7	95	10	ACCIARRI	00D L3	$\tan\beta > 0.7$ , all $\Delta m_+$ , all $m_0$	
> 69.4	95	11	ACCIARRI	00K L3	$e^+e^- \rightarrow \tilde{\chi}^\pm \tilde{\chi}^\mp$ , all $\Delta m_+$ , heavy scalars	
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●						
		12	AALTONEN	13Q CDF	$\tilde{\chi}_1^\pm \rightarrow \tau X$ , simplified gravity- and gauge-mediated models	
		13	AAD	12AS ATLS	$3\ell^\pm + \cancel{E}_T$ , pMSSM	
		14	AAD	12T ATLS	$\ell^\pm \ell^\mp + \cancel{E}_T, \ell^\pm \ell^\pm + \cancel{E}_T, pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$	
>163	95	15	CHATRCHYAN 11B	CMS	$\tilde{W}^0 \rightarrow \gamma \tilde{G}, \tilde{W}^\pm \rightarrow \ell^\pm \tilde{G}$ , GMSB	
		16	CHATRCHYAN 11V	CMS	$\tan\beta=3, m_0=60 \text{ GeV}, A_0=0, \mu > 0$	
>129	95	17	AALTONEN	09G CDF	$p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$	
>138	95	18	ABAZOV	09T D0	$p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$	
		19	AALTONEN	08AE CDF	$p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$	
		20	AALTONEN	08L CDF	$p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$	
>229	95	21	ABAZOV	08F D0	$p\bar{p} \rightarrow \tilde{\chi}\tilde{\chi}, \tilde{\chi}=\tilde{\chi}_2^0, \tilde{\chi}_1^\pm, \tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ , GMSB	
		22	AALTONEN	07J CDF	$p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$	
		23	ABULENCIA	07H CDF	$\cancel{R}, LL\bar{E}$	
		24	ABULENCIA	07N CDF	$p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$	
		25	ABAZOV	06D D0	$\cancel{R}, LL\bar{E}$	
>195	95	26	ABAZOV	05A D0	$p\bar{p} \rightarrow \tilde{\chi}\tilde{\chi}, \tilde{\chi}=\tilde{\chi}_2^0, \tilde{\chi}_1^\pm, \tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ , GMSB	
>167	95	27	ACOSTA	05E CDF	$p\bar{p} \rightarrow \tilde{\chi}\tilde{\chi}, \tilde{\chi}=\tilde{\chi}_2^0, \tilde{\chi}_1^\pm, \tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ , GMSB	
> 66	95	28,29	ABDALLAH	04H DLPH	AMSB, $\mu > 0$	
>102.5	95	30,31	ABDALLAH	04M DLPH	$\cancel{R}(\overline{UDD})$	
>100		32	ABDALLAH	03D DLPH	$e^+e^- \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_1^\mp (\tilde{\chi}_1^\pm \rightarrow \tilde{\tau}_1 \nu_\tau, \tilde{\tau}_1 \rightarrow \tau \tilde{G})$	
>103		33	HEISTER	03G ALEP	$\cancel{R}$ decays, $m_0 > 500 \text{ GeV}$	
>102.7	95	34	ACHARD	02 L3	$\cancel{R}$ , MSUGRA	
		35	GHODBANE	02 THEO		

> 94.3	95	36 ABREU	01C	DLPH	$\tilde{\chi}^{\pm} \rightarrow \tau J$
> 93.8	95	37 ACCIARRI	01	L3	$\tilde{R}$ , all $m_0$ , $0.7 \leq \tan\beta \leq 40$
>100	95	38 BARATE	01B	ALEP	$\tilde{R}$ decays, $m_0 > 500$ GeV
> 91.8	95	39 ABREU	00V	DLPH	$e^+ e^- \rightarrow \tilde{\chi}_1^{\pm} \tilde{\chi}_1^{\pm} (\tilde{\chi}_1^{\pm} \rightarrow \tilde{\tau}_1 \nu_{\tau}, \tilde{\tau}_1 \rightarrow \tau \tilde{G})$
		40 CHO	00B	THEO	EW analysis
> 76	95	41 ABBIENDI	99T	OPAL	$\tilde{R}$ , $m_0=500$ GeV
> 51	95	42 MALTONI	99B	THEO	EW analysis, $\Delta m_+ \sim 1$ GeV
> 81.5	95	43 ABE	98J	CDF	$p\bar{p} \rightarrow \tilde{\chi}_1^{\pm} \tilde{\chi}_2^0$
		44 ACKERSTAFF	98K	OPAL	$\tilde{\chi}^+ \rightarrow \ell^+ \cancel{E}$
> 65.7	95	45 ACKERSTAFF	98L	OPAL	$\Delta m_+ > 3$ GeV, $\Delta m_{\nu} > 2$ GeV
		46 ACKERSTAFF	98V	OPAL	light gluino
		47 CARENA	97	THEO	$g_{\mu} - 2$
		48 KALINOWSKI	97	THEO	$W \rightarrow \tilde{\chi}_1^{\pm} \tilde{\chi}_1^0$
		49 ABE	96K	CDF	$p\bar{p} \rightarrow \tilde{\chi}_1^{\pm} \tilde{\chi}_2^0$

<sup>1</sup> AAD 13 searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for charginos and neutralinos decaying to a final state with three leptons ( $e$  and  $\mu$ ) and missing transverse energy. No excess beyond the Standard Model expectation is observed. Exclusion limits are derived in the phenomenological MSSM, see Fig. 2 and 3, and in simplified models, see Fig. 4. For the simplified models with intermediate slepton decays, degenerate  $\tilde{\chi}_1^{\pm}$  and  $\tilde{\chi}_2^0$  masses up to 500 GeV are excluded at 95% C.L. for very large mass differences with the  $\tilde{\chi}_1^0$ . Supersedes AAD 12AS.

<sup>2</sup> AAD 13B searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for gauginos decaying to a final state with two leptons ( $e$  and  $\mu$ ) and missing transverse energy. No excess beyond the Standard Model expectation is observed. Limits are derived in a simplified model of wino-like chargino pair production, where the chargino always decays to the lightest neutralino via an intermediate on-shell charged slepton, see Fig. 2(b). Chargino masses between 110 and 340 GeV are excluded at 95% C.L. for  $m_{\tilde{\chi}_1^0} = 10$  GeV. Exclusion limits are also derived in the phenomenological MSSM, see Fig. 3.

<sup>3</sup> AAD 12CT searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for events containing four or more leptons (electrons or muons) and either moderate values of missing transverse momentum or large effective mass. No significant excess is found in the data. Limits are presented in a simplified model of  $R$ -parity violating supersymmetry in which charginos are pair-produced and then decay into a  $W$ -boson and a  $\tilde{\chi}_1^0$ , which in turn decays through an RPV coupling into two charged leptons ( $e^{\pm} e^{\mp}$  or  $e^{\pm} \mu^{\mp}$ ) and a neutrino. In this model, chargino masses up to 540 GeV are excluded at 95% C.L. for  $m_{\tilde{\chi}_1^0}$  above 300 GeV, see Fig. 3a. The limit deteriorates for lighter  $\tilde{\chi}_1^0$ . Limits are also set in an  $R$ -parity violating mSUGRA model, see Fig. 3b.

<sup>4</sup> CHATRCHYAN 12BJ searched in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for direct electroweak production of charginos and neutralinos in events with at least two leptons, jets and missing transverse momentum. No significant excesses over the expected SM backgrounds are observed and 95% C.L. limits on the production cross section of  $\tilde{\chi}_1^{\pm} \tilde{\chi}_2^0$  pair production were set in a number of simplified models, see Figs. 7 to 12.

<sup>5</sup> ABBIENDI 04H search for charginos and neutralinos in events with acoplanar leptons+jets and multi-jet final states in the 192–209 GeV data, combined with the results on leptonic final states from ABBIENDI 04. The results hold for a scan over the parameter space covering the region  $0 < M_2 < 5000$  GeV,  $-1000 < \mu < 1000$  GeV and  $\tan\beta$  from 1 to 40. This limit supersedes ABBIENDI 00H.

<sup>6</sup> ABBIENDI 03H used  $e^+ e^-$  data at  $\sqrt{s} = 188$ –209 GeV to search for chargino pair production in the case of small  $\Delta m_+$ . They select events with an energetic photon, large

$\cancel{E}$  and little hadronic or leptonic activity. The bound applies to higgsino-like charginos with zero lifetime and a 100% branching ratio  $\tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 W^*$ . The mass limit for gaugino-like charginos, in case of non-universal gaugino masses, is of 92 GeV for  $m_{\tilde{\nu}} = 1000$  GeV and is lowered to 74 GeV for  $m_{\tilde{\nu}} \geq 100$  GeV. Limits in the plane  $(m_{\tilde{\chi}_1^\pm}, \Delta m_+)$  are shown in Fig. 7. Exclusion regions are also derived for the AMSB scenario in the  $(m_{3/2}, \tan\beta)$  plane, see their Fig. 9.

- <sup>7</sup> ABDALLAH 03M searches for the production of charginos using data from  $\sqrt{s} = 192$  to 208 GeV to investigate topologies with multiple leptons, jets plus leptons, multi-jets, or isolated photons. The first limit holds for  $\tan\beta \geq 1$  and is obtained at  $\Delta m_+ = 3$  GeV in the higgsino region. For  $\Delta m_+ \geq 10$  (5) GeV and large  $m_0$ , the limit improves to 102.7 (101.7) GeV. For the region of small  $\Delta m_+$ , all data from  $\sqrt{s} = 130$  to 208 GeV are used to investigate final states with heavy stable charged particles, decay vertices inside the detector and soft topologies with a photon from initial state radiation. The second limit is obtained in the higgsino region, assuming gaugino mass universality at the GUT scale and  $1 < \tan\beta < 50$ . For the case of non-universality of gaugino masses, the parameter space is scanned in the domain  $1 < \tan\beta < 50$  and, for  $\Delta m_+ < 3$  GeV, for values of  $M_1$ ,  $M_2$  and  $\mu$  such that  $M_2 \leq 2M_1 \leq 10M_2$  and  $|\mu| \geq M_2$ . The third limit is obtained in the gaugino region. See Fig. 36 for the dependence of the low  $\Delta m_+$  limits on  $\Delta m_+$ . These limits include and update the results of ABREU 00J and ABREU 00T.
- <sup>8</sup> ABDALLAH 03M uses data from  $\sqrt{s} = 192$ –208 GeV to obtain limits in the framework of the MSSM with gaugino and sfermion mass universality at the GUT scale. An indirect limit on the mass of charginos is derived by constraining the MSSM parameter space by the results from direct searches for neutralinos (including cascade decays), for charginos and for sleptons. These limits are valid for values of  $M_2 < 1$  TeV,  $|\mu| \leq 2$  TeV with the  $\tilde{\chi}_1^0$  as LSP. Constraints from the Higgs search in the  $m_h^{\max}$  scenario assuming  $m_t = 174.3$  GeV are included. The quoted limit applies if there is no mixing in the third family or when  $m_{\tilde{\tau}_1} - m_{\tilde{\chi}_1^0} > 6$  GeV. If mixing is included the limit degrades to 90 GeV. See Fig. 43 for the mass limits as a function of  $\tan\beta$ . These limits update the results of ABREU 00w.
- <sup>9</sup> HEISTER 02J search for chargino production with small  $\Delta m_+$  in final states with a hard isolated initial state radiation photon and few low-momentum particles, using 189–208 GeV data. This search is sensitive in the intermediate  $\Delta m_+$  region. Combined with searches for  $\cancel{E}$  topologies and for stable charged particles, the above bound is obtained for  $m_0$  larger than few hundred GeV,  $1 < \tan\beta < 300$  and holds for any chargino field contents. For light scalars, the general limit reduces to the one from the  $Z^0$ , but under the assumption of gaugino and sfermion mass unification the above bound is recovered. See Figs. 4–6 for the more general dependence of the limits on  $\Delta m_+$ . Updates BARATE 98X.
- <sup>10</sup> ACCIARRI 00D data collected at  $\sqrt{s}=189$  GeV. The results hold over the full parameter space defined by  $0.7 \leq \tan\beta \leq 60$ ,  $0 \leq M_2 \leq 2$  TeV,  $|\mu| \leq 2$  TeV  $m_0 \leq 500$  GeV. The results of slepton searches from ACCIARRI 99W are used to help set constraints in the region of small  $m_0$ . See their Figs. 5 for the  $\tan\beta$  and  $M_2$  dependence on the limits. See the text for the impact of a large  $B(\tilde{\chi}^\pm \rightarrow \tau \tilde{\nu}_\tau)$  on the result. The region of small  $\Delta m_+$  is excluded by the analysis of ACCIARRI 00K. Updates ACCIARRI 98F.
- <sup>11</sup> ACCIARRI 00K searches for the production of charginos with small  $\Delta m_+$  using data from  $\sqrt{s}=189$  GeV. They investigate soft final states with a photon from initial state radiation. The results are combined with the limits on prompt decays from ACCIARRI 00D and from heavy stable charged particles from ACCIARRI 99L (see Heavy Charged Lepton Searches). The production and decay branching ratios are evaluated within the MSSM, assuming heavy sfermions. The parameter space is scanned in the domain  $1 < \tan\beta < 50$ ,  $0.3 < M_1/M_2 < 50$ , and  $0 < |\mu| < 2$  TeV. The limit is obtained in the higgsino region and improves to 78.6 GeV for gaugino-like charginos. The limit is unchanged for light scalar quarks. For light  $\tilde{\tau}$  or  $\tilde{\nu}_\tau$ , the limit is unchanged in the gaugino-like region and is lowered by 0.8 GeV in the higgsino-like case. For light  $\tilde{\mu}$  or  $\tilde{\nu}_\mu$ , the limit is unchanged in

- the higgsino-like region and is lowered by 0.9 GeV in the gaugino-like region. No direct mass limits are obtained for light  $\tilde{e}$  or  $\tilde{\nu}_e$ .
- 12 AALTONEN 13Q searched in  $6.0 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for evidence of chargino-neutralino associated production in like-sign dilepton final states. One lepton is identified as the hadronic decay of a tau lepton, while the other is an electron or muon. Good agreement with the Standard Model predictions is observed and limits are set on the chargino-neutralino cross section for simplified gravity- and gauge-mediated models, see their Figs 2 and 3.
  - 13 AAD 12AS searched in  $2.06 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for charginos and neutralinos decaying to a final state with three leptons ( $e$  and  $\mu$ ) and missing transverse energy. No excess beyond the Standard Model expectation is observed. Exclusion limits are derived in the phenomenological MSSM, see Fig. 2 (top), and in simplified models, see Fig. 2 (bottom).
  - 14 AAD 12T looked in  $1 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for the production of supersymmetric particles decaying into final states with missing transverse momentum and exactly two isolated leptons ( $e$  or  $\mu$ ). Opposite-sign and same-sign dilepton events were separately studied. Additionally, in opposite-sign events, a search was made for an excess of same-flavor over different-flavor lepton pairs. No excess over the expected background is observed and limits are placed on the effective production cross section of opposite-sign dilepton events with  $\cancel{E}_T > 250 \text{ GeV}$  and on same-sign dilepton events with  $\cancel{E}_T > 100 \text{ GeV}$ . The latter limit is interpreted in a simplified electroweak gaugino production model as a lower chargino mass limit.
  - 15 CHATRCHYAN 11B looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s}=7 \text{ TeV}$  for events with an isolated lepton ( $e$  or  $\mu$ ), a photon and  $\cancel{E}_T$  which may arise in a generalized gauge mediated model from the decay of Wino-like NLSPs. No evidence for an excess over the expected background is observed. Limits are derived in the plane of squark/gluino mass versus Wino mass (see Fig. 4). Mass degeneracy of the produced squarks and gluinos is assumed.
  - 16 CHATRCHYAN 11V looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 3$  isolated leptons ( $e$ ,  $\mu$  or  $\tau$ ), with or without jets and  $\cancel{E}_T$ . No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM  $(m_0, m_{1/2})$  plane for  $\tan\beta = 3$  (see Fig. 5).
  - 17 AALTONEN 09G searched in  $976 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with trileptons ( $\mu\mu\mu$  or  $\mu\mu e$ ) with a low, 5 GeV,  $p_T$  threshold, and large  $\cancel{E}_T$  from the decay of  $\tilde{\chi}_1^\pm \tilde{\chi}_2^0 X$ . The selected number of events is consistent with the SM background expectation. The results are combined with the analysis of AALTONEN 07J to set a limit on the  $\tilde{\chi}_1^\pm$  mass for a mSUGRA scenario with no slepton mixing.
  - 18 ABAZOV 09T searched in  $2.3 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with trileptons ( $e$ ,  $\mu$  or hadronically decaying  $\tau$ ) from the decay of  $\tilde{\chi}_1^\pm \tilde{\chi}_2^0 X$  and large  $\cancel{E}_T$ . No evidence for a signal is observed. The data are used to constrain the cross section times branching ratio as a function of the  $\tilde{\chi}_1^\pm$  mass under the assumption that  $m_{\tilde{\chi}_1^\pm} = m_{\tilde{\chi}_2^0} = 2 m_{\tilde{\chi}_1^0}$ ,  $\tan\beta = 3$ ,  $\mu > 0$  and that the sleptons are heavier than the  $\tilde{\chi}_1^\pm$ , see their Fig. 8. A chargino lighter than 138 GeV is excluded in the “3l-max” scenario. Exclusion regions in the  $(m_0, m_{1/2})$  plane are shown in their Fig. 9 for a mSUGRA scenario with  $\tan\beta = 3$ ,  $A_0 = 0$  and  $\mu > 0$ . The  $\tan\beta$  dependence of this exclusion is illustrated in Fig. 10. Supersedes the results of ABAZOV 05U.
  - 19 AALTONEN 08AE searched in  $2.0 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with trileptons ( $e$ ,  $\mu$  or a charged isolated track from  $\tau$ ) from the decay of  $p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0 X$  and large  $\cancel{E}_T$ . The selected number of events is consistent with the SM background expectation. The data are used to constrain the cross section times branching ratio as a function of the  $\tilde{\chi}_1^\pm$  mass. Exclusion regions in the  $(m_0, m_{1/2})$  plane are shown in their Fig. 2 for a mSUGRA scenario. When the  $\tilde{\chi}_1^\pm$  is nearly mass degenerate with the  $\tilde{\tau}_1$  the

leptons are too soft and no limit is obtained. For the case  $m_0 = 60$  GeV a lower limit of 145 GeV on the chargino mass is obtained in this mSUGRA scenario.

- 20 AALTONEN 08L searched in  $0.7$  to  $1.0 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV for events with one high- $p_T$  electron or muon and two additional leptons ( $e$  or  $\mu$ ) from the decay of  $\tilde{\chi}_1^\pm \tilde{\chi}_2^0 X$ . The selected number of events is consistent with the SM background expectation. The data are used to constrain the cross section times branching ratio as a function of the  $\tilde{\chi}_1^\pm$  mass. The results are compared to three MSSM scenarios. An exclusion on chargino and neutralino production is only obtained in a scenario of no mixing between sleptons, yielding nearly equal branching ratios to all three lepton flavors. It amounts to  $m_{\tilde{\chi}_1^\pm} > 151$  GeV, while the analysis is not sensitive to chargino masses below about 110 GeV. The analyses have been combined with the analyses of AALTONEN 07J and ABULENCIA 07N. The observed limits for the combination are less stringent than the one obtained for the high- $p_T$  analysis due to slight excesses in the other channels.
- 21 ABAZOV 08F looked in  $1.1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV for diphoton events with large  $E_T$ . They may originate from the production of  $\tilde{\chi}^\pm$  in pairs or associated to a  $\tilde{\chi}_2^0$ , decaying to a  $\tilde{\chi}_1^0$  which itself decays promptly in GMSB to  $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ . No significant excess was found compared to the background expectation. A limit is derived on the masses of SUSY particles in the GMSB framework for  $M = 2\Lambda$ ,  $N = 1$ ,  $\tan\beta = 15$  and  $\mu > 0$ , see Figure 2. It also excludes  $\Lambda < 91.5$  TeV. Supersedes the results of ABAZOV 05A.
- 22 AALTONEN 07J searched in  $0.7$  to  $1.1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV for events with either two same sign leptons ( $e$  or  $\mu$ ) or trileptons from the decay of  $\tilde{\chi}_1^\pm \tilde{\chi}_2^0 X$  and large  $E_T$ . The selected number of events is consistent with the SM background expectation. The data are used to constrain the cross section times branching ratio as a function of the  $\tilde{\chi}_1^\pm$  mass. The results, shown in their Fig. 2, are compared to several MSSM scenarios. The strongest exclusion is in the case of no mixing between sleptons, yielding nearly equal branching ratios to all three lepton flavors, and amounting to  $m_{\tilde{\chi}_1^\pm} > 129$  GeV. This analysis includes the same sign dilepton analysis of ABULENCIA 07N.
- 23 ABULENCIA 07H searched in  $346 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV for events with at least three leptons ( $e$  or  $\mu$ ) from the decay of  $\tilde{\chi}_1^0$  via  $LL\bar{E}$  couplings. The results are consistent with the hypothesis of no signal. Upper limits on the cross-section are extracted and a limit is derived in the framework of mSUGRA on the masses of  $\tilde{\chi}_1^0$  and  $\tilde{\chi}_1^\pm$ , see e.g. their Fig. 3 and Tab. II.
- 24 ABULENCIA 07N searched in  $1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV for events with two same sign leptons ( $e$  or  $\mu$ ) from the decay of  $\tilde{\chi}_1^\pm \tilde{\chi}_2^0 X$  and large  $E_T$ . A slight excess of 13 events is observed over a SM background expectation of  $7.8 \pm 1.1$ . However, the kinematic distributions do not show any anomalous deviation from expectations in any particular region of parameter space.
- 25 ABAZOV 06D looked in  $360 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV for events with three leptons originating from the pair production of charginos and neutralinos, followed by  $R$  decays mediated by  $LL\bar{E}$  couplings. One coupling is assumed to be dominant at a time. No significant excess was found compared to the background expectation in the  $eel$ ,  $\mu\mu\ell$  nor  $ee\tau$  ( $\ell = e, \mu$ ) final states. Upper limits on the cross-section are extracted in a specific MSUGRA model and a MSSM model without unification of  $M_1$  and  $M_2$  at the GUT scale. A limit is derived on the masses of charginos and neutralinos for both scenarios assuming  $\lambda_{ijk}$  couplings such that the decay length is less than 1 cm, see their Table III and Fig. 4.
- 26 ABAZOV 05A looked in  $263 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV for diphoton events with large  $E_T$ . They may originate from the production of  $\tilde{\chi}^\pm$  in pairs or associated to a  $\tilde{\chi}_2^0$ , decaying to a  $\tilde{\chi}_1^0$  which itself decays promptly in GMSB to  $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ . No

- significant excess was found at large  $\cancel{E}_T$  compared to the background expectation. A limit is derived on the masses of SUSY particles in the GMSB framework for  $M = 2 \Lambda$ ,  $N = 1$ ,  $\tan\beta = 15$  and  $\mu > 0$ , see Figure 2. It also excludes  $\Lambda < 79.6$  TeV. Very similar results are obtained for different choices of parameters, see their Table 2. Supersedes the results of ABBOTT 98.
- 27 ACOSTA 05E looked in  $202 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.96$  TeV for diphoton events with large  $\cancel{E}_T$ . They may originate from the production of  $\tilde{\chi}^\pm$  in pairs or associated to a  $\tilde{\chi}_2^0$ , decaying to a  $\tilde{\chi}_1^0$  which itself decays promptly in GMSB to  $\gamma\tilde{G}$ . No events are selected at large  $\cancel{E}_T$  compared to the background expectation. A limit is derived on the masses of SUSY particles in the GMSB framework for  $M = 2 \Lambda$ ,  $N = 1$ ,  $\tan\beta = 15$  and  $\mu > 0$ , see Figure 2. It also excludes  $\Lambda < 69$  TeV. Supersedes the results of ABE 99I.
- 28 ABDALLAH 04H use data from LEP 1 and  $\sqrt{s} = 192\text{--}208$  GeV. They re-use results or re-analyze the data from ABDALLAH 03M to put limits on the parameter space of anomaly-mediated supersymmetry breaking (AMSB), which is scanned in the region  $1 < m_{3/2} < 50$  TeV,  $0 < m_0 < 1000$  GeV,  $1.5 < \tan\beta < 35$ , both signs of  $\mu$ . The constraints are obtained from the searches for mass degenerate chargino and neutralino, for SM-like and invisible Higgs, for leptonically decaying charginos and from the limit on non-SM  $Z$  width of 3.2 MeV. The limit is for  $m_t = 174.3$  GeV (see Table 2 for other  $m_t$  values).
- 29 The limit improves to 73 GeV for  $\mu < 0$ .
- 30 ABDALLAH 04M use data from  $\sqrt{s} = 192\text{--}208$  GeV to derive limits on sparticle masses under the assumption of  $\cancel{R}$  with  $LL\bar{E}$  or  $\overline{UDD}$  couplings. The results are valid in the ranges  $90 < m_0 < 500$  GeV,  $0.7 < \tan\beta < 30$ ,  $-200 < \mu < 200$  GeV,  $0 < M_2 < 400$  GeV. Supersedes the result of ABREU 01D and ABREU 00U.
- 31 The limit improves to 103 GeV for  $LL\bar{E}$  couplings.
- 32 ABDALLAH 03D use data from  $\sqrt{s} = 183\text{--}208$  GeV. They look for final states with two acoplanar leptons, expected in GMSB when the  $\tilde{\tau}_1$  is the NLSP and assuming a short-lived  $\tilde{\chi}_1^\pm$ . Limits are obtained in the plane  $(m(\tilde{\tau}), m(\tilde{\chi}_1^\pm))$  for different domains of  $m(\tilde{G})$ , after combining these results with the search for slepton pair production from the same paper. The limit above is valid if the  $\tilde{\tau}_1$  is the NLSP for all values of  $m(\tilde{G})$  provided  $m(\tilde{\chi}_1^\pm) - m(\tilde{\tau}_1) \geq 0.3$  GeV. For larger  $m(\tilde{G}) > 100$  eV the limit improves to 102 GeV, see their Fig. 11. In the co-NLSP scenario, the limits are 96 and 102 GeV for all  $m(\tilde{G})$  and  $m(\tilde{G}) > 100$  eV, respectively. Supersedes the results of ABREU 01G.
- 33 HEISTER 03G searches for the production of charginos prompt decays. in the case of  $\cancel{R}$  prompt decays with  $LL\bar{E}$ ,  $LQ\bar{D}$  or  $\overline{UDD}$  couplings at  $\sqrt{s}=189\text{--}209$  GeV. The search is performed for indirect decays, assuming one coupling at a time to be non-zero. The limit holds for  $\tan\beta=1.41$ . Excluded regions in the  $(\mu, M_2)$  plane are shown in their Fig. 3.
- 34 ACHARD 02 searches for the production of sparticles in the case of  $\cancel{R}$  prompt decays with  $LL\bar{E}$  or  $\overline{UDD}$  couplings at  $\sqrt{s}=189\text{--}208$  GeV. The search is performed for direct and indirect decays, assuming one coupling at the time to be nonzero. The MSUGRA limit results from a scan over the MSSM parameter space with the assumption of gaugino and scalar mass unification at the GUT scale, imposing simultaneously the exclusions from neutralino, chargino, sleptons, and squarks analyses. The limit of  $\tilde{\chi}_1^\pm$  holds for  $\overline{UDD}$  couplings and increases to 103.0 GeV for  $LL\bar{E}$  couplings. For L3 limits from  $LQ\bar{D}$  couplings, see ACCIARRI 01.
- 35 GHODBANE 02 reanalyzes DELPHI data at  $\sqrt{s}=189$  GeV in the presence of complex phases for the MSSM parameters.
- 36 ABREU 01C looked for  $\tau$  pairs with  $\cancel{R}$  at  $\sqrt{s}=183\text{--}189$  GeV to search for the associated production of charginos, followed by the decay  $\tilde{\chi}^\pm \rightarrow \tau J$ ,  $J$  being an invisible massless particle. See Fig. 6 for the regions excluded in the  $(\mu, M_2)$  plane.
- 37 ACCIARRI 01 searches for multi-lepton and/or multi-jet final states from  $\cancel{R}$  prompt decays with  $LL\bar{E}$ ,  $LQ\bar{D}$ , or  $\overline{UDD}$  couplings at  $\sqrt{s}=189$  GeV. The search is performed for direct and indirect decays of neutralinos, charginos, and scalar leptons, with the  $\tilde{\chi}_1^0$  or a  $\tilde{\ell}$  as LSP and assuming one coupling to be nonzero at a time. Mass limits are derived

- using simultaneously the constraints from the neutralino, chargino, and slepton analyses; and the  $Z^0$  width measurements from ACCIARRI 00C in a scan of the parameter space assuming MSUGRA with gaugino and scalar mass universality. Updates and supersedes the results from ACCIARRI 99I.
- 38 BARATE 01B searches for the production of charginos in the case of  $\bar{R}$  prompt decays with  $LL\bar{E}$ ,  $LQ\bar{D}$ , or  $UDD$  couplings at  $\sqrt{s}=189\text{--}202$  GeV. The search is performed for indirect decays, assuming one coupling at a time to be nonzero. Updates BARATE 00H.
- 39 ABREU 00V use data from  $\sqrt{s}=183\text{--}189$  GeV. They look for final states with two acoplanar leptons, expected in GMSB when the  $\tilde{\tau}_1$  is the NLSP and assuming a short-lived  $\tilde{\chi}_1^\pm$ . Limits are obtained in the plane  $(m_{\tilde{\tau}}, m_{\tilde{\chi}_1^\pm})$  for different domains of  $m_{\tilde{G}}$ , after combining these results with the search for slepton pair production in the SUGRA framework from ABREU 01 to cover prompt decays and on stable particle searches from ABREU 00Q. The limit above is valid for all values of  $m_{\tilde{G}}$ .
- 40 CHO 00B studied constraints on the MSSM spectrum from precision EW observables. Global fits favour charginos with masses at the lower bounds allowed by direct searches. Allowing for variations of the squark and slepton masses does not improve the fits.
- 41 ABBIENDI 99T searches for the production of neutralinos in the case of  $R$ -parity violation with  $LL\bar{E}$ ,  $LQ\bar{D}$ , or  $UDD$  couplings using data from  $\sqrt{s}=183$  GeV. They investigate topologies with multiple leptons, jets plus leptons, or multiple jets, assuming one coupling at the time to be non-zero and giving rise to direct or indirect decays. Mixed decays (where one particle has a direct, the other an indirect decay) are also considered for the  $UDD$  couplings. Upper limits on the cross section are derived which, combined with the constraint from the  $Z^0$  width, allow to exclude regions in the  $M_2$  versus  $\mu$  plane for any coupling. Limits on the chargino mass are obtained for non-zero  $LL\bar{E}$  couplings  $> 10^{-5}$  and assuming decays via a  $W^*$ .
- 42 MALTONI 99B studied the effect of light chargino-neutralino to the electroweak precision data with a particular focus on the case where they are nearly degenerate ( $\Delta m_+ \sim 1$  GeV) which is difficult to exclude from direct collider searches. The quoted limit is for higgsino-like case while the bound improves to 56 GeV for wino-like case. The values of the limits presented here are obtained in an update to MALTONI 99B, as described in MALTONI 00.
- 43 ABE 98J searches for trilepton final states ( $\ell=e,\mu$ ). Efficiencies are calculated using mass relations in the Minimal Supergravity scenario, exploring the domain of parameter space defined by  $1.1 < \tan\beta < 8$ ,  $-1000 < \mu(\text{GeV}) < -200$ , and  $m_{\tilde{q}}/m_{\tilde{g}}=1\text{--}2$ . In this region  $m_{\tilde{\chi}_1^\pm} \sim m_{\tilde{\chi}_2^0}$  and  $m_{\tilde{\chi}_1^\pm} \sim 2m_{\tilde{\chi}_1^0}$ . Results are presented in Fig. 1 as upper bounds on  $\sigma(p\bar{p} \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0) \times B(3\ell)$ . Limits range from 0.8 pb ( $m_{\tilde{\chi}_1^\pm}=50$  GeV) to 0.23 pb ( $m_{\tilde{\chi}_1^\pm}=100$  GeV) at 95%CL. The gaugino mass unification hypothesis and the assumed mass relation between squarks and gluinos define the value of the leptonic branching ratios. The quoted result corresponds to the best limit within the selected range of parameters, obtained for  $m_{\tilde{q}} > m_{\tilde{g}}$ ,  $\tan\beta=2$ , and  $\mu=-600$  GeV. Mass limits for different values of  $\tan\beta$  and  $\mu$  are given in Fig. 2.
- 44 ACKERSTAFF 98K looked for dilepton+ $\cancel{E}_T$  final states at  $\sqrt{s}=130\text{--}172$  GeV. Limits on  $\sigma(e^+e^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^-) \times B^2(\ell)$ , with  $B(\ell)=B(\chi^+ \rightarrow \ell^+ \nu_\ell \chi_1^0)$  ( $B(\ell)=B(\chi^+ \rightarrow \ell^+ \tilde{\nu}_\ell)$ ), are given in Fig. 16 (Fig. 17).
- 45 ACKERSTAFF 98L limit is obtained for  $0 < M_2 < 1500$ ,  $|\mu| < 500$  and  $\tan\beta > 1$ , but remains valid outside this domain. The dependence on the trilinear-coupling parameter  $A$  is studied, and found negligible. The limit holds for the smallest value of  $m_0$  consistent with scalar lepton constraints (ACKERSTAFF 97H) and for all values of  $m_0$  where the condition  $\Delta m_{\tilde{\nu}} > 2.0$  GeV is satisfied.  $\Delta m_{\tilde{\nu}} > 10$  GeV if  $\tilde{\chi}^\pm \rightarrow \ell \tilde{\nu}_\ell$ . The limit improves to 84.5 GeV for  $m_0=1$  TeV. Data taken at  $\sqrt{s}=130\text{--}172$  GeV.

- 46 ACKERSTAFF 98V excludes the light gluino with universal gaugino mass where charginos, neutralinos decay as  $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0 \rightarrow q\bar{q}\tilde{g}$  from total hadronic cross sections at  $\sqrt{s}=130\text{--}172$  GeV. See paper for the case of nonuniversal gaugino mass.
- 47 CARENA 97 studied the constraints on chargino and sneutrino masses from muon  $g-2$ . The bound can be important for large  $\tan\beta$ .
- 48 KALINOWSKI 97 studies the constraints on the chargino-neutralino parameter space from limits on  $\Gamma(W \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_1^0)$  achievable at LEP2. This is relevant when  $\tilde{\chi}_1^\pm$  is "invisible," i.e., if  $\tilde{\chi}_1^\pm$  dominantly decays into  $\tilde{\nu}_\ell \ell^\pm$  with little energy for the lepton. Small otherwise allowed regions could be excluded.
- 49 ABE 96K looked for trilepton events from chargino-neutralino production. The bound on  $m_{\tilde{\chi}_1^\pm}$  can reach up to 47 GeV for specific choices of parameters. The limits on the combined production cross section times 3-lepton branching ratios range between 1.4 and 0.4 pb, for  $45 < m_{\tilde{\chi}_1^\pm} \text{ (GeV)} < 100$ . See the paper for more details on the parameter dependence of the results.

### Long-lived $\tilde{\chi}^\pm$ (Chargino) MASS LIMITS

Limits on charginos which leave the detector before decaying.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
>103	95	1 AAD	13H ATLS	long-lived $\tilde{\chi}^\pm \rightarrow \tilde{\chi}_1^0 \pi^\pm$ , mAMSB, $\Delta m_{\tilde{\chi}_1^0} = 160$ MeV
> 92	95	2 AAD	12BJ ATLS	long-lived $\tilde{\chi}^\pm \rightarrow \pi^\pm \tilde{\chi}_1^0$ , mAMSB
>171	95	3 ABAZOV	09M D0	$\tilde{H}$
>102	95	4 ABBIENDI	03L OPAL	$m_{\tilde{\nu}} > 500$ GeV
none 2–93.0	95	5 ABREU	00T DLPH	$\tilde{H}^\pm$ or $m_{\tilde{\nu}} > m_{\tilde{\chi}^\pm}$
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
>270	95	6 AAD	13BD ATLS	disappearing-track signature, AMSB
>278	95	7 ABAZOV	13B D0	long-lived $\tilde{\chi}^\pm$ , gaugino-like
>244	95	7 ABAZOV	13B D0	long-lived $\tilde{\chi}^\pm$ , higgsino-like
	95	8 ABAZOV	12L D0	long-lived $\tilde{\chi}^\pm$ , gaugino-like
	95	9 ABAZOV	12L D0	long-lived $\tilde{\chi}^\pm$ , higgsino-like
> 83	95	10 BARATE	97K ALEP	
> 28.2	95	ADACHI	90C TOPZ	

<sup>1</sup> AAD 13H searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for direct electroweak production of long-lived charginos in the context of AMSB scenarios. The search is based on the signature of a high-momentum isolated track with few associated hits in the outer part of the tracking system, arising from a chargino decay into a neutralino and a low-momentum pion. The  $p_T$  spectrum of the tracks was found to be consistent with the SM expectations. Constraints on the lifetime and the production cross section were obtained, see Fig. 6. In the minimal AMSB framework with  $\tan\beta = 5$ , and  $\mu > 0$ , a chargino having a mass below 103 (85) GeV for a chargino-neutralino mass splitting  $\Delta m_{\tilde{\chi}_1^0}$  of 160 (170) MeV is excluded at the 95% C.L. See Fig. 7 for more precise bounds.

<sup>2</sup> AAD 12BJ looked in  $1.02 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for signatures of decaying charginos resulting in isolated tracks with few associated hits in the outer region of the tracking system. The  $p_T$  spectrum of the tracks was found to be consistent with the SM expectations. Constraints on the lifetime and the production cross section were obtained. In the minimal AMSB framework with  $m_{3/2} < 32$  TeV,  $m_0 < 1.5$  TeV,  $\tan\beta = 5$ , and  $\mu > 0$ , a chargino having a mass below 92 GeV and a lifetime between 0.5 ns and 2 ns is excluded at the 95% C.L. See their Fig. 8 for more precise bounds.

- <sup>3</sup> ABAZOV 09M searched in  $1.1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with direct production of a pair of charged massive stable particles identified by their TOF. The number of the observed events is consistent with the predicted background. The data are used to constrain the production cross section as a function of the  $\tilde{\chi}_1^\pm$  mass, see their Fig. 2. The quoted limit improves to 206 GeV for gaugino-like charginos.
- <sup>4</sup> ABBIENDI 03L used  $e^+e^-$  data at  $\sqrt{s} = 130\text{--}209 \text{ GeV}$  to select events with two high momentum tracks with anomalous  $dE/dx$ . The excluded cross section is compared to the theoretical expectation as a function of the heavy particle mass in their Fig. 3. The bounds are valid for colorless fermions with lifetime longer than  $10^{-6} \text{ s}$ . Supersedes the results from ACKERSTAFF 98P.
- <sup>5</sup> ABREU 00T searches for the production of heavy stable charged particles, identified by their ionization or Cherenkov radiation, using data from  $\sqrt{s} = 130$  to  $189 \text{ GeV}$ . These limits include and update the results of ABREU 98P.
- <sup>6</sup> AAD 13BD searched in  $20.3 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  for events containing tracks with no associated hits in the outer region of the tracking system resulting from the decay of charginos that are nearly mass degenerate with the lightest neutralino, as is often the case in AMSB scenarios. No significant excess above the background expectation is observed for candidate tracks with large transverse momentum. Constraints on chargino properties are obtained and in the minimal AMSB model, a chargino mass below  $270 \text{ GeV}$  is excluded at 95% C.L., see their Fig. 7.
- <sup>7</sup> ABAZOV 13B looked in  $6.3 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for charged massive long-lived particles in events with muon-like particles that have both speed and ionization energy loss inconsistent with muons produced in beam collisions. In the absence of an excess, limits are set at 95% C.L. on gaugino- and higgsino-like charginos, see their Table 20 and Fig. 23.
- <sup>8</sup> ABAZOV 12L looked in  $5.2 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for charged massive long-lived particles in events in which one or more particles are reconstructed as muons but have speed and ionization energy loss inconsistent with muons produced in beam collisions. Long-lived pair-produced gaugino-like charginos are excluded below  $267 \text{ GeV}$  at 95% C.L. using the nominal value of the NLO production cross section.
- <sup>9</sup> ABAZOV 12L looked in  $5.2 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for charged massive long-lived particles in events in which one or more particles are reconstructed as muons but have speed and ionization energy loss inconsistent with muons produced in beam collisions. Long-lived pair-produced Higgsino-like charginos are excluded below  $217 \text{ GeV}$  at 95% C.L. using the nominal value of the NLO production cross section.
- <sup>10</sup> BARATE 97K uses  $e^+e^-$  data collected at  $\sqrt{s} = 130\text{--}172 \text{ GeV}$ . Limit valid for  $\tan\beta = \sqrt{2}$  and  $m_{\tilde{\nu}} > 100 \text{ GeV}$ . The limit improves to  $86 \text{ GeV}$  for  $m_{\tilde{\nu}} > 250 \text{ GeV}$ .

## $\tilde{\nu}$ (Sneutrino) MASS LIMIT

The limits may depend on the number,  $N(\tilde{\nu})$ , of sneutrinos assumed to be degenerate in mass. Only  $\tilde{\nu}_L$  (not  $\tilde{\nu}_R$ ) is assumed to exist. It is possible that  $\tilde{\nu}$  could be the lightest supersymmetric particle (LSP).

We report here, but do not include in the Listings, the limits obtained from the fit of the final results obtained by the LEP Collaborations on the invisible width of the Z boson ( $\Delta\Gamma_{\text{inv.}} < 2.0 \text{ MeV}$ , LEP-SLC 06):  $m_{\tilde{\nu}} > 43.7 \text{ GeV}$  ( $N(\tilde{\nu})=1$ ) and  $m_{\tilde{\nu}} > 44.7 \text{ GeV}$  ( $N(\tilde{\nu})=3$ ).

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
<b>&gt; 94</b>	95	<sup>1</sup> AAD	11Z ATLS	$\tilde{\nu}_\tau \rightarrow e\mu, \bar{R}$
		<sup>2</sup> ABDALLAH	03M DLPH	$1 \leq \tan\beta \leq 40,$ $m_{\tilde{e}_R} - m_{\tilde{\chi}_1^0} > 10 \text{ GeV}$
> 84	95	<sup>3</sup> HEISTER	02N ALEP	$\tilde{\nu}_e$ , any $\Delta m$
> 37.1	95	<sup>4</sup> ADRIANI	93M L3	$\Gamma(Z \rightarrow \text{invisible}); N(\tilde{\nu})=1$

> 41	95	5	DECAMP	92	ALEP	$\Gamma(Z \rightarrow \text{invisible}); N(\tilde{\nu})=3$
> 36	95		ABREU	91F	DLPH	$\Gamma(Z \rightarrow \text{invisible}); N(\tilde{\nu})=1$
> 31.2	95	6	ALEXANDER	91F	OPAL	$\Gamma(Z \rightarrow \text{invisible}); N(\tilde{\nu})=1$
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●						
		7	AAD	13AI	ATLS	$\tilde{\nu}_\tau \rightarrow e\mu, e\tau, \mu\tau, \mathcal{R}$
		8	AAD	11H	ATLS	$\tilde{\nu}_\tau \rightarrow e\mu, \mathcal{R}$
		9	AALTONEN	10Z	CDF	$\tilde{\nu}_\tau \rightarrow e\mu, e\tau, \mu\tau, \mathcal{R}$
		10	ABAZOV	10M	D0	$\tilde{\nu}_\tau \rightarrow e\mu, \mathcal{R}$
		11	AALTONEN	09V	CDF	$\tilde{\nu} \rightarrow \mu\mu, \mathcal{R} LQ\bar{D}$
		12	ABAZOV	08Q	D0	$\tilde{\nu}_\tau \rightarrow e\mu, \mathcal{R}$
		13	SCHAEF	07A	ALEP	$\tilde{\nu}_{\mu,\tau}, \mathcal{R}, (s+t)\text{-channel}$
		14	ABAZOV	06I	D0	$\mathcal{R}, \lambda'_{211}$
		15	ABDALLAH	06C	DLPH	$\tilde{\nu}_\ell, \mathcal{R}, (s+t)\text{-channel}$
		16	ABULENCIA	06M	CDF	$\tilde{\nu}_\tau \rightarrow e\mu, \mathcal{R}$
		17	ABULENCIA	05A	CDF	$\tilde{\nu} \rightarrow ee, \mu\mu, \mathcal{R} LQ\bar{D}$
		18	ACOSTA	05R	CDF	$\tilde{\nu} \rightarrow \tau\tau, \mathcal{R}, LQ\bar{D}$
		19	ABBIENDI	04F	OPAL	$\mathcal{R}, \tilde{\nu}_{e,\mu,\tau}$
> 95	95	20,21	ABDALLAH	04H	DLPH	AMSB, $\mu > 0$
> 98	95	22	ABDALLAH	04M	DLPH	$\mathcal{R}(LL\bar{E}), \tilde{\nu}_e, \text{indirect}, \Delta m > 5 \text{ GeV}$
> 85	95	22	ABDALLAH	04M	DLPH	$\mathcal{R}(LL\bar{E}), \tilde{\nu}_\mu, \text{indirect}, \Delta m > 5 \text{ GeV}$
> 85	95	22	ABDALLAH	04M	DLPH	$\mathcal{R}(LL\bar{E}), \tilde{\nu}_\tau, \text{indirect}, \Delta m > 5 \text{ GeV}$
		23	ABDALLAH	03F	DLPH	$\tilde{\nu}_{\mu,\tau}, \mathcal{R} LL\bar{E} \text{ decays}$
		24	ACOSTA	03E	CDF	$\tilde{\nu}, \mathcal{R}, LQ\bar{D} \text{ production and } LL\bar{E} \text{ decays}$
> 88	95	25	HEISTER	03G	ALEP	$\tilde{\nu}_e, \mathcal{R} \text{ decays}, \mu = -200 \text{ GeV}, \tan\beta = 2$
> 65	95	25	HEISTER	03G	ALEP	$\tilde{\nu}_{\mu,\tau}, \mathcal{R} \text{ decays}$
		26	ABAZOV	02H	D0	$\mathcal{R}, \lambda'_{211}$
> 95	95	27	ACHARD	02	L3	$\tilde{\nu}_e, \mathcal{R} \text{ decays}, \mu = -200 \text{ GeV}, \tan\beta = \sqrt{2}$
> 65	95	27	ACHARD	02	L3	$\tilde{\nu}_{\nu,\tau}, \mathcal{R} \text{ decays}$
>149	95	27	ACHARD	02	L3	$\tilde{\nu}, \mathcal{R} \text{ decays}, \text{MSUGRA}$
		28	HEISTER	02F	ALEP	$e\gamma \rightarrow \tilde{\nu}_{\mu,\tau} \ell_k, \mathcal{R} LL\bar{E}$
none 100–264	95	29	ABBIENDI	00R	OPAL	$\tilde{\nu}_{\mu,\tau}, \mathcal{R}, (s+t)\text{-channel}$
none 100–200	95	30	ABBIENDI	00R	OPAL	$\tilde{\nu}_\tau, \mathcal{R}, s\text{-channel}$
		31	ABREU	00S	DLPH	$\tilde{\nu}_\ell, \mathcal{R}, (s+t)\text{-channel}$
none 50–210	95	32	ACCIARRI	00P	L3	$\tilde{\nu}_{\mu,\tau}, \mathcal{R}, s\text{-channel}$
none 50–210	95	33	BARATE	00I	ALEP	$\tilde{\nu}_{\mu,\tau}, \mathcal{R}, (s+t)\text{-channel}$
none 90–210	95	34	BARATE	00I	ALEP	$\tilde{\nu}_{\mu,\tau}, \mathcal{R}, s\text{-channel}$
none 100–160	95	35	ABBIENDI	99	OPAL	$\tilde{\nu}_e, \mathcal{R}, t\text{-channel}$
$\neq m_Z$	95	36	ACCIARRI	97U	L3	$\tilde{\nu}_\tau, \mathcal{R}, s\text{-channel}$
none 125–180	95	36	ACCIARRI	97U	L3	$\tilde{\nu}_\tau, \mathcal{R}, s\text{-channel}$
		37	CARENA	97	THEO	$g_\mu - 2$
> 46.0	95	38	BUSKULIC	95E	ALEP	$N(\tilde{\nu})=1, \tilde{\nu} \rightarrow \nu\nu\ell\bar{\ell}'$
none 20–25000		39	BECK	94	COSM	Stable $\tilde{\nu}$ , dark matter
<600		40	FALK	94	COSM	$\tilde{\nu}$ LSP, cosmic abundance
none 3–90	90	41	SATO	91	KAMI	Stable $\tilde{\nu}_e$ or $\tilde{\nu}_\mu$ , dark matter
none 4–90	90	41	SATO	91	KAMI	Stable $\tilde{\nu}_\tau$ , dark matter

- <sup>1</sup> AAD 11Z looked in  $1.07 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with one electron and one muon of opposite charge from the production of  $\tilde{\nu}_\tau$  via an  $\mathcal{R} \lambda'_{311}$  coupling and followed by a decay via  $\lambda_{312}$  into  $e + \mu$ . No evidence for an  $(e, \mu)$  resonance over the SM expectation is observed, and a limit is derived in the plane of  $\lambda'_{311}$  versus  $m_{\tilde{\nu}}$  for three values of  $\lambda_{312}$ , see their Fig. 2. Masses  $m_{\tilde{\nu}} < 1.32$  (1.45) TeV are excluded for  $\lambda'_{311} = 0.10$  and  $\lambda_{312} = 0.05$  ( $\lambda'_{311} = 0.11$  and  $\lambda_{312} = 0.07$ ).
- <sup>2</sup> ABDALLAH 03M uses data from  $\sqrt{s} = 192\text{--}208 \text{ GeV}$  to obtain limits in the framework of the MSSM with gaugino and sfermion mass universality at the GUT scale. An indirect limit on the mass is derived by constraining the MSSM parameter space by the results from direct searches for neutralinos (including cascade decays) and for sleptons. These limits are valid for values of  $M_2 < 1 \text{ TeV}$ ,  $|\mu| \leq 1 \text{ TeV}$  with the  $\tilde{\chi}_1^0$  as LSP. The quoted limit is obtained when there is no mixing in the third family. See Fig. 43 for the mass limits as a function of  $\tan\beta$ . These limits update the results of ABREU 00W.
- <sup>3</sup> HEISTER 02N derives a bound on  $m_{\tilde{\nu}_e}$  by exploiting the mass relation between the  $\tilde{\nu}_e$  and  $\tilde{e}$ , based on the assumption of universal GUT scale gaugino and scalar masses  $m_{1/2}$  and  $m_0$  and the search described in the  $\tilde{e}$  section. In the MSUGRA framework with radiative electroweak symmetry breaking, the limit improves to  $m_{\tilde{\nu}_e} > 130 \text{ GeV}$ , assuming a trilinear coupling  $A_0=0$  at the GUT scale. See Figs. 5 and 7 for the dependence of the limits on  $\tan\beta$ .
- <sup>4</sup> ADRIANI 93M limit from  $\Delta\Gamma(Z)(\text{invisible}) < 16.2 \text{ MeV}$ .
- <sup>5</sup> DECAMP 92 limit is from  $\Gamma(\text{invisible})/\Gamma(\ell\ell) = 5.91 \pm 0.15$  ( $N_\nu = 2.97 \pm 0.07$ ).
- <sup>6</sup> ALEXANDER 91F limit is for one species of  $\tilde{\nu}$  and is derived from  $\Gamma(\text{invisible, new})/\Gamma(\ell\ell) < 0.38$ .
- <sup>7</sup> AAD 13AI searched in  $4.6 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for evidence of heavy particles decaying into  $e\mu$ ,  $e\tau$  or  $\mu\tau$  final states. No significant excess above the Standard Model expectation is observed, and 95% C.L. exclusions are placed on the cross section times branching ratio for the production of an R-parity-violating supersymmetric tau sneutrino, see their Fig. 2. For couplings  $\lambda'_{311} = 0.10$  and  $\lambda_{i3k} = 0.05$ , the lower limits on the  $\tilde{\nu}_\tau$  mass are 1610, 1110, 1100 GeV in the  $e\mu$ ,  $e\tau$ , and  $\mu\tau$  channels, respectively.
- <sup>8</sup> AAD 11H looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with one electron and one muon of opposite charge from the production of  $\tilde{\nu}_\tau$  via an  $\mathcal{R} \lambda'_{311}$  coupling and followed by a decay via  $\lambda_{312}$  into  $e + \mu$ . No evidence for an excess over the SM expectation is observed, and a limit is derived in the plane of  $\lambda'_{311}$  versus  $m_{\tilde{\nu}}$  for several values of  $\lambda_{312}$ , see their Fig. 2. Superseded by AAD 11Z.
- <sup>9</sup> AALTONEN 10Z searched in  $1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events from the production  $d\bar{d} \rightarrow \tilde{\nu}_\tau$  with the subsequent decays  $\tilde{\nu}_\tau \rightarrow e\mu$ ,  $\mu\tau$ ,  $e\tau$  in the MSSM framework with  $\mathcal{R}$ . Two isolated leptons of different flavor and opposite charges are required, with  $\tau$ s identified by their hadronic decay. No statistically significant excesses are observed over the SM background. Upper limits on  $\lambda_{311}^{\prime 2}$  times the branching ratio are listed in their Table III for various  $\tilde{\nu}_\tau$  masses. Limits on the cross section times branching ratio for  $\lambda'_{311} = 0.10$  and  $\lambda_{i3k} = 0.05$ , displayed in Fig. 2, are used to set limits on the  $\tilde{\nu}_\tau$  mass of 558 GeV for the  $e\mu$ , 441 GeV for the  $\mu\tau$  and 442 GeV for the  $e\tau$  channels.
- <sup>10</sup> ABAZOV 10M looked in  $5.3 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with exactly one pair of high  $p_T$  isolated  $e\mu$  and a veto against hard jets. No evidence for an excess over the SM expectation is observed, and a limit at 95% C.L. on the cross section times branching ratio is derived, see their Fig. 3. These limits are translated into limits on couplings as a function of  $m_{\tilde{\nu}_\tau}$  as shown on their Fig. 4. As an example, for  $m_{\tilde{\nu}_\tau} = 100 \text{ GeV}$  and  $\lambda_{312} \leq 0.07$ , couplings  $\lambda'_{311} > 7.7 \times 10^{-4}$  are excluded.
- <sup>11</sup> AALTONEN 09V searched in  $2.3 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with an oppositely charged pair originating from the  $\mathcal{R}$  production of a sneutrino decaying to

- dimuons. A limit is derived on the cross section times branching ratio,  $B$ , of  $\tilde{\nu} \rightarrow \mu\mu$  for several values of the coupling  $\lambda'$ , see their Fig. 3. For  $\lambda'^2 B = 0.01$ , the range  $100 \text{ GeV} \leq m_{\tilde{\nu}} \leq 810 \text{ GeV}$  is excluded.
- 12 ABAZOV 08Q searched in  $1.04 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for an excess of events with oppositely charged  $e\mu$  pairs. They might be expected in a SUSY model with  $\cancel{R}$  where a sneutrino is produced by  $LQ\bar{D}$  couplings and decays via  $LL\bar{E}$  couplings, focusing on  $\tilde{\nu}_\tau$ , hence on the  $\lambda'_{311}$  and  $\lambda_{312}$  constants. No significant excess was found compared to the background expectation. Upper limits on the cross-section times branching ratio are extracted and displayed in their Fig. 2. Exclusion regions are determined for the  $\tilde{\nu}_\tau$  mass as a function of both couplings, see their Fig. 3. As an indication, for  $\tilde{\nu}_\tau$  masses of  $100 \text{ GeV}$  and  $\lambda_{312} = 0.01$ , values of  $\lambda'_{311} \geq 1.6 \times 10^{-3}$  are excluded at the 95% C.L. Superseded by ABAZOV 10M.
  - 13 SCHAEL 07A searches for the  $s$ - or  $t$ -channel exchange of sneutrinos in the case of  $\cancel{R}$  with  $LL\bar{E}$  couplings by studying di-lepton production at  $\sqrt{s} = 189\text{--}209 \text{ GeV}$ . Limits are obtained on the couplings as a function of the  $\tilde{\nu}$  mass, see their Figs. 22-24. The results of this analysis are combined with BARATE 00I.
  - 14 ABAZOV 06I looked in  $380 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with at least 2 muons and 2 jets for  $s$ -channel production of  $\tilde{\mu}$  or  $\tilde{\nu}$  and subsequent decay via  $\cancel{R}$  couplings  $LQ\bar{D}$ . The data are in agreement with the SM expectation. They set limits on resonant slepton production and derive exclusion contours on  $\lambda'_{211}$  in the mass plane of  $\tilde{\ell}$  versus  $\tilde{\chi}_1^0$  assuming a MSUGRA model with  $\tan\beta = 5$ ,  $\mu < 0$  and  $A_0 = 0$ , see their Fig. 3. For  $\lambda'_{211} \geq 0.09$  slepton masses up to  $358 \text{ GeV}$  are excluded. Supersedes the results of ABAZOV 02H.
  - 15 ABDALLAH 06C searches for anomalies in the production cross sections and forward-backward asymmetries of the  $\ell^+\ell^-(\gamma)$  final states ( $\ell=e,\mu,\tau$ ) from  $675 \text{ pb}^{-1}$  of  $e^+e^-$  data at  $\sqrt{s}=130\text{--}207 \text{ GeV}$ . Limits are set on the  $s$ - and  $t$ -channel exchange of sneutrinos in the presence of  $\cancel{R}$  with  $\lambda LL\bar{E}$  couplings. For points between the energies at which data were taken, information is obtained from events in which a photon was radiated. Exclusion limits in the  $(\lambda, m_{\tilde{\nu}})$  plane are given in Fig. 16. These limits include and update the results of ABREU 00S.
  - 16 ABULENCIA 06M searched in  $344 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for an excess of events with oppositely charged  $e\mu$  pairs. They might be expected in a SUSY model with  $\cancel{R}$  where a sneutrino is produced by  $LQ\bar{D}$  couplings and decays via  $LL\bar{E}$  couplings, focusing on  $\tilde{\nu}_\tau$ , hence on the  $\lambda'_{311}$  and  $\lambda_{132}$  constants. No significant excess was found compared to the background expectation. Upper limits on the cross-section times branching ratio are extracted and exclusion regions determined for the  $\tilde{\nu}_\tau$  mass as a function of both couplings, see their Fig. 3. As an indication,  $\tilde{\nu}_\tau$  masses are excluded up to  $300 \text{ GeV}$  for  $\lambda'_{311} \geq 0.01$  and  $\lambda_{132} \geq 0.02$ . Superseded by AALTONEN 10Z.
  - 17 ABULENCIA 05A looked in  $\sim 200 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for dimuon and dielectron events. They may originate from the  $\cancel{R}$  production of a sneutrino decaying to dileptons. No significant excess rate was found compared to the background expectation. A limit is derived on the cross section times branching ratio,  $B$ , of  $\tilde{\nu} \rightarrow ee, \mu\mu$  of  $25 \text{ fb}$  at high mass, see their Figure 2. Sneutrino masses are excluded at 95% CL below  $680, 620, 460 \text{ GeV}$  ( $ee$  channel) and  $665, 590, 450 \text{ GeV}$  ( $\mu\mu$  channel) for a  $\lambda'$  coupling and branching ratio such that  $\lambda'^2 B = 0.01, 0.005, 0.001$ , respectively.
  - 18 ACOSTA 05R looked in  $195 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for ditau events with one identified hadronic tau decay and one other tau decay. They may originate from the  $\cancel{R}$  production of a sneutrino decaying to  $\tau\tau$ . No significant excess rate was found compared to the background expectation, dominated by Drell-Yan. A limit is derived on the cross section times branching ratio,  $B$ , of  $\tilde{\nu} \rightarrow \tau\tau$ , see their Figure 3. Sneutrino masses below  $377 \text{ GeV}$  are excluded at 95% CL for a  $\lambda'$  coupling to  $d\bar{d}$  and branching ratio such that  $\lambda'^2 B = 0.01$ .
  - 19 ABBIENDI 04F use data from  $\sqrt{s} = 189\text{--}209 \text{ GeV}$ . They derive limits on sparticle masses under the assumption of  $\cancel{R}$  with  $LL\bar{E}$  or  $LQ\bar{D}$  couplings. The results are valid for  $\tan\beta$

- = 1.5,  $\mu = -200$  GeV, and a BR for the decay given by CMSSM, assuming no sensitivity to other decays. Limits are quoted for  $m_{\tilde{\chi}^0_1} = 60$  GeV and degrade for low-mass  $\tilde{\chi}^0_1$ . For  $\tilde{\nu}_e$  the direct (indirect) limits with  $LL\bar{E}$  couplings are 89 (95) GeV and with  $LQ\bar{D}$  they are 89 (88) GeV. For  $\tilde{\nu}_{\mu,\tau}$  the direct (indirect) limits with  $LL\bar{E}$  couplings are 79 (81) GeV and with  $LQ\bar{D}$  they are 74 (no limit) GeV. Supersedes the results of ABBIENDI 00.
- <sup>20</sup> ABDALLAH 04H use data from LEP 1 and  $\sqrt{s} = 192\text{--}208$  GeV. They re-use results or re-analyze the data from ABDALLAH 03M to put limits on the parameter space of anomaly-mediated supersymmetry breaking (AMSB), which is scanned in the region  $1 < m_{3/2} < 50$  TeV,  $0 < m_0 < 1000$  GeV,  $1.5 < \tan\beta < 35$ , both signs of  $\mu$ . The constraints are obtained from the searches for mass degenerate chargino and neutralino, for SM-like and invisible Higgs, for leptonically decaying charginos and from the limit on non-SM Z width of 3.2 MeV. The limit is for  $m_t = 174.3$  GeV (see Table 2 for other  $m_t$  values).
- <sup>21</sup> The limit improves to 114 GeV for  $\mu < 0$ .
- <sup>22</sup> ABDALLAH 04M use data from  $\sqrt{s} = 189\text{--}208$  GeV. The results are valid for  $\mu = -200$  GeV,  $\tan\beta = 1.5$ ,  $\Delta m > 5$  GeV and assuming a BR of 1 for the given decay. The limit quoted is for indirect decays using the neutralino constraint of 39.5 GeV, also derived in ABDALLAH 04M. For indirect decays the limit on  $\tilde{\nu}_e$  decreases to 96 GeV if the constraint from the neutralino is not used and for direct decays it remains 96 GeV. For indirect decays the limit on  $\tilde{\nu}_\mu$  decreases to 82 GeV if the constraint from the neutralino is not used and to 83 GeV for direct decays. For indirect decays the limit on  $\tilde{\nu}_\tau$  decreases to 82 GeV if the constraint from the neutralino is not used and improves to 91 GeV for direct decays. Supersedes the results of ABREU 00U.
- <sup>23</sup> ABDALLAH 03F looked for events of the type  $e^+e^- \rightarrow \tilde{\nu} \rightarrow \tilde{\chi}^0_\nu, \tilde{\chi}^\pm \ell^\mp$  followed by  $\tilde{R}$  decays of the  $\tilde{\chi}^0$  via  $\lambda_{1j1}$  ( $j = 2,3$ ) couplings in the data at  $\sqrt{s} = 183\text{--}208$  GeV. From a scan over the SUGRA parameters, they derive upper limits on the  $\lambda_{1j1}$  couplings as a function of the sneutrino mass, see their Figs. 5–8.
- <sup>24</sup> ACOSTA 03E search for  $e\mu, e\tau$  and  $\mu\tau$  final states, and sets limits on the product of production cross-section and decay branching ratio for a  $\tilde{\nu}$  in RPV models (see Fig. 3).
- <sup>25</sup> HEISTER 03G searches for the production of sneutrinos in the case of  $\tilde{R}$  prompt decays with  $LL\bar{E}$ ,  $LQ\bar{D}$  or  $UDD$  couplings at  $\sqrt{s} = 189\text{--}209$  GeV. The search is performed for direct and indirect decays, assuming one coupling at a time to be non-zero. The limit holds for indirect  $\tilde{\nu}$  decays via  $UDD$  couplings and  $\Delta m > 10$  GeV. Stronger limits are reached for  $(\tilde{\nu}_e, \tilde{\nu}_{\mu,\tau})$  for  $LL\bar{E}$  direct (100,90) GeV or indirect (98,89) GeV and for  $LQ\bar{D}$  direct (–,79) GeV or indirect (91,78) GeV couplings. For  $LL\bar{E}$  indirect decays, use is made of the bound  $m(\tilde{\chi}^0_1) > 23$  GeV from BARATE 98S. Supersedes the results from BARATE 01B.
- <sup>26</sup> ABAZOV 02H looked in  $94 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8$  TeV for events with at least 2 muons and 2 jets for  $s$ -channel production of  $\tilde{\mu}$  or  $\tilde{\nu}$  and subsequent decay via  $\tilde{R}$  couplings  $LQ\bar{D}$ . A scan over the MSUGRA parameters is performed to exclude regions of the  $(m_0, m_{1/2})$  plane, examples being shown in Fig. 2.
- <sup>27</sup> ACHARD 02 searches for the associated production of sneutrinos in the case of  $\tilde{R}$  prompt decays with  $LL\bar{E}$  or  $UDD$  couplings at  $\sqrt{s}=189\text{--}208$  GeV. The search is performed for direct and indirect decays, assuming one coupling at the time to be nonzero. The limit holds for direct decays via  $LL\bar{E}$  couplings. Stronger limits are reached for  $(\tilde{\nu}_e, \tilde{\nu}_{\mu,\tau})$  for  $LL\bar{E}$  indirect (99,78) GeV and for  $UDD$  direct or indirect (99,70) GeV decays. The MSUGRA limit results from a scan over the MSSM parameter space with the assumption of gaugino and scalar mass unification at the GUT scale, imposing simultaneously the exclusions from neutralino, chargino, sleptons, and squarks analyses. The limit holds for  $UDD$  couplings and increases to 152.7 GeV for  $LL\bar{E}$  couplings.
- <sup>28</sup> HEISTER 02F searched for single sneutrino production via  $e\gamma \rightarrow \tilde{\nu}_j \ell_k$  mediated by  $\tilde{R}$   $LL\bar{E}$  couplings, decaying directly or indirectly via a  $\tilde{\chi}^0_1$  and assuming a single coupling to be nonzero at a time. Final states with three leptons and possible  $\cancel{E}_T$  due to neutrinos were selected in the 189–209 GeV data. Limits on the couplings  $\lambda_{1jk}$  as function of

- the sneutrino mass are shown in Figs. 10–14. The couplings  $\lambda_{232}$  and  $\lambda_{233}$  are not accessible and  $\lambda_{121}$  and  $\lambda_{131}$  are measured with better accuracy in sneutrino resonant production. For all tested couplings, except  $\lambda_{133}$ , the limits are significantly improved compared to the low-energy limits.
- 29 ABBIENDI 00R studied the effect of  $s$ - and  $t$ -channel  $\tau$  or  $\mu$  sneutrino exchange in  $e^+e^- \rightarrow e^+e^-$  at  $\sqrt{s}=130\text{--}189$  GeV, via the  $R$ -parity violating coupling  $\lambda_{1j1}L_1L_je_1$  ( $i=2$  or  $3$ ). The limits quoted here hold for  $\lambda_{1j1} > 0.13$ , and supersede the results of ABBIENDI 99. See Fig. 11 for limits on  $m_{\tilde{\nu}}$  versus coupling.
  - 30 ABBIENDI 00R studied the effect of  $s$ -channel  $\tau$  sneutrino exchange in  $e^+e^- \rightarrow \mu^+\mu^-$  at  $\sqrt{s}=130\text{--}189$  GeV, in presence of the  $R$ -parity violating couplings  $\lambda_{i3j}L_iL_3e_j$  ( $i=1$  and  $2$ ), with  $\lambda_{131}=\lambda_{232}$ . The limits quoted here hold for  $\lambda_{131} > 0.09$ , and supersede the results of ABBIENDI 99. See Fig. 12 for limits on  $m_{\tilde{\nu}}$  versus coupling.
  - 31 ABREU 00S searches for anomalies in the production cross sections and forward-backward asymmetries of the  $\ell^+\ell^-(\gamma)$  final states ( $\ell=e,\mu,\tau$ ) from  $e^+e^-$  collisions at  $\sqrt{s}=130\text{--}189$  GeV. Limits are set on the  $s$ - and  $t$ -channel exchange of sneutrinos in the presence of  $\mathcal{R}$  with  $\lambda LL\bar{E}$  couplings. For points between the energies at which data were taken, information is obtained from events in which a photon was radiated. Exclusion limits in the  $(\lambda, m_{\tilde{\nu}})$  plane are given in Fig. 5. These limits include and update the results of ABREU 99A.
  - 32 ACCIARRI 00P use the dilepton total cross sections and asymmetries at  $\sqrt{s}=m_Z$  and  $\sqrt{s}=130\text{--}189$  GeV data to set limits on the effect of  $\mathcal{R} LL\bar{E}$  couplings giving rise to  $\mu$  or  $\tau$  sneutrino exchange. See their Fig. 5 for limits on the sneutrino mass versus couplings.
  - 33 BARATE 00i studied the effect of  $s$ -channel and  $t$ -channel  $\tau$  or  $\mu$  sneutrino exchange in  $e^+e^- \rightarrow e^+e^-$  at  $\sqrt{s}=130\text{--}183$  GeV, via the  $R$ -parity violating coupling  $\lambda_{1j1}L_1L_je_1^C$  ( $i=2$  or  $3$ ). The limits quoted here hold for  $\lambda_{1j1} > 0.1$ . See their Fig. 15 for limits as a function of the coupling. Superseded by SCHÄEL 07A.
  - 34 BARATE 00i studied the effect of  $s$ -channel  $\tau$  sneutrino exchange in  $e^+e^- \rightarrow \mu^+\mu^-$  at  $\sqrt{s}=130\text{--}183$  GeV, in presence of the  $R$ -parity violating coupling  $\lambda_{i3j}L_iL_3e_i^C$  ( $i=1$  and  $2$ ). The limits quoted here hold for  $\sqrt{|\lambda_{131}\lambda_{232}|} > 0.2$ . See their Fig. 16 for limits as a function of the coupling. Superseded by SCHÄEL 07A.
  - 35 ABBIENDI 99 studied the effect of  $t$ -channel electron sneutrino exchange in  $e^+e^- \rightarrow \tau^+\tau^-$  at  $\sqrt{s}=130\text{--}183$  GeV, in presence of the  $R$ -parity violating couplings  $\lambda_{131}L_1L_3e_1^C$ . The limits quoted here hold for  $\lambda_{131} > 0.6$ .
  - 36 ACCIARRI 97U studied the effect of the  $s$ -channel tau-sneutrino exchange in  $e^+e^- \rightarrow e^+e^-$  at  $\sqrt{s}=m_Z$  and  $\sqrt{s}=130\text{--}172$  GeV, via the  $R$ -parity violating coupling  $\lambda_{131}L_1L_3e_1^C$ . The limits quoted here hold for  $\lambda_{131} > 0.05$ . Similar limits were studied in  $e^+e^- \rightarrow \mu^+\mu^-$  together with  $\lambda_{232}L_2L_3e_2^C$  coupling.
  - 37 CARENA 97 studied the constraints on chargino and sneutrino masses from muon  $g-2$ . The bound can be important for large  $\tan\beta$ .
  - 38 BUSKULIC 95E looked for  $Z \rightarrow \tilde{\nu}\bar{\nu}$ , where  $\tilde{\nu} \rightarrow \nu\chi_1^0$  and  $\chi_1^0$  decays via  $R$ -parity violating interactions into two leptons and a neutrino.
  - 39 BECK 94 limit can be inferred from limit on Dirac neutrino using  $\sigma(\tilde{\nu}) = 4\sigma(\nu)$ . Also private communication with H.V. Klapdor-Kleingrothaus.
  - 40 FALK 94 puts an upper bound on  $m_{\tilde{\nu}}$  when  $\tilde{\nu}$  is LSP by requiring its relic density does not overclose the Universe.
  - 41 SATO 91 search for high-energy neutrinos from the sun produced by annihilation of sneutrinos in the sun. Sneutrinos are assumed to be stable and to constitute dark matter in our galaxy. SATO 91 follow the analysis of NG 87, OLIVE 88, and GAISSER 86.
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## CHARGED SLEPTONS

This section contains limits on charged scalar leptons ( $\tilde{\ell}$ , with  $\ell=e,\mu,\tau$ ). Studies of width and decays of the  $Z$  boson (use is made here of  $\Delta\Gamma_{\text{inv}} < 2.0$  MeV, LEP 00) conclusively rule out  $m_{\tilde{\ell}_R} < 40$  GeV (41

GeV for  $\tilde{\ell}_L$ ), independently of decay modes, for each individual slepton. The limits improve to 43 GeV (43.5 GeV for  $\tilde{\ell}_L$ ) assuming all 3 flavors to be degenerate. Limits on higher mass sleptons depend on model assumptions and on the mass splitting  $\Delta m = m_{\tilde{\ell}} - m_{\tilde{\chi}_1^0}$ . The mass and composition

of  $\tilde{\chi}_1^0$  may affect the selectron production rate in  $e^+e^-$  collisions through  $t$ -channel exchange diagrams. Production rates are also affected by the potentially large mixing angle of the lightest mass eigenstate  $\tilde{\ell}_1 = \tilde{\ell}_R \sin\theta_\ell + \tilde{\ell}_L \cos\theta_\ell$ . It is generally assumed that only  $\tilde{\tau}$  may have significant mixing. The coupling to the  $Z$  vanishes for  $\theta_\ell=0.82$ . In the high-energy limit of  $e^+e^-$  collisions the interference between  $\gamma$  and  $Z$  exchange leads to a minimal cross section for  $\theta_\ell=0.91$ , a value which is sometimes used in the following entries relative to data taken at LEP2. When limits on  $m_{\tilde{\ell}_R}$  are quoted, it is understood that limits on  $m_{\tilde{\ell}_L}$  are usually at least as strong.

Possibly open decays involving gauginos other than  $\tilde{\chi}_1^0$  will affect the detection efficiencies. Unless otherwise stated, the limits presented here result from the study of  $\tilde{\ell}^+\tilde{\ell}^-$  production, with production rates and decay properties derived from the MSSM. Limits made obsolete by the recent analyses of  $e^+e^-$  collisions at high energies can be found in previous Editions of this Review.

For decays with final state gravitinos ( $\tilde{G}$ ),  $m_{\tilde{G}}$  is assumed to be negligible relative to all other masses.

### $\tilde{e}$ (Selectron) MASS LIMIT

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
		1 AAD	13B ATLS	$2\ell^\pm + \cancel{E}_T$ , SMS, pMSSM
> 97.5		2 ABBIENDI	04 OPAL	$\tilde{e}_R, \Delta m > 11$ GeV, $ \mu  > 100$ GeV, $\tan\beta=1.5$
> 94.4		3 ACHARD	04 L3	$\tilde{e}_R, \Delta m > 10$ GeV, $ \mu  > 200$ GeV, $\tan\beta \geq 2$
> 71.3		3 ACHARD	04 L3	$\tilde{e}_R$ , all $\Delta m$
none 30–94	95	4 ABDALLAH	03M DLPH	$\Delta m > 15$ GeV, $\tilde{e}_R^+ \tilde{e}_R^-$
> 94	95	5 ABDALLAH	03M DLPH	$\tilde{e}_R, 1 \leq \tan\beta \leq 40, \Delta m > 10$ GeV
> 95	95	6 HEISTER	02E ALEP	$\Delta m > 15$ GeV, $\tilde{e}_R^+ \tilde{e}_R^-$
<b>&gt; 73</b>	95	7 HEISTER	02N ALEP	$\tilde{e}_R$ , any $\Delta m$
<b>&gt;107</b>	95	7 HEISTER	02N ALEP	$\tilde{e}_L$ , any $\Delta m$
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
> 89	95	8 ABBIENDI	04F OPAL	$\cancel{R}, \tilde{e}_L$
> 92	95	9 ABDALLAH	04M DLPH	$\cancel{R}, \tilde{e}_R$ , indirect, $\Delta m > 5$ GeV
> 93	95	10 HEISTER	03G ALEP	$\tilde{e}_R, \cancel{R}$ decays, $\mu = -200$ GeV, $\tan\beta=2$
> 69	95	11 ACHARD	02 L3	$\tilde{e}_R, \cancel{R}$ decays, $\mu = -200$ GeV, $\tan\beta=\sqrt{2}$
> 92	95	12 BARATE	01 ALEP	$\Delta m > 10$ GeV, $\tilde{e}_R^+ \tilde{e}_R^-$

> 77	95	13	ABBIENDI	00J	OPAL	$\Delta m > 5 \text{ GeV}, \tilde{e}_R^+ \tilde{e}_R^-$
> 83	95	14	ABREU	00U	DLPH	$\tilde{e}_R, \cancel{E} (LL\bar{E})$
> 67	95	15	ABREU	00V	DLPH	$\tilde{e}_R \tilde{e}_R (\tilde{e}_R \rightarrow e \tilde{G}), m_{\tilde{G}} > 10 \text{ eV}$
> 85	95	16	BARATE	00G	ALEP	$\tilde{\ell}_R \rightarrow \ell \tilde{G}, \text{ any } \tau(\tilde{\ell}_R)$
> 29.5	95	17	ACCIARRI	99I	L3	$\tilde{e}_R, \cancel{E}, \tan\beta \geq 2$
> 56	95	18	ACCIARRI	98F	L3	$\Delta m > 5 \text{ GeV}, \tilde{e}_R^+ \tilde{e}_R^-, \tan\beta \geq 1.41$
> 77	95	19	BARATE	98K	ALEP	Any $\Delta m, \tilde{e}_R^+ \tilde{e}_R^-, \tilde{e}_R \rightarrow e \gamma \tilde{G}$
> 77	95	20	BREITWEG	98	ZEUS	$m_{\tilde{q}} = m_{\tilde{e}}, m(\tilde{\chi}_1^0) = 40 \text{ GeV}$
> 63	95	21	AID	96C	H1	$m_{\tilde{q}} = m_{\tilde{e}}, m_{\tilde{\chi}_1^0} = 35 \text{ GeV}$

<sup>1</sup> AAD 13B searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for sleptons decaying to a final state with two leptons ( $e$  and  $\mu$ ) and missing transverse energy. No excess beyond the Standard Model expectation is observed. Limits are derived in a simplified model of direct left-handed slepton pair production, where left-handed slepton masses between 85 and 195 GeV are excluded at 95% C.L. for  $m_{\tilde{\chi}_1^0} = 20 \text{ GeV}$ . See also Fig. 2(a). Exclusion

limits are also derived in the phenomenological MSSM, see Fig. 3.

<sup>2</sup> ABBIENDI 04 search for  $\tilde{e}_R \tilde{e}_R$  production in acoplanar di-electron final states in the 183–208 GeV data. See Fig. 13 for the dependence of the limits on  $m_{\tilde{\chi}_1^0}$  and for the

limit at  $\tan\beta=35$ . This limit supersedes ABBIENDI 00G.

<sup>3</sup> ACHARD 04 search for  $\tilde{e}_R \tilde{e}_L$  and  $\tilde{e}_R \tilde{e}_R$  production in single- and acoplanar di-electron final states in the 192–209 GeV data. Absolute limits on  $m_{\tilde{e}_R}$  are derived from a scan over the MSSM parameter space with universal GUT scale gaugino and scalar masses  $m_{1/2}$  and  $m_0$ ,  $1 \leq \tan\beta \leq 60$  and  $-2 \leq \mu \leq 2 \text{ TeV}$ . See Fig. 4 for the dependence of the limits on  $m_{\tilde{\chi}_1^0}$ . This limit supersedes ACCIARRI 99W.

<sup>4</sup> ABDALLAH 03M looked for acoplanar dielectron  $+\cancel{E}$  final states at  $\sqrt{s} = 189\text{--}208 \text{ GeV}$ . The limit assumes  $\mu = -200 \text{ GeV}$  and  $\tan\beta=1.5$  in the calculation of the production cross section and  $B(\tilde{e} \rightarrow e \tilde{\chi}_1^0)$ . See Fig. 15 for limits in the  $(m_{\tilde{e}_R}, m_{\tilde{\chi}_1^0})$  plane. These limits include and update the results of ABREU 01

<sup>5</sup> ABDALLAH 03M uses data from  $\sqrt{s} = 192\text{--}208 \text{ GeV}$  to obtain limits in the framework of the MSSM with gaugino and sfermion mass universality at the GUT scale. An indirect limit on the mass is derived by constraining the MSSM parameter space by the results from direct searches for neutralinos (including cascade decays) and for sleptons. These limits are valid for values of  $M_2 < 1 \text{ TeV}$ ,  $|\mu| \leq 1 \text{ TeV}$  with the  $\tilde{\chi}_1^0$  as LSP. The quoted limit is obtained when there is no mixing in the third family. See Fig. 43 for the mass limits as a function of  $\tan\beta$ . These limits update the results of ABREU 00W.

<sup>6</sup> HEISTER 02E looked for acoplanar dielectron  $+\cancel{E}_T$  final states from  $e^+ e^-$  interactions between 183 and 209 GeV. The mass limit assumes  $\mu < -200 \text{ GeV}$  and  $\tan\beta=2$  for the production cross section and  $B(\tilde{e} \rightarrow e \tilde{\chi}_1^0)=1$ . See their Fig. 4 for the dependence of the limit on  $\Delta m$ . These limits include and update the results of BARATE 01.

<sup>7</sup> HEISTER 02N search for  $\tilde{e}_R \tilde{e}_L$  and  $\tilde{e}_R \tilde{e}_R$  production in single- and acoplanar di-electron final states in the 183–208 GeV data. Absolute limits on  $m_{\tilde{e}_R}$  are derived from a scan over the MSSM parameter space with universal GUT scale gaugino and scalar masses  $m_{1/2}$  and  $m_0$ ,  $1 \leq \tan\beta \leq 50$  and  $-10 \leq \mu \leq 10 \text{ TeV}$ . The region of small  $|\mu|$ , where cascade decays are important, is covered by a search for  $\tilde{\chi}_1^0 \tilde{\chi}_3^0$  in final states with leptons and possibly photons. Limits on  $m_{\tilde{e}_L}$  are derived by exploiting the mass relation between the  $\tilde{e}_L$  and  $\tilde{e}_R$ , based on universal  $m_0$  and  $m_{1/2}$ . When the constraint from the mass limit of the lightest Higgs from HEISTER 02 is included, the bounds improve to  $m_{\tilde{e}_R} > 77(75) \text{ GeV}$  and  $m_{\tilde{e}_L} > 115(115) \text{ GeV}$  for a top mass of 175(180) GeV. In the

- MSUGRA framework with radiative electroweak symmetry breaking, the limits improve further to  $m_{\tilde{e}_R} > 95$  GeV and  $m_{\tilde{e}_L} > 152$  GeV, assuming a trilinear coupling  $A_0=0$  at the GUT scale. See Figs. 4, 5, 7 for the dependence of the limits on  $\tan\beta$ .
- 8 ABBIENDI 04F use data from  $\sqrt{s} = 189\text{--}209$  GeV. They derive limits on sparticle masses under the assumption of  $\tilde{R}$  with  $LL\bar{E}$  or  $LQ\bar{D}$  couplings. The results are valid for  $\tan\beta = 1.5$ ,  $\mu = -200$  GeV, with, in addition,  $\Delta m > 5$  GeV for indirect decays via  $LQ\bar{D}$ . The limit quoted is for indirect decays via  $LL\bar{E}$  or  $LQ\bar{D}$  couplings. For indirect decays, the limits on the  $\tilde{e}_R$  mass are respectively 99 and 92 GeV for  $LL\bar{E}$  and  $LQ\bar{D}$  couplings and  $m_{\tilde{\chi}_1^0} = 10$  GeV and degrade slightly for larger  $\tilde{\chi}_1^0$  mass. Supersedes the results of ABBIENDI 00.
  - 9 ABDALLAH 04M use data from  $\sqrt{s} = 192\text{--}208$  GeV to derive limits on sparticle masses under the assumption of  $\tilde{R}$  with  $LL\bar{E}$  or  $U\bar{D}\bar{D}$  couplings. The results are valid for  $\mu = -200$  GeV,  $\tan\beta = 1.5$ ,  $\Delta m > 5$  GeV and assuming a BR of 1 for the given decay. The limit quoted is for indirect  $U\bar{D}\bar{D}$  decays using the neutralino constraint of 39.5 GeV for  $LL\bar{E}$  and of 38.0 GeV for  $U\bar{D}\bar{D}$  couplings, also derived in ABDALLAH 04M. For indirect decays via  $LL\bar{E}$  the limit improves to 95 GeV if the constraint from the neutralino is used and to 94 GeV if it is not used. For indirect decays via  $U\bar{D}\bar{D}$  couplings it remains unchanged when the neutralino constraint is not used. Supersedes the result of ABREU 00U.
  - 10 HEISTER 03G searches for the production of selectrons in the case of  $\tilde{R}$  prompt decays with  $LL\bar{E}$ ,  $LQ\bar{D}$  or  $U\bar{D}\bar{D}$  couplings at  $\sqrt{s} = 189\text{--}209$  GeV. The search is performed for direct and indirect decays, assuming one coupling at a time to be non-zero. The limit holds for indirect decays mediated by  $LQ\bar{D}$  couplings with  $\Delta m > 10$  GeV. Limits are also given for  $LL\bar{E}$  direct ( $m_{\tilde{e},R} > 96$  GeV) and indirect decays ( $m_{\tilde{e},R} > 96$  GeV for  $m(\tilde{\chi}_1^0) > 23$  GeV from BARATE 98s) and for  $U\bar{D}\bar{D}$  indirect decays ( $m_{\tilde{e},R} > 94$  GeV with  $\Delta m > 10$  GeV). Supersedes the results from BARATE 01B.
  - 11 ACHARD 02 searches for the production of selectrons in the case of  $\tilde{R}$  prompt decays with  $LL\bar{E}$  or  $U\bar{D}\bar{D}$  couplings at  $\sqrt{s}=189\text{--}208$  GeV. The search is performed for direct and indirect decays, assuming one coupling at the time to be nonzero. The limit holds for direct decays via  $LL\bar{E}$  couplings. Stronger limits are reached for  $LL\bar{E}$  indirect (79 GeV) and for  $U\bar{D}\bar{D}$  direct or indirect (96 GeV) decays.
  - 12 BARATE 01 looked for acoplanar dielectron +  $\cancel{E}_T$  final states at 189 to 202 GeV. The limit assumes  $\mu=-200$  GeV and  $\tan\beta=2$  for the production cross section and 100% branching ratio for  $\tilde{e} \rightarrow e\tilde{\chi}_1^0$ . See their Fig. 1 for the dependence of the limit on  $\Delta m$ . These limits include and update the results of BARATE 99Q.
  - 13 ABBIENDI 00J looked for acoplanar dielectron +  $\cancel{E}_T$  final states at  $\sqrt{s}=161\text{--}183$  GeV. The limit assumes  $\mu < -100$  GeV and  $\tan\beta=1.5$  for the production cross section and decay branching ratios, evaluated within the MSSM, and zero efficiency for decays other than  $\tilde{e} \rightarrow e\tilde{\chi}_1^0$ . See their Fig. 12 for the dependence of the limit on  $\Delta m$  and  $\tan\beta$ .
  - 14 ABREU 00U studies decays induced by  $R$ -parity violating  $LL\bar{E}$  couplings, using data from  $\sqrt{s}=189$  GeV. They investigate topologies with multiple leptons, assuming one coupling at the time to be nonzero and giving rise to indirect decays. The limits assume a neutralino mass limit of 30 GeV, also derived in ABREU 00U. Updates ABREU 00I. Superseded by ABDALLAH 04M.
  - 15 ABREU 00V use data from  $\sqrt{s}=130\text{--}189$  GeV to search for tracks with large impact parameter or visible decay vertices. Limits are obtained as a function of  $m_{\tilde{G}}$ , from a scan of the GMSB parameters space, after combining these results with the search for slepton pair production in the SUGRA framework from ABREU 01 to cover prompt decays and on stable particle searches from ABREU 00Q. For limits at different  $m_{\tilde{G}}$ , see their Fig. 12.
  - 16 BARATE 00G combines the search for acoplanar dileptons, leptons with large impact parameters, kinks, and stable heavy-charged tracks, assuming 3 flavors of degenerate sleptons, produced in the  $s$  channel. Data collected at  $\sqrt{s}=189$  GeV.
  - 17 ACCIARRI 99i establish indirect limits on  $m_{\tilde{e}_R}$  from the regions excluded in the  $M_2$  versus  $m_0$  plane by their chargino and neutralino searches at  $\sqrt{s}=130\text{--}183$  GeV. The

- situations where the  $\tilde{\chi}_1^0$  is the LSP (indirect decays) and where a  $\tilde{\ell}$  is the LSP (direct decays) were both considered. The weakest limit, quoted above, comes from direct decays with  $\overline{UDD}$  couplings;  $LL\overline{E}$  couplings or indirect decays lead to a stronger limit.
- 18 ACCIARRI 98F looked for acoplanar dielectron+ $\cancel{E}_T$  final states at  $\sqrt{s}=130\text{--}172$  GeV. The limit assumes  $\mu=-200$  GeV, and zero efficiency for decays other than  $\tilde{e}_R \rightarrow e\tilde{\chi}_1^0$ . See their Fig. 6 for the dependence of the limit on  $\Delta m$ .
- 19 BARATE 98K looked for  $e^+e^-\gamma\gamma + \cancel{E}$  final states at  $\sqrt{s}=161\text{--}184$  GeV. The limit assumes  $\mu=-200$  GeV and  $\tan\beta=2$  for the evaluation of the production cross section. See Fig. 4 for limits on the  $(m_{\tilde{e}_R}, m_{\tilde{\chi}_1^0})$  plane and for the effect of cascade decays.
- 20 BREITWEG 98 used positron+jet events with missing energy and momentum to look for  $e^+q \rightarrow \tilde{e}\tilde{q}$  via gaugino-like neutralino exchange with decays into  $(e\tilde{\chi}_1^0)(q\tilde{\chi}_1^0)$ . See paper for dependences in  $m(\tilde{q})$ ,  $m(\tilde{\chi}_1^0)$ .
- 21 AID 96C used positron+jet events with missing energy and momentum to look for  $e^+q \rightarrow \tilde{e}\tilde{q}$  via neutralino exchange with decays into  $(e\tilde{\chi}_1^0)(q\tilde{\chi}_1^0)$ . See the paper for dependences on  $m_{\tilde{q}}$ ,  $m_{\tilde{\chi}_1^0}$ .

### $\tilde{\mu}$ (Smuon) MASS LIMIT

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
		1 AAD	13B ATLS	$2\ell^\pm + \cancel{E}_T$ , SMS, pMSSM
>91.0		2 ABBIENDI	04 OPAL	$\Delta m > 3$ GeV, $\tilde{\mu}_R^+\tilde{\mu}_R^-$ , $ \mu  > 100$ GeV, $\tan\beta=1.5$
>86.7		3 ACHARD	04 L3	$\Delta m > 10$ GeV, $\tilde{\mu}_R^+\tilde{\mu}_R^-$ , $ \mu  > 200$ GeV, $\tan\beta \geq 2$
none 30–88	95	4 ABDALLAH	03M DLPH	$\Delta m > 5$ GeV, $\tilde{\mu}_R^+\tilde{\mu}_R^-$
<b>&gt;94</b>	95	5 ABDALLAH	03M DLPH	$\tilde{\mu}_{R,1} \leq \tan\beta \leq 40$ , $\Delta m > 10$ GeV
>88	95	6 HEISTER	02E ALEP	$\Delta m > 15$ GeV, $\tilde{\mu}_R^+\tilde{\mu}_R^-$
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
		7 ABAZOV	06i D0	$\cancel{R}$ , $\lambda'_{211}$
>74	95	8 ABBIENDI	04F OPAL	$\cancel{R}$ , $\tilde{\mu}_L$
>87	95	9 ABDALLAH	04M DLPH	$\cancel{R}$ , $\tilde{\mu}_R$ , indirect, $\Delta m > 5$ GeV
>81	95	10 HEISTER	03G ALEP	$\tilde{\mu}_L, \cancel{R}$ decays
		11 ABAZOV	02H D0	$\cancel{R}$ , $\lambda'_{211}$
>61	95	12 ACHARD	02 L3	$\tilde{\mu}_R$ , $\cancel{R}$ decays
>85	95	13 BARATE	01 ALEP	$\Delta m > 10$ GeV, $\tilde{\mu}_R^+\tilde{\mu}_R^-$
>65	95	14 ABBIENDI	00J OPAL	$\Delta m > 2$ GeV, $\tilde{\mu}_R^+\tilde{\mu}_R^-$
>80	95	15 ABREU	00V DLPH	$\tilde{\mu}_R\tilde{\mu}_R$ ( $\tilde{\mu}_R \rightarrow \mu\tilde{G}$ ), $m_{\tilde{G}} > 8$ eV
>77	95	16 BARATE	98K ALEP	Any $\Delta m$ , $\tilde{\mu}_R^+\tilde{\mu}_R^-$ , $\tilde{\mu}_R \rightarrow \mu\gamma\tilde{G}$

<sup>1</sup> AAD 13B searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for sleptons decaying to a final state with two leptons ( $e$  and  $\mu$ ) and missing transverse energy. No excess beyond the Standard Model expectation is observed. Limits are derived in a simplified model of direct left-handed slepton pair production, where left-handed slepton masses between 85 and 195 GeV are excluded at 95% C.L. for  $m_{\tilde{\chi}_1^0} = 20$  GeV. See also Fig. 2(a). Exclusion limits are also derived in the phenomenological MSSM, see Fig. 3.

<sup>2</sup> ABBIENDI 04 search for  $\tilde{\mu}_R\tilde{\mu}_R$  production in acoplanar di-muon final states in the 183–208 GeV data. See Fig. 14 for the dependence of the limits on  $m_{\tilde{\chi}_1^0}$  and for the

- limit at  $\tan\beta=35$ . Under the assumption of 100% branching ratio for  $\tilde{\mu}_R \rightarrow \mu \tilde{\chi}_1^0$ , the limit improves to 94.0 GeV for  $\Delta m > 4$  GeV. See Fig. 11 for the dependence of the limits on  $m_{\tilde{\chi}_1^0}$  at several values of the branching ratio. This limit supersedes ABBIENDI 00G.
- <sup>3</sup> ACHARD 04 search for  $\tilde{\mu}_R \tilde{\mu}_R$  production in acoplanar di-muon final states in the 192–209 GeV data. Limits on  $m_{\tilde{\mu}_R}$  are derived from a scan over the MSSM parameter space with universal GUT scale gaugino and scalar masses  $m_{1/2}$  and  $m_0$ ,  $1 \leq \tan\beta \leq 60$  and  $-2 \leq \mu \leq 2$  TeV. See Fig. 4 for the dependence of the limits on  $m_{\tilde{\chi}_1^0}$ . This limit supersedes ACCIARRI 99W.
  - <sup>4</sup> ABDALLAH 03M looked for acoplanar dimuon  $+ \cancel{E}$  final states at  $\sqrt{s} = 189$ –208 GeV. The limit assumes  $B(\tilde{\mu} \rightarrow \mu \tilde{\chi}_1^0) = 100\%$ . See Fig. 16 for limits on the  $(m_{\tilde{\mu}_R}, m_{\tilde{\chi}_1^0})$  plane. These limits include and update the results of ABREU 01.
  - <sup>5</sup> ABDALLAH 03M uses data from  $\sqrt{s} = 192$ –208 GeV to obtain limits in the framework of the MSSM with gaugino and sfermion mass universality at the GUT scale. An indirect limit on the mass is derived by constraining the MSSM parameter space by the results from direct searches for neutralinos (including cascade decays) and for sleptons. These limits are valid for values of  $M_2 < 1$  TeV,  $|\mu| \leq 1$  TeV with the  $\tilde{\chi}_1^0$  as LSP. The quoted limit is obtained when there is no mixing in the third family. See Fig. 43 for the mass limits as a function of  $\tan\beta$ . These limits update the results of ABREU 00W.
  - <sup>6</sup> HEISTER 02E looked for acoplanar dimuon  $+ \cancel{E}_T$  final states from  $e^+e^-$  interactions between 183 and 209 GeV. The mass limit assumes  $B(\tilde{\mu} \rightarrow \mu \tilde{\chi}_1^0)=1$ . See their Fig. 4 for the dependence of the limit on  $\Delta m$ . These limits include and update the results of BARATE 01.
  - <sup>7</sup> ABAZOV 06i looked in  $380 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV for events with at least 2 muons and 2 jets for s-channel production of  $\tilde{\mu}$  or  $\tilde{\nu}$  and subsequent decay via  $\cancel{R}$  couplings  $LQ\bar{D}$ . The data are in agreement with the SM expectation. They set limits on resonant slepton production and derive exclusion contours on  $\lambda'_{211}$  in the mass plane of  $\tilde{\ell}$  versus  $\tilde{\chi}_1^0$  assuming a MSUGRA model with  $\tan\beta = 5$ ,  $\mu < 0$  and  $A_0 = 0$ , see their Fig. 3. For  $\lambda'_{211} \geq 0.09$  slepton masses up to 358 GeV are excluded. Supersedes the results of ABAZOV 02H.
  - <sup>8</sup> ABBIENDI 04F use data from  $\sqrt{s} = 189$ –209 GeV. They derive limits on sparticle masses under the assumption of  $\cancel{R}$  with  $LL\bar{E}$  or  $LQ\bar{D}$  couplings. The results are valid for  $\tan\beta = 1.5$ ,  $\mu = -200$  GeV, with, in addition,  $\Delta m > 5$  GeV for indirect decays via  $LQ\bar{D}$ . The limit quoted applies to direct decays with  $LL\bar{E}$  couplings and improves to 75 GeV for  $LQ\bar{D}$  couplings. The limits on the  $\tilde{\mu}_R$  mass for indirect decays are respectively 94 and 87 GeV for  $LL\bar{E}$  and  $LQ\bar{D}$  couplings and  $m_{\tilde{\chi}_1^0} = 10$  GeV. Supersedes the results of ABBIENDI 00.
  - <sup>9</sup> ABDALLAH 04M use data from  $\sqrt{s} = 192$ –208 GeV to derive limits on sparticle masses under the assumption of  $\cancel{R}$  with  $LL\bar{E}$  or  $U\bar{D}\bar{D}$  couplings. The results are valid for  $\mu = -200$  GeV,  $\tan\beta = 1.5$ ,  $\Delta m \geq 5$  GeV and assuming a BR of 1 for the given decay. The limit quoted is for indirect  $U\bar{D}\bar{D}$  decays using the neutralino constraint of 39.5 GeV for  $LL\bar{E}$  and of 38.0 GeV for  $U\bar{D}\bar{D}$  couplings, also derived in ABDALLAH 04M. For indirect decays via  $LL\bar{E}$  the limit improves to 90 GeV if the constraint from the neutralino is used and remains at 87 GeV if it is not used. For indirect decays via  $U\bar{D}\bar{D}$  couplings it degrades to 85 GeV when the neutralino constraint is not used. Supersedes the result of ABREU 00U.
  - <sup>10</sup> HEISTER 03G searches for the production of smuons in the case of  $\cancel{R}$  prompt decays with  $LL\bar{E}$ ,  $LQ\bar{D}$  or  $U\bar{D}\bar{D}$  couplings at  $\sqrt{s} = 189$ –209 GeV. The search is performed for direct and indirect decays, assuming one coupling at a time to be non-zero. The limit holds for direct decays mediated by  $\cancel{R}$   $LQ\bar{D}$  couplings and improves to 90 GeV for indirect decays (for  $\Delta m > 10$  GeV). Limits are also given for  $LL\bar{E}$  direct ( $m_{\tilde{\mu}_R} > 87$  GeV) and indirect decays ( $m_{\tilde{\mu}_R} > 96$  GeV for  $m(\tilde{\chi}_1^0) > 23$  GeV from BARATE 98s) and for  $U\bar{D}\bar{D}$  indirect decays ( $m_{\tilde{\mu}_R} > 85$  GeV for  $\Delta m > 10$  GeV). Supersedes the results from BARATE 01B.

- 11 ABAZOV 02H looked in  $94 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8 \text{ TeV}$  for events with at least 2 muons and 2 jets for  $s$ -channel production of  $\tilde{\mu}$  or  $\tilde{\nu}$  and subsequent decay via  $\tilde{R}$  couplings  $LQ\bar{D}$ . A scan over the MSUGRA parameters is performed to exclude regions of the  $(m_0, m_{1/2})$  plane, examples being shown in Fig. 2.
- 12 ACHARD 02 searches for the production of smuons in the case of  $\tilde{R}$  prompt decays with  $LL\bar{E}$  or  $UDD$  couplings at  $\sqrt{s}=189\text{--}208 \text{ GeV}$ . The search is performed for direct and indirect decays, assuming one coupling at the time to be nonzero. The limit holds for direct decays via  $LL\bar{E}$  couplings. Stronger limits are reached for  $LL\bar{E}$  indirect (87 GeV) and for  $UDD$  direct or indirect (86 GeV) decays.
- 13 BARATE 01 looked for acoplanar dimuon +  $\cancel{E}_T$  final states at 189 to 202 GeV. The limit assumes 100% branching ratio for  $\tilde{\mu} \rightarrow \mu\tilde{\chi}_1^0$ . See their Fig. 1 for the dependence of the limit on  $\Delta m$ . These limits include and update the results of BARATE 99Q.
- 14 ABBIENDI 00J looked for acoplanar dimuon +  $\cancel{E}_T$  final states at  $\sqrt{s}= 161\text{--}183 \text{ GeV}$ . The limit assumes  $B(\tilde{\mu} \rightarrow \mu\tilde{\chi}_1^0)=1$ . Using decay branching ratios derived from the MSSM, a lower limit of 65 GeV is obtained for  $\mu < -100 \text{ GeV}$  and  $\tan\beta=1.5$ . See their Figs. 10 and 13 for the dependence of the limit on the branching ratio and on  $\Delta m$ .
- 15 ABREU 00V use data from  $\sqrt{s}= 130\text{--}189 \text{ GeV}$  to search for tracks with large impact parameter or visible decay vertices. Limits are obtained as function of  $m_{\tilde{G}}$ , after combining these results with the search for slepton pair production in the SUGRA framework from ABREU 01 to cover prompt decays and on stable particle searches from ABREU 00Q. For limits at different  $m_{\tilde{G}}$ , see their Fig. 12.
- 16 BARATE 98K looked for  $\mu^+\mu^-\gamma\gamma + \cancel{E}$  final states at  $\sqrt{s}= 161\text{--}184 \text{ GeV}$ . See Fig. 4 for limits on the  $(m_{\tilde{\mu}_R}, m_{\tilde{\chi}_1^0})$  plane and for the effect of cascade decays.

## $\tilde{\tau}$ (Stau) MASS LIMIT

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
>85.2		1 ABBIENDI 04	OPAL	$\Delta m > 6 \text{ GeV}, \theta_\tau=\pi/2,  \mu  > 100 \text{ GeV}, \tan\beta=1.5$
>78.3		2 ACHARD 04	L3	$\Delta m > 15 \text{ GeV}, \theta_\tau=\pi/2,  \mu  > 200 \text{ GeV}, \tan\beta \geq 2$
<b>&gt;81.9</b>	95	3 ABDALLAH 03M	DLPH	$\Delta m > 15 \text{ GeV}, \text{all } \theta_\tau$
none $m_\tau - 26.3$	95	3 ABDALLAH 03M	DLPH	$\Delta m > m_\tau, \text{all } \theta_\tau$
>79	95	4 HEISTER 02E	ALEP	$\Delta m > 15 \text{ GeV}, \theta_\tau=\pi/2$
>76	95	4 HEISTER 02E	ALEP	$\Delta m > 15 \text{ GeV}, \theta_\tau=0.91$
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
		5 AAD 12AF	ATLS	$2\tau + \text{jets} + \cancel{E}_T, \text{GMSB}$
		6 AAD 12AG	ATLS	$\geq 1\tau_h + \text{jets} + \cancel{E}_T, \text{GMSB}$
		7 AAD 12CM	ATLS	$\geq 1\tau + \text{jets} + \cancel{E}_T, \text{GMSB}$
>87.4	95	8 ABBIENDI 06B	OPAL	$\tilde{\tau}_R \rightarrow \tau\tilde{G}, \text{all } \tau(\tilde{\tau}_R)$
>74	95	9 ABBIENDI 04F	OPAL	$\tilde{R}, \tilde{\tau}_L$
>68	95	10,11 ABDALLAH 04H	DLPH	AMSB, $\mu > 0$
>90	95	12 ABDALLAH 04M	DLPH	$\tilde{R}, \tilde{\tau}_R, \text{indirect}, \Delta m > 5 \text{ GeV}$
>82.5		13 ABDALLAH 03D	DLPH	$\tilde{\tau}_R \rightarrow \tau\tilde{G}, \text{all } \tau(\tilde{\tau}_R)$
>70	95	14 HEISTER 03G	ALEP	$\tilde{\tau}_R, \tilde{R} \text{ decay}$
>61	95	15 ACHARD 02	L3	$\tilde{\tau}_R, \tilde{R} \text{ decays}$
>77	95	16 HEISTER 02R	ALEP	$\tau_1, \text{any lifetime}$
>70	95	17 BARATE 01	ALEP	$\Delta m > 10 \text{ GeV}, \theta_\tau=\pi/2$
>68	95	17 BARATE 01	ALEP	$\Delta m > 10 \text{ GeV}, \theta_\tau=0.91$

>64	95	18	ABBIENDI	00J	OPAL	$\Delta m > 10$ GeV, $\tilde{\tau}_R^+ \tilde{\tau}_R^-$
>84	95	19	ABREU	00V	DLPH	$\tilde{\ell}_R \tilde{\ell}_R$ ( $\tilde{\ell}_R \rightarrow \ell \tilde{G}$ ), $m_{\tilde{G}} > 9$
>73	95	20	ABREU	00V	DLPH	$\tilde{\tau}_1 \tilde{\tau}_1$ ( $\tilde{\tau}_1 \rightarrow \tau \tilde{G}$ ), all $\tau(\tilde{\tau}_1)$
>52		21	BARATE	98K	ALEP	Any $\Delta m, \theta_{\tau} = \pi/2, \tilde{\tau}_R \rightarrow \tau \gamma \tilde{G}$

<sup>1</sup> ABBIENDI 04 search for  $\tilde{\tau}\tilde{\tau}$  production in acoplanar di-tau final states in the 183–208 GeV data. See Fig. 15 for the dependence of the limits on  $m_{\tilde{\chi}_1^0}$  and for the limit

at  $\tan\beta=35$ . Under the assumption of 100% branching ratio for  $\tilde{\tau}_R \rightarrow \tau \tilde{\chi}_1^0$ , the limit improves to 89.8 GeV for  $\Delta m > 8$  GeV. See Fig. 12 for the dependence of the limits on  $m_{\tilde{\chi}_1^0}$  at several values of the branching ratio and for their dependence on  $\theta_{\tau}$ . This limit supersedes ABBIENDI 00G.

<sup>2</sup> ACHARD 04 search for  $\tilde{\tau}\tilde{\tau}$  production in acoplanar di-tau final states in the 192–209 GeV data. Limits on  $m_{\tilde{\tau}_R}$  are derived from a scan over the MSSM parameter space with universal GUT scale gaugino and scalar masses  $m_{1/2}$  and  $m_0$ ,  $1 \leq \tan\beta \leq 60$  and  $-2 \leq \mu \leq 2$  TeV. See Fig. 4 for the dependence of the limits on  $m_{\tilde{\chi}_1^0}$ .

<sup>3</sup> ABDALLAH 03M looked for acoplanar ditau +  $\cancel{E}_T$  final states at  $\sqrt{s} = 130$ –208 GeV. A dedicated search was made for low mass  $\tilde{\tau}$ s decoupling from the  $Z^0$ . The limit assumes  $B(\tilde{\tau} \rightarrow \tau \tilde{\chi}_1^0) = 100\%$ . See Fig. 20 for limits on the  $(m_{\tilde{\tau}}, m_{\tilde{\chi}_1^0})$  plane and as function

of the  $\tilde{\chi}_1^0$  mass and of the branching ratio. The limit in the low-mass region improves to 29.6 and 31.1 GeV for  $\tilde{\tau}_R$  and  $\tilde{\tau}_L$ , respectively, at  $\Delta m > m_{\tilde{\tau}}$ . The limit in the high-mass region improves to 84.7 GeV for  $\tilde{\tau}_R$  and  $\Delta m > 15$  GeV. These limits include and update the results of ABREU 01.

<sup>4</sup> HEISTER 02E looked for acoplanar ditau +  $\cancel{E}_T$  final states from  $e^+e^-$  interactions between 183 and 209 GeV. The mass limit assumes  $B(\tilde{\tau} \rightarrow \tau \tilde{\chi}_1^0) = 1$ . See their Fig. 4 for the dependence of the limit on  $\Delta m$ . These limits include and update the results of BARATE 01.

<sup>5</sup> AAD 12AF searched in  $2 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for events with two tau leptons, jets and large  $\cancel{E}_T$  in a GMSB framework. No significant excess above the expected background was found and an upper limit on the visible cross section for new phenomena is set. A 95% C.L. lower limit of 32 TeV on the mGMSB breaking scale  $\Lambda$  is set for  $M_{mess} = 250$  TeV,  $N_S = 3$ ,  $\mu > 0$  and  $C_{grav} = 1$ , independent of  $\tan\beta$ .

<sup>6</sup> AAD 12AG searched in  $2.05 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for events with at least one hadronically decaying tau lepton, jets, and large  $\cancel{E}_T$  in a GMSB framework. No significant excess above the expected background was found and an upper limit on the visible cross section for new phenomena is set. A 95% C.L. lower limit of 30 TeV on the mGMSB breaking scale  $\Lambda$  is set for  $M_{mess} = 250$  TeV,  $N_S = 3$ ,  $\mu > 0$  and  $C_{grav} = 1$ , independent of  $\tan\beta$ . For large values of  $\tan\beta$ , the limit on  $\Lambda$  increases to 43 TeV.

<sup>7</sup> AAD 12CM searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s}=7$  TeV for events with at least one tau lepton, zero or one additional light lepton ( $e/\mu$ ) jets, and large  $\cancel{E}_T$  in a GMSB framework. No significant excess above the expected background was found and an upper limit on the visible cross section for new phenomena is set. A 95% C. L. lower limit of 54 TeV on the mGMSB breaking scale  $\Lambda$  is set for  $M_{mess} = 250$  TeV,  $N_S = 3$ ,  $\mu > 0$  and  $C_{grav} = 1$ , for  $\tan\beta > 20$ . Here the  $\tilde{\tau}_1$  is the NLSP.

<sup>8</sup> ABBIENDI 06B use  $600 \text{ pb}^{-1}$  of data from  $\sqrt{s} = 189$ –209 GeV. They look for events from pair-produced staus in a GMSB scenario with  $\tilde{\tau}$  NLSP including prompt  $\tilde{\tau}$  decays to ditau +  $\cancel{E}$  final states, large impact parameters, kinked tracks and heavy stable charged particles. Limits on the cross-section are computed as a function of  $m(\tilde{\tau})$  and the lifetime, see their Fig. 7. The limit is compared to the  $\sigma \cdot BR^2$  from a scan over the GMSB parameter space.

<sup>9</sup> ABBIENDI 04F use data from  $\sqrt{s} = 189$ –209 GeV. They derive limits on sparticle masses under the assumption of  $\cancel{R}$  with  $LL\bar{E}$  or  $LQ\bar{D}$  couplings. The results are valid for  $\tan\beta$

- $= 1.5$ ,  $\mu = -200$  GeV, with, in addition,  $\Delta m > 5$  GeV for indirect decays via  $LQ\bar{D}$ . The limit quoted applies to direct decays with  $LL\bar{E}$  couplings and improves to 75 GeV for  $LQ\bar{D}$  couplings. The limit on the  $\tilde{\tau}_R$  mass for indirect decays is 92 GeV for  $LL\bar{E}$  couplings at  $m_{\tilde{\chi}_1^0} = 10$  GeV and no exclusion is obtained for  $LQ\bar{D}$  couplings. Supersedes the results of ABBIENDI 00.
- <sup>10</sup> ABDALLAH 04H use data from LEP 1 and  $\sqrt{s} = 192\text{--}208$  GeV. They re-use results or re-analyze the data from ABDALLAH 03M to put limits on the parameter space of anomaly-mediated supersymmetry breaking (AMSB), which is scanned in the region  $1 < m_{3/2} < 50$  TeV,  $0 < m_0 < 1000$  GeV,  $1.5 < \tan\beta < 35$ , both signs of  $\mu$ . The constraints are obtained from the searches for mass degenerate chargino and neutralino, for SM-like and invisible Higgs, for leptonically decaying charginos and from the limit on non-SM  $Z$  width of 3.2 MeV. The limit is for  $m_t = 174.3$  GeV (see Table 2 for other  $m_t$  values).
- <sup>11</sup> The limit improves to 75 GeV for  $\mu < 0$ .
- <sup>12</sup> ABDALLAH 04M use data from  $\sqrt{s} = 192\text{--}208$  GeV to derive limits on sparticle masses under the assumption of  $\cancel{R}$  with  $LL\bar{E}$  couplings. The results are valid for  $\mu = -200$  GeV,  $\tan\beta = 1.5$ ,  $\Delta m > 5$  GeV and assuming a BR of 1 for the given decay. The limit quoted is for indirect decays using the neutralino constraint of 39.5 GeV, also derived in ABDALLAH 04M. For indirect decays via  $LL\bar{E}$  the limit decreases to 86 GeV if the constraint from the neutralino is not used. Supersedes the result of ABREU 00U.
- <sup>13</sup> ABDALLAH 03D use data from  $\sqrt{s} = 130\text{--}208$  GeV to search for tracks with large impact parameter or visible decay vertices and for heavy charged stable particles. Limits are obtained as function of  $m(\tilde{G})$ , after combining these results with the search for slepton pair production in the SUGRA framework from ABDALLAH 03M to cover prompt decays. The above limit is reached for the stau decaying promptly,  $m(\tilde{G}) < 6$  eV, and is computed for stau mixing yielding the minimal cross section. Stronger limits are obtained for longer lifetimes, See their Fig. 9. Supersedes the results of ABREU 01G.
- <sup>14</sup> HEISTER 03G searches for the production of stau in the case of  $\cancel{R}$  prompt decays with  $LL\bar{E}$ ,  $LQ\bar{D}$  or  $UDD$  couplings at  $\sqrt{s} = 189\text{--}209$  GeV. The search is performed for direct and indirect decays, assuming one coupling at a time to be non-zero. The limit holds for indirect decays mediated by  $\cancel{R}$   $UDD$  couplings with  $\Delta m > 10$  GeV. Limits are also given for  $LL\bar{E}$  direct ( $m_{\tilde{\tau}_R} > 87$  GeV) and indirect decays ( $m_{\tilde{\tau}_R} > 95$  GeV for  $m(\tilde{\chi}_1^0) > 23$  GeV from BARATE 98S) and for  $LQ\bar{D}$  indirect decays ( $m_{\tilde{\tau}_R} > 76$  GeV). Supersedes the results from BARATE 01B.
- <sup>15</sup> ACHARD 02 searches for the production of staus in the case of  $\cancel{R}$  prompt decays with  $LL\bar{E}$  or  $UDD$  couplings at  $\sqrt{s}=189\text{--}208$  GeV. The search is performed for direct and indirect decays, assuming one coupling at the time to be nonzero. The limit holds for direct decays via  $LL\bar{E}$  couplings. Stronger limits are reached for  $LL\bar{E}$  indirect (86 GeV) and for  $UDD$  direct or indirect (75 GeV) decays.
- <sup>16</sup> HEISTER 02R search for signals of GMSB in the 189–209 GeV data. For the  $\tilde{\chi}_1^0$  NLSP scenario, they looked for topologies consisting of  $\gamma\gamma\cancel{E}$  or a single  $\gamma$  not pointing to the interaction vertex. For the  $\tilde{\ell}$  NLSP case, the topologies consist of  $\ell\ell\cancel{E}$ , including leptons with large impact parameters, kinks, or stable particles. Limits are derived from a scan over the GMSB parameters (see their Table 5 for the ranges). The limit remains valid whichever is the NLSP. The absolute mass bound on the  $\tilde{\chi}_1^0$  for any lifetime includes indirect limits from the slepton search HEISTER 02E performed within the MSUGRA framework. A bound for any NLSP and any lifetime of 77 GeV has also been derived by using the constraints from the neutral Higgs search in HEISTER 02. In the co-NLSP scenario, limits  $m_{\tilde{e}_R} > 83$  GeV (neglecting  $t$ -channel exchange) and  $m_{\tilde{\mu}_R} > 88$  GeV are obtained independent of the lifetime. Supersedes the results from BARATE 00G.
- <sup>17</sup> BARATE 01 looked for acoplanar ditau +  $\cancel{E}_T$  final states at 189 to 202 GeV. A slight excess (with 1.2% probability) of events is observed relative to the expected SM background. The limit assumes 100% branching ratio for  $\tilde{\tau} \rightarrow \tau\tilde{\chi}_1^0$ . See their Fig. 1 for the dependence of the limit on  $\Delta m$ . These limits include and update the results of BARATE 99Q.

- <sup>18</sup> ABBIENDI 00J looked for acoplanar ditau +  $\cancel{E}_T$  final states at  $\sqrt{s}=161\text{--}183$  GeV. The limit assumes  $B(\tilde{\tau} \rightarrow \tau \tilde{\chi}_1^0)=1$ . Using decay branching ratios derived from the MSSM, a lower limit of 60 GeV at  $\Delta m > 9$  GeV is obtained for  $\mu < -100$  GeV and  $\tan\beta=1.5$ . See their Figs. 11 and 14 for the dependence of the limit on the branching ratio and on  $\Delta m$ .
- <sup>19</sup> ABREU 00V use data from  $\sqrt{s}=130\text{--}189$  GeV to search for tracks with large impact parameter or visible decay vertices. Limits are obtained as function of  $m_{\tilde{G}}$ , after combining these results with the search for slepton pair production in the SUGRA framework from ABREU 01 to cover prompt decays and on stable particle searches from ABREU 00Q. The above limit assumes the degeneracy of stau and smuon. For limits at different  $m_{\tilde{G}}$ , see their Fig. 12.
- <sup>20</sup> ABREU 00V use data from  $\sqrt{s}=130\text{--}189$  GeV to search for tracks with large impact parameter or visible decay vertices. Limits are obtained as function of  $m_{\tilde{G}}$ , after combining these results with the search for slepton pair production in the SUGRA framework from ABREU 01 to cover prompt decays and on stable particle searches from ABREU 00Q. The above limit is reached for the stau mixing yielding the minimal cross section and decaying promptly. Stronger limits are obtained for longer lifetimes or for  $\tilde{\tau}_R$ ; see their Fig. 11. For  $10 \leq m_{\tilde{G}} \leq 310$  eV, the whole range  $2 \leq m_{\tilde{\tau}_1} \leq 80$  GeV is excluded. Supersedes the results of ABREU 99C and ABREU 99F.
- <sup>21</sup> BARATE 98K looked for  $\tau^+ \tau^- \gamma \gamma + \cancel{E}$  final states at  $\sqrt{s}=161\text{--}184$  GeV. See Fig. 4 for limits on the  $(m_{\tilde{\tau}_R}, m_{\tilde{\chi}_1^0})$  plane and for the effect of cascade decays.

### Degenerate Charged Sleptons

Unless stated otherwise in the comment lines or in the footnotes, the following limits assume 3 families of degenerate charged sleptons.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
>93	95	<sup>1</sup> BARATE	01	ALEP $\Delta m > 10$ GeV, $\tilde{\ell}_R^+ \tilde{\ell}_R^-$
>70	95	<sup>1</sup> BARATE	01	ALEP all $\Delta m$ , $\tilde{\ell}_R^+ \tilde{\ell}_R^-$
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
>91.9	95	<sup>2</sup> ABBIENDI	06B	OPAL $\tilde{\ell}_R \rightarrow \ell \tilde{G}$ , all $\ell(\tilde{\ell}_R)$
>88		<sup>3</sup> ABDALLAH	03D	DLPH $\tilde{\ell}_R \rightarrow \ell \tilde{G}$ , all $\ell(\tilde{\ell}_R)$
>82.7	95	<sup>4</sup> ACHARD	02	L3 $\tilde{\ell}_R, \cancel{R}$ decays, MSUGRA
>83	95	<sup>5</sup> ABBIENDI	01	OPAL $e^+ e^- \rightarrow \tilde{\ell}_1^+ \tilde{\ell}_1^-$ , GMSB, $\tan\beta=2$
		<sup>6</sup> ABREU	01	DLPH $\tilde{\ell} \rightarrow \ell \tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow \gamma \tilde{\chi}_1^0$ , $\ell=e, \mu$
>68.8	95	<sup>7</sup> ACCIARRI	01	L3 $\tilde{\ell}_R, \cancel{R}, 0.7 \leq \tan\beta \leq 40$
>84	95	<sup>8,9</sup> ABREU	00V	DLPH $\tilde{\ell}_R \tilde{\ell}_R (\tilde{\ell}_R \rightarrow \ell \tilde{G}), m_{\tilde{G}} > 9$ eV

<sup>1</sup> BARATE 01 looked for acoplanar dilepton +  $\cancel{E}_T$  and single electron (for  $\tilde{e}_R \tilde{e}_L$ ) final states at 189 to 202 GeV. The limit assumes  $\mu=-200$  GeV and  $\tan\beta=2$  for the production cross section and decay branching ratios, evaluated within the MSSM, and zero efficiency for decays other than  $\tilde{\ell} \rightarrow \ell \tilde{\chi}_1^0$ . The slepton masses are determined from the GUT relations without stau mixing. See their Fig. 1 for the dependence of the limit on  $\Delta m$ .

<sup>2</sup> ABBIENDI 06B use  $600 \text{ pb}^{-1}$  of data from  $\sqrt{s}=189\text{--}209$  GeV. They look for events from pair-produced staus in a GMSB scenario with  $\tilde{\ell}$  co-NLSP including prompt  $\tilde{\ell}$  decays to dileptons +  $\cancel{E}$  final states, large impact parameters, kinked tracks and heavy stable charged particles. Limits on the cross-section are computed as a function of  $m(\tilde{\ell})$  and the lifetime, see their Fig. 7. The limit is compared to the  $\sigma \cdot BR^2$  from a scan over the GMSB parameter space. The highest mass limit is reached for  $\tilde{\mu}_R$ , from which the quoted mass limit is derived by subtracting  $m_\tau$ .

- <sup>3</sup> ABDALLAH 03D use data from  $\sqrt{s} = 130\text{--}208$  GeV to search for tracks with large impact parameter or visible decay vertices and for heavy charged stable particles. Limits are obtained as function of  $m(\tilde{G})$ , after combining these results with the search for slepton pair production in the SUGRA framework from ABDALLAH 03M to cover prompt decays. The above limit is reached for prompt decays and assumes the degeneracy of the sleptons. For limits at different  $m(\tilde{G})$ , see their Fig. 9. Supersedes the results of ABREU 01G.
- <sup>4</sup> ACHARD 02 searches for the production of sparticles in the case of  $\tilde{R}$  prompt decays with  $LL\bar{E}$  or  $UDD$  couplings at  $\sqrt{s}=189\text{--}208$  GeV. The search is performed for direct and indirect decays, assuming one coupling at the time to be nonzero. The MSUGRA limit results from a scan over the MSSM parameter space with the assumption of gaugino and scalar mass unification at the GUT scale and no mixing in the slepton sector, imposing simultaneously the exclusions from neutralino, chargino, sleptons, and squarks analyses. The limit holds for  $LL\bar{E}$  couplings and increases to 88.7 GeV for  $UDD$  couplings. For L3 limits from  $LQ\bar{D}$  couplings, see ACCIARRI 01.
- <sup>5</sup> ABBIENDI 01 looked for final states with  $\gamma\gamma\cancel{E}$ ,  $ll\cancel{E}$ , with possibly additional activity and four leptons +  $\cancel{E}$  to search for prompt decays of  $\tilde{\chi}_1^0$  or  $\tilde{\ell}_1$  in GMSB. They derive limits in the plane  $(m_{\tilde{\chi}_1^0}, m_{\tilde{\tau}_1})$ , see Fig. 6, allowing either the  $\tilde{\chi}_1^0$  or a  $\tilde{\ell}_1$  to be the NLSP. Two scenarios are considered:  $\tan\beta=2$  with the 3 sleptons degenerate in mass and  $\tan\beta=20$  where the  $\tilde{\tau}_1$  is lighter than the other sleptons. Data taken at  $\sqrt{s}=189$  GeV. For  $\tan\beta=20$ , the obtained limits are  $m_{\tilde{\tau}_1} > 69$  GeV and  $m_{\tilde{e}_1, \tilde{\mu}_1} > 88$  GeV.
- <sup>6</sup> ABREU 01 looked for acoplanar dilepton + diphoton +  $\cancel{E}$  final states from  $\tilde{\ell}$  cascade decays at  $\sqrt{s}=130\text{--}189$  GeV. See Fig. 9 for limits on the  $(\mu, M_2)$  plane for  $m_{\tilde{\ell}}=80$  GeV,  $\tan\beta=1.0$ , and assuming degeneracy of  $\tilde{\mu}$  and  $\tilde{e}$ .
- <sup>7</sup> ACCIARRI 01 searches for multi-lepton and/or multi-jet final states from  $\tilde{R}$  prompt decays with  $LL\bar{E}$ ,  $LQ\bar{D}$ , or  $UDD$  couplings at  $\sqrt{s}=189$  GeV. The search is performed for direct and indirect decays of neutralinos, charginos, and scalar leptons, with the  $\tilde{\chi}_1^0$  or a  $\tilde{\ell}$  as LSP and assuming one coupling to be nonzero at a time. Mass limits are derived using simultaneously the constraints from the neutralino, chargino, and slepton analyses; and the  $Z^0$  width measurements from ACCIARRI 00C in a scan of the parameter space assuming MSUGRA with gaugino and scalar mass universality. Updates and supersedes the results from ACCIARRI 99I.
- <sup>8</sup> ABREU 00V use data from  $\sqrt{s}=130\text{--}189$  GeV to search for tracks with large impact parameter or visible decay vertices. Limits are obtained as function of  $m_{\tilde{G}}$ , after combining these results with the search for slepton pair production in the SUGRA framework from ABREU 01 to cover prompt decays and on stable particle searches from ABREU 00Q. For limits at different  $m_{\tilde{G}}$ , see their Fig. 12.
- <sup>9</sup> The above limit assumes the degeneracy of stau and smuon.

## Long-lived $\tilde{\ell}$ (Slepton) MASS LIMIT

Limits on scalar leptons which leave detector before decaying. Limits from  $Z$  decays are independent of lepton flavor. Limits from continuum  $e^+e^-$  annihilation are also independent of flavor for smuons and staus. Selectron limits from  $e^+e^-$  collisions in the continuum depend on MSSM parameters because of the additional neutralino exchange contribution.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
> 98	95	<sup>1</sup> ABBIENDI	03L OPAL	$\tilde{\mu}_R, \tilde{\tau}_R$
<b>none 2–87.5</b>	95	<sup>2</sup> ABREU	00Q DLPH	$\tilde{\mu}_R, \tilde{\tau}_R$
> 81.2	95	<sup>3</sup> ACCIARRI	99H L3	$\tilde{\mu}_R, \tilde{\tau}_R$
> 81	95	<sup>4</sup> BARATE	98K ALEP	$\tilde{\mu}_R, \tilde{\tau}_R$

• • • We do not use the following data for averages, fits, limits, etc. • • •

>300	95	<sup>5</sup> AAD	13AA ATLS	long-lived $\tilde{\tau}$ , GMSB, $\tan\beta = 5\text{--}20$
		<sup>6</sup> ABAZOV	13B D0	long-lived $\tilde{\tau}$ , $100 < m_{\tilde{\tau}} < 300$ GeV
>339	95	<sup>7,8</sup> CHATRCHYAN	13AB CMS	long-lived $\tilde{\tau}$ , direct $\tilde{\tau}_1$ pair prod., minimal GMSB, SPS line 7
>500	95	<sup>7,9</sup> CHATRCHYAN	13AB CMS	long-lived $\tilde{\tau}$ , $\tilde{\tau}_1$ from direct pair prod. and from decay of heavier SUSY particles, minimal GMSB, SPS line 7
>314	95	<sup>10</sup> CHATRCHYAN	12L CMS	long-lived $\tilde{\tau}$ , $\tilde{\tau}_1$ from decay of heavier SUSY particles, minimal GMSB, SPS line 7
>136	95	<sup>11</sup> AAD	11P ATLS	stable $\tilde{\tau}$ , GMSB scenario, $\tan\beta=5$

<sup>1</sup> ABBIENDI 03L used  $e^+e^-$  data at  $\sqrt{s} = 130\text{--}209$  GeV to select events with two high momentum tracks with anomalous  $dE/dx$ . The excluded cross section is compared to the theoretical expectation as a function of the heavy particle mass in their Fig. 3. The limit improves to 98.5 GeV for  $\tilde{\mu}_L$  and  $\tilde{\tau}_L$ . The bounds are valid for colorless spin 0 particles with lifetimes longer than  $10^{-6}$  s. Supersedes the results from ACKERSTAFF 98P.

<sup>2</sup> ABREU 00Q searches for the production of pairs of heavy, charged stable particles in  $e^+e^-$  annihilation at  $\sqrt{s} = 130\text{--}189$  GeV. The upper bound improves to 88 GeV for  $\tilde{\mu}_L$ ,  $\tilde{\tau}_L$ . These limits include and update the results of ABREU 98P.

<sup>3</sup> ACCIARRI 99H searched for production of pairs of back-to-back heavy charged particles at  $\sqrt{s}=130\text{--}183$  GeV. The upper bound improves to 82.2 GeV for  $\tilde{\mu}_L$ ,  $\tilde{\tau}_L$ .

<sup>4</sup> The BARATE 98K mass limit improves to 82 GeV for  $\tilde{\mu}_L, \tilde{\tau}_L$ . Data collected at  $\sqrt{s}=161\text{--}184$  GeV.

<sup>5</sup> AAD 13AA searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for events containing long-lived massive particles in a GMSB framework. No significant excess above the expected background was found. A 95% C.L. lower limit of 300 GeV is placed on long-lived  $\tilde{\tau}$ 's in the GMSB model with  $M_{mess} = 250$  TeV,  $N_G = 3$ ,  $\mu > 0$ , for  $\tan\beta = 5\text{--}20$ . The lower limit on the GMSB breaking scale  $\Lambda$  was found to be 99–110 TeV, for  $\tan\beta$  values between 5 and 40, see Fig. 4 (top). Also, directly produced long-lived sleptons, or sleptons decaying to long-lived ones, are excluded at 95% C.L. up to a  $\tilde{\tau}$  mass of 278 GeV for models with slepton splittings smaller than 50 GeV.

<sup>6</sup> ABAZOV 13B looked in  $6.3 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV for charged massive long-lived particles in events with muon-like particles that have both speed and ionization energy loss inconsistent with muons produced in beam collisions. In the absence of an excess, limits are set at 95% C.L. on the production cross section of stau leptons in the mass range 100–300 GeV, see their Table 20 and Fig. 23.

<sup>7</sup> CHATRCHYAN 13AB looked in  $5.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV and in  $18.8 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8$  TeV for events with heavy stable particles, identified by their anomalous  $dE/dx$  in the tracker or additionally requiring that it be identified as muon in the muon chambers, from pair production of  $\tilde{\tau}_1$ 's. No evidence for an excess over the expected background is observed. Supersedes CHATRCHYAN 12L.

<sup>8</sup> CHATRCHYAN 13AB limits are derived for pair production of  $\tilde{\tau}_1$  as a function of mass in minimal GMSB scenarios along the Snowmass Points and Slopes (SPS) line 7 (see Fig. 8 and Table 7). The limit given here is valid for direct pair  $\tilde{\tau}_1$  production.

<sup>9</sup> CHATRCHYAN 13AB limits are derived for the production of  $\tilde{\tau}_1$  as a function of mass in minimal GMSB scenarios along the Snowmass Points and Slopes (SPS) line 7 (see Fig. 8 and Table 7). The limit given here is valid for the production of  $\tilde{\tau}_1$  from both direct pair production and from the decay of heavier supersymmetric particles.

- <sup>10</sup> CHATRCHYAN 12L looked in  $5.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with heavy stable particles, identified by their anomalous  $dE/dx$  in the tracker or additionally requiring that it be identified as muon in the muon chambers, from pair production of  $\tilde{\tau}_1$ 's. No evidence for an excess over the expected background is observed. Limits are derived for the production of  $\tilde{\tau}_1$  as a function of mass in minimal GMSB scenarios along the Snowmass Points and Slopes (SPS) line 7 (see Fig. 3). The limit given here is valid for the production of  $\tilde{\tau}_1$  in the decay of heavier supersymmetric particles.
- <sup>11</sup> AAD 11P looked in  $37 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with two heavy stable particles, reconstructed in the Inner tracker and the Muon System and identified by their time of flight in the Muon System. No evidence for an excess over the SM expectation is observed. Limits on the mass are derived, see Fig. 3, for  $\tilde{\tau}$  in a GMSB scenario and for sleptons produced by electroweak processes only, in which case the limit degrades to 110 GeV.

## $\tilde{q}$ (Squark) MASS LIMIT

For  $m_{\tilde{q}} > 60\text{--}70 \text{ GeV}$ , it is expected that squarks would undergo a cascade decay via a number of neutralinos and/or charginos rather than undergo a direct decay to photinos as assumed by some papers. Limits obtained when direct decay is assumed are usually higher than limits when cascade decays are included.

Limits from  $e^+e^-$  collisions depend on the mixing angle of the lightest mass eigenstate  $\tilde{q}_1 = \tilde{q}_R \sin\theta_q + \tilde{q}_L \cos\theta_q$ . It is usually assumed that only the sbottom and stop squarks have non-trivial mixing angles (see the stop and sbottom sections). Here, unless otherwise noted, squarks are always taken to be either left/right degenerate, or purely of left or right type. Data from  $Z$  decays have set squark mass limits above 40 GeV, in the case of  $\tilde{q} \rightarrow q\tilde{\chi}_1^0$  decays if  $\Delta m = m_{\tilde{q}} - m_{\tilde{\chi}_1^0} \gtrsim 5 \text{ GeV}$ . For smaller values of  $\Delta m$ , current constraints on the invisible width of the  $Z$  ( $\Delta\Gamma_{\text{inv}} < 2.0 \text{ MeV}$ , LEP 00) exclude  $m_{\tilde{u}_{L,R}} < 44 \text{ GeV}$ ,  $m_{\tilde{d}_R} < 33 \text{ GeV}$ ,  $m_{\tilde{d}_L} < 44 \text{ GeV}$  and, assuming all squarks degenerate,  $m_{\tilde{q}} < 45 \text{ GeV}$ .

Limits made obsolete by the most recent analyses of  $e^+e^-$ ,  $p\bar{p}$ , and  $ep$  collisions can be found in previous Editions of this *Review*.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
<b>&gt;1110</b>	<b>(CL = 95%)</b>	<b>OUR EVALUATION</b>		
>1360	95	<sup>1</sup> AAD	13L ATLS	jets + $\cancel{E}_T$ , CMSSM, $m_{\tilde{g}} = m_{\tilde{q}}$
>1200	95	<sup>2</sup> AAD	13Q ATLS	$\gamma + b + \cancel{E}_T$ , higgsino-like neutralino, $m_{\tilde{\chi}_1^0} > 220 \text{ GeV}$ , GMSB
>1250	95	<sup>3</sup> CHATRCHYAN 13	CMS	$\ell^\pm \ell^\mp + \text{jets} + \cancel{E}_T$ , CMSSM
		<sup>4</sup> CHATRCHYAN 13G	CMS	$0, 1, 2, \geq 3 \text{ } b\text{-jets} + \cancel{E}_T$ , CMSSM, $m_{\tilde{q}} = m_{\tilde{g}}$
>1430	95	<sup>5</sup> CHATRCHYAN 13H	CMS	$2\gamma + \geq 4 \text{ jets} + \text{low } \cancel{E}_T$ , stealth SUSY model
<b>&gt; 750</b>	95	<sup>6</sup> CHATRCHYAN 13T	CMS	jets + $\cancel{E}_T$ , $\tilde{q} \rightarrow q\tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} = 0 \text{ GeV}$
> 820	95	<sup>7</sup> AAD	12AX ATLS	$\ell + \text{jets} + \cancel{E}_T$ , CMSSM, $m_{\tilde{q}} = m_{\tilde{g}}$
>1200	95	<sup>8</sup> AAD	12CJ ATLS	$\ell^\pm + \text{jets} + \cancel{E}_T$ , CMSSM, $m_{\tilde{q}} = m_{\tilde{g}}$
> 870	95	<sup>9</sup> AAD	12CP ATLS	$2\gamma + \cancel{E}_T$ , GMSB, bino NLSP, $m_{\tilde{\chi}_1^0} > 50 \text{ GeV}$
> 950	95	<sup>10</sup> AAD	12W ATLS	jets + $\cancel{E}_T$ , CMSSM, $m_{\tilde{q}} = m_{\tilde{g}}$
		<sup>11</sup> CHATRCHYAN 12	CMS	$e, \mu, \text{jets}$ , razor, CMSSM

> 760	95	12	CHATRCHYAN 12AE	CMS	jets + $\cancel{E}_T$ , $\tilde{q} \rightarrow q\tilde{\chi}_1^0$ , $m_{\tilde{\chi}_1^0} <$
					$200 \text{ GeV}$
					$\geq 3\ell^\pm, \cancel{R}$
>1110	95	14	CHATRCHYAN 12AL	CMS	jets + $\cancel{E}_T$ , CMSSM
>1180	95	14	CHATRCHYAN 12AT	CMS	jets + $\cancel{E}_T$ , CMSSM, $m_{\tilde{q}}=m_{\tilde{g}}$
> 690	95	15	AAD	11B ATLS	$\ell^\pm \ell^\pm + \cancel{E}_T$ , $m_{\tilde{g}}=m_{\tilde{q}}+10\text{GeV}$ , $m_{\tilde{\chi}_1^0}=100\text{GeV}$ , $\tan\beta=4$
> 550	95	15	AAD	11B ATLS	$\ell^+ \ell^- + \cancel{E}_T$ , $m_{\tilde{g}}=m_{\tilde{q}}+10\text{GeV}$ , $m_{\tilde{\chi}_1^0}=100\text{GeV}$ , $\tan\beta=4$
> 558	95	16	AAD	11C ATLS	$\ell^+ \ell^- + \text{jets} + \cancel{E}_T$ , $m_{\tilde{g}}=m_{\tilde{q}}+10\text{GeV}$ , $m_{\tilde{\chi}_1^0}=100\text{GeV}$ , $\tan\beta=4$
> 700	95	17	AAD	11G ATLS	$\ell + \text{jets} + \cancel{E}_T$ , $\tan\beta=3$ , $A_0=0$ , $\mu >$ $0$ , $m_{\tilde{g}}=m_{\tilde{q}}$
> 870	95	18	AAD	11N ATLS	jets + $\cancel{E}_T$ , degenerate $m_{\tilde{q}}$ of first two generations, $m_{\tilde{\chi}_1^0}=0$ , all other supersymmetric particles heavy, $m_{\tilde{q}}=m_{\tilde{g}}$
> 775	95	18	AAD	11N ATLS	jets + $\cancel{E}_T$ , CMSSM, $m_{\tilde{q}}=m_{\tilde{g}}$
>1100	95	19	CHATRCHYAN 11W	CMS	jets + $\cancel{E}_T$ , CMSSM
> 392	95	20	AALTONEN	09S CDF	jets + $\cancel{E}_T$ , $m_{\tilde{q}}=m_{\tilde{g}}$
> 379	95	21	ABAZOV	08G D0	jets + $\cancel{E}_T$ , $\tan\beta=3$ , $\mu < 0$ , $A_0=0$ , any $m_{\tilde{g}}$
> 99.5		22	ACHARD	04 L3	$\Delta m > 10 \text{ GeV}$ , $e^+ e^- \rightarrow$ $\tilde{q}_{L,R} \tilde{q}_{L,R}^*$
> 97		22	ACHARD	04 L3	$\Delta m > 10 \text{ GeV}$ , $e^+ e^- \rightarrow \tilde{q}_R \tilde{q}_R^*$
> 138	95	23	ABBOTT	01D D0	$\ell\ell + \text{jets} + \cancel{E}_T$ , $\tan\beta < 10$ , $m_0 <$ $300 \text{ GeV}$ , $\mu < 0$ , $A_0=0$
> 255	95	23	ABBOTT	01D D0	$\tan\beta=2$ , $m_{\tilde{g}}=m_{\tilde{q}}$ , $\mu < 0$ , $A_0=0$ , $\ell\ell + \text{jets} + \cancel{E}_T$
> 97	95	24	BARATE	01 ALEP	$e^+ e^- \rightarrow \tilde{q}\tilde{q}^*$ , $\Delta m > 6 \text{ GeV}$
> 224	95	25	ABE	96D CDF	$m_{\tilde{g}} \leq m_{\tilde{q}}$ ; with cascade decays, $\ell\ell + \text{jets} + \cancel{E}_T$
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●					
> 700	95	26	CHATRCHYAN 13AO	CMS	$\ell^\pm \ell^\mp + \text{jets} + \cancel{E}_T$ , CMSSM, $m_0 < 700 \text{ GeV}$
>1350	95	27	CHATRCHYAN 13AV	CMS	jets (+ leptons) + $\cancel{E}_T$ , CMSSM, $m_{\tilde{g}} = m_{\tilde{q}}$
> 800	95	28	CHATRCHYAN 13W	CMS	$\geq 1$ photons + jets + $\cancel{E}_T$ , GGM, wino-like NLSP, $m_{\tilde{\chi}_1^0}$ $= 375 \text{ GeV}$
>1000	95	28	CHATRCHYAN 13W	CMS	$\geq 2$ photons + jets + $\cancel{E}_T$ , GGM, bino-like NLSP, $m_{\tilde{\chi}_1^0}$ $= 375 \text{ GeV}$
> 340	95	29	DREINER	12A THEO	$m_{\tilde{q}} \sim m_{\tilde{\chi}_1^0}$

> 650	95	30 DREINER	12A THEO	$m_{\tilde{q}} = m_{\tilde{g}} \sim m_{\tilde{\chi}_1^0}$
		31 AAD	11AE ATLS	$\ell^\pm \ell^\pm$
		32 AAD	11AF ATLS	$\geq 6$ jets + $\cancel{E}_T$ , CMSSM
> 290	95	33 AARON	11 H1	$e^- p \rightarrow \tilde{d}_R, \cancel{R}, LQ\bar{D}, \lambda'=0.3$
> 275	95	33 AARON	11 H1	$e^+ p \rightarrow \tilde{u}_L, \cancel{R}, LQ\bar{D}, \lambda'=0.3$
> 330	95	34 AARON	11C H1	$\tilde{u}, \cancel{R}, LQ\bar{D}, \lambda'=0.3$
		35 CHATRCHYAN	11AC CMS	jets + $\cancel{E}_T$ , CMSSM
		36 CHATRCHYAN	11C CMS	$\tilde{q} \rightarrow X \tilde{\chi}_2^0 \rightarrow X \ell^+ \ell^- \tilde{\chi}_1^0$
		37 CHATRCHYAN	11G CMS	$\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$
		38 CHATRCHYAN	11Q CMS	$\ell$ + jets + $\cancel{E}_T$
> 830	95	39 CHATRCHYAN	11V CMS	GMSB scenario, $\bar{\ell}$ co-NLSP
		40 CHATRCHYAN	11V CMS	$\cancel{R}$
		41 KHACHATRY...	11I CMS	jets + $\cancel{E}_T$
		42 ABAZOV	09S D0	jets + $\tau$ + $\cancel{E}_T$ , $\tan\beta=15, \mu < 0,$ $A_0 = -2m_0$
> 490	95	43 SCHAEAL	07A ALEP	$\tilde{d}_R, \cancel{R}, \lambda=0.3$
> 544	95	43 SCHAEAL	07A ALEP	$\tilde{s}_R, \cancel{R}, \lambda=0.3$
> 273	95	44 CHEKANOV	05A ZEUS	$\tilde{q} \rightarrow \mu q, \cancel{R}, LQ\bar{D}, \lambda=0.3$
> 270	95	44 CHEKANOV	05A ZEUS	$\tilde{q} \rightarrow \tau q, \cancel{R}, LQ\bar{D}, \lambda=0.3$
> 275		45 AKTAS	04D H1	$e^\pm p \rightarrow \tilde{U}_L, \cancel{R}, LQ\bar{D}$
> 280		45 AKTAS	04D H1	$e^\pm p \rightarrow \tilde{D}_R, \cancel{R}, LQ\bar{D}$
		46 ADLOFF	03 H1	$e^\pm p \rightarrow \tilde{q}, \cancel{R}, LQ\bar{D}$
> 276	95	47 CHEKANOV	03B ZEUS	$\tilde{d} \rightarrow e^- u, \nu d, \cancel{R}, LQ\bar{D}, \lambda > 0.1$
> 260	95	47 CHEKANOV	03B ZEUS	$\tilde{u} \rightarrow e^+ d, \cancel{R}, LQ\bar{D}, \lambda > 0.1$
> 82.5	95	48 HEISTER	03G ALEP	$\tilde{u}_R, \cancel{R}$ decay
> 77	95	48 HEISTER	03G ALEP	$\tilde{d}_R, \cancel{R}$ decay
> 240	95	49 ABAZOV	02F D0	$\tilde{q}, \cancel{R} \lambda'_{2jk}$ indirect decays, $\tan\beta=2, \text{ any } m_{\tilde{g}}$
> 265	95	49 ABAZOV	02F D0	$\tilde{q}, \cancel{R} \lambda'_{2jk}$ indirect decays, $\tan\beta=2, m_{\tilde{q}}=m_{\tilde{g}}$
		50 ABAZOV	02G D0	$p\bar{p} \rightarrow \tilde{g}\tilde{g}, \tilde{g}\tilde{q}$
none 80–121	95	51 ABBIENDI	02 OPAL	$e\gamma \rightarrow \tilde{u}_L, \cancel{R}, LQ\bar{D}, \lambda=0.3$
none 80–158	95	51 ABBIENDI	02 OPAL	$e\gamma \rightarrow \tilde{d}_R, \cancel{R}, LQ\bar{D}, \lambda=0.3$
none 80–185	95	52 ABBIENDI	02B OPAL	$e\gamma \rightarrow \tilde{u}_L, \cancel{R}, LQ\bar{D}, \lambda=0.3$
none 80–196	95	52 ABBIENDI	02B OPAL	$e\gamma \rightarrow \tilde{d}_R, \cancel{R}, LQ\bar{D}, \lambda=0.3$
> 79	95	53 ACHARD	02 L3	$\tilde{u}_R, \cancel{R}$ decays
> 55	95	53 ACHARD	02 L3	$\tilde{d}_R, \cancel{R}$ decays
> 263	95	54 CHEKANOV	02 ZEUS	$\tilde{u}_L \rightarrow \mu q, \cancel{R}, LQ\bar{D}, \lambda=0.3$
> 258	95	54 CHEKANOV	02 ZEUS	$\tilde{u}_L \rightarrow \tau q, \cancel{R}, LQ\bar{D}, \lambda=0.3$
> 82	95	55 BARATE	01B ALEP	$\tilde{u}_R, \cancel{R}$ decays
> 68	95	55 BARATE	01B ALEP	$\tilde{d}_R, \cancel{R}$ decays
none 150–204	95	56 BREITWEG	01 ZEUS	$e^+ p \rightarrow \tilde{d}_R, \cancel{R}, LQ\bar{D}, \lambda=0.3$
> 200	95	57 ABBOTT	00C D0	$\tilde{u}_L, \cancel{R}, \lambda'_{2jk}$ decays
> 180	95	57 ABBOTT	00C D0	$\tilde{d}_R, \cancel{R}, \lambda'_{2jk}$ decays
> 390	95	58 ACCIARRI	00P L3	$e^+ e^- \rightarrow q\bar{q}, \cancel{R}, \lambda=0.3$
> 148	95	59 AFFOLDER	00K CDF	$\tilde{d}_L, \cancel{R} \lambda'_{ij3}$ decays

> 200	95	60 BARATE	00I	ALEP	$e^+e^- \rightarrow q\bar{q}, \cancel{R}, \lambda=0.3$
none 150–269	95	61 BREITWEG	00E	ZEUS	$e^+p \rightarrow \tilde{u}_L, \cancel{R}, LQ\bar{D}, \lambda=0.3$
> 240	95	62 ABBOTT	99	D0	$\tilde{q} \rightarrow \tilde{\chi}_2^0 X \rightarrow \tilde{\chi}_1^0 \gamma X, m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0} > 20 \text{ GeV}$
> 320	95	62 ABBOTT	99	D0	$\tilde{q} \rightarrow \tilde{\chi}_1^0 X \rightarrow \tilde{G} \gamma X$
> 243	95	63 ABBOTT	99K	D0	any $m_{\tilde{g}}, \cancel{R}, \tan\beta=2, \mu < 0$
> 250	95	64 ABBOTT	99L	D0	$\tan\beta=2, \mu < 0, A=0, \text{jets} + \cancel{E}_T$
> 200	95	65 ABE	99M	CDF	$p\bar{p} \rightarrow \tilde{q}\tilde{q}, \cancel{R}$
none 80–134	95	66 ABREU	99G	DLPH	$e\gamma \rightarrow \tilde{u}_L, \cancel{R}, LQ\bar{D}, \lambda=0.3$
none 80–161	95	66 ABREU	99G	DLPH	$e\gamma \rightarrow \tilde{d}_R, \cancel{R}, LQ\bar{D}, \lambda=0.3$
> 225	95	67 ABBOTT	98E	D0	$\tilde{u}_L, \cancel{R}, \lambda'_{1jk}$ decays
> 204	95	67 ABBOTT	98E	D0	$\tilde{d}_R, \cancel{R}, \lambda'_{1jk}$ decays
> 79	95	67 ABBOTT	98E	D0	$\tilde{d}_L, \cancel{R}, \lambda'_{ijk}$ decays
> 202	95	68 ABE	98S	CDF	$\tilde{u}_L, \cancel{R}, \lambda'_{2jk}$ decays
> 160	95	68 ABE	98S	CDF	$\tilde{d}_R, \cancel{R}, \lambda'_{2jk}$ decays
> 140	95	69 ACKERSTAFF	98V	OPAL	$e^+e^- \rightarrow q\bar{q}, \cancel{R}, \lambda=0.3$
> 77	95	70 BREITWEG	98	ZEUS	$m_{\tilde{q}}=m_{\tilde{e}}, m(\tilde{\chi}_1^0)=40 \text{ GeV}$
		71 DATTA	97	THEO	$\tilde{\nu}$ 's lighter than $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$
> 216	95	72 DERRICK	97	ZEUS	$ep \rightarrow \tilde{q}, \tilde{q} \rightarrow \mu j \text{ or } \tau j, \cancel{R}$
none 130–573	95	73 HEWETT	97	THEO	$q\tilde{g} \rightarrow \tilde{q}, \tilde{q} \rightarrow q\tilde{g}, \text{ with a light gluino}$
none 190–650	95	74 TEREKHOV	97	THEO	$qg \rightarrow \tilde{q}\tilde{g}, \tilde{q} \rightarrow q\tilde{g}, \text{ with a light gluino}$
> 63	95	75 AID	96C	H1	$m_{\tilde{q}}=m_{\tilde{e}}, m_{\tilde{\chi}_1^0}=35 \text{ GeV}$
none 330–400	95	76 TEREKHOV	96	THEO	$ug \rightarrow \tilde{u}\tilde{g}, \tilde{u} \rightarrow u\tilde{g} \text{ with a light gluino}$
> 176	95	77 ABACHI	95C	D0	Any $m_{\tilde{g}} < 300 \text{ GeV}; \text{ with cascade decays}$
		78 ABE	95T	CDF	$\tilde{q} \rightarrow \tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \gamma$
> 90	90	79 ABE	92L	CDF	Any $m_{\tilde{g}} < 410 \text{ GeV}; \text{ with cascade decay}$
> 100		80 ROY	92	RVUE	$p\bar{p} \rightarrow \tilde{q}\tilde{q}; \cancel{R}$
		81 NOJIRI	91	COSM	

<sup>1</sup>AAD 13L searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for the production of squarks and gluinos in events containing jets, missing transverse momentum and no high- $p_T$  electrons or muons. No excess over the expected SM background is observed. In mSUGRA/CMSSM models with  $\tan\beta = 10, A_0 = 0$  and  $\mu > 0$ , squarks and gluinos of equal mass are excluded for masses below 1360 GeV at 95% C.L. In a simplified model containing only squarks of the first two generations, a gluino octet and a massless neutralino, squark masses below 1320 GeV are excluded at 95% C.L. for gluino masses below 2 TeV. See Figures 10–15 for more precise bounds.

- <sup>2</sup> AAD 13Q searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing a high- $p_T$  isolated photon, at least one jet identified as originating from a bottom quark, and high missing transverse momentum. Such signatures may originate from supersymmetric models with gauge-mediated supersymmetry breaking in events in which one of a pair of higgsino-like neutralinos decays into a photon and a gravitino while the other decays into a Higgs boson and a gravitino. No significant excess above the expected background was found and limits were set on the squark mass as a function of the neutralino mass in a generalized GMSB model (GGM) with a higgsino-like neutralino NLSP, see their Fig. 4. For neutralino masses greater than 220 GeV, squark masses below 1020 GeV are excluded at 95% C.L.
- <sup>3</sup> CHATRCHYAN 13 looked in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with two opposite-sign leptons ( $e, \mu, \tau$ ), jets and missing transverse energy. No excess beyond the Standard Model expectation is observed. Exclusion limits are derived in the mSUGRA/CMSSM model with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , see Fig. 6.
- <sup>4</sup> CHATRCHYAN 13G searched in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for the production of squarks and gluinos in events containing 0,1,2,  $\geq 3$   $b$ -jets, missing transverse momentum and no electrons or muons. No excess over the expected SM background is observed. In mSUGRA/CMSSM models with  $\tan\beta = 10$ ,  $A_0 = 0$ , and  $\mu > 0$ , squarks and gluinos of equal mass are excluded for masses below 1250 GeV at 95% C.L. Exclusions are also derived in various simplified models, see Fig. 7.
- <sup>5</sup> CHATRCHYAN 13H searched in  $4.96 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with two photons,  $\geq 4$  jets and low  $\cancel{E}_T$  due to  $\tilde{q} \rightarrow \gamma \tilde{\chi}_1^0$  decays in a stealth SUSY framework, where the  $\tilde{\chi}_1^0$  decays through a singlino ( $\tilde{S}$ ) intermediate state to  $\gamma S \tilde{G}$ , with the singlet state  $S$  decaying to two jets. No significant excess above the expected background was found and limits were set in a particular  $R$ -parity conserving stealth SUSY model. The model assumes  $m_{\tilde{\chi}_1^0} = 0.5 m_{\tilde{q}}$ ,  $m_{\tilde{S}} = 100 \text{ GeV}$  and  $m_S = 90 \text{ GeV}$ . Under these assumptions, squark masses less than 1430 GeV were excluded at the 95% C.L.
- <sup>6</sup> CHATRCHYAN 13T searched in  $11.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  for events with at least two energetic jets and significant  $\cancel{E}_T$ , using the  $\alpha_T$  variable to discriminate between processes with genuine and misreconstructed  $\cancel{E}_T$ . No significant excess above the Standard Model expectations is observed. Limits are set on squark masses in simplified models where the decay  $\tilde{q} \rightarrow q \tilde{\chi}_1^0$  takes place with a branching ratio of 100%, assuming an eightfold degeneracy of the masses of the first two generation squarks, see Fig. 8 and Table 9. Also limits in the case of a single light squark are given.
- <sup>7</sup> AAD 12AX searched in  $1.04 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for supersymmetry in events containing jets, missing transverse momentum and one isolated electron or muon. No excess over the expected SM background is observed and model-independent limits are set on the cross section of new physics contributions to the signal regions. In mSUGRA/CMSSM models with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , squarks and gluinos of equal mass are excluded for masses below 820 GeV at 95% C.L. Limits are also set on simplified models for squark production and decay via an intermediate chargino and on supersymmetric models with bilinear  $R$ -parity violation. Supersedes AAD 11G.
- <sup>8</sup> AAD 12CJ searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing one or more isolated leptons (electrons or muons), jets and  $\cancel{E}_T$ . The observations are in good agreement with the SM expectations and exclusion limits have been set in number of SUSY models. In the mSUGRA/CMSSM model with  $\tan\beta = 10$ ,  $A_0 = 0$ , and  $\mu > 0$ , 95% C.L. exclusion limits have been derived for  $m_{\tilde{q}} < 1200 \text{ GeV}$ , assuming equal squark and gluino masses. In minimal GMSB, values of the effective SUSY breaking scale  $\Lambda < 50 \text{ TeV}$  are excluded at 95% C.L. for  $\tan\beta < 45$ . Also exclusion limits in a number of simplified models have been presented, see Figs. 10 and 12.

- <sup>9</sup> AAD 12CP searched in  $4.8 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with two photons and large  $\cancel{E}_T$  due to  $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$  decays in a GMSB framework. No significant excess above the expected background was found and limits were set on the squark mass as a function of the neutralino mass in a generalized GMSB model (GGM) with a bino-like neutralino NLSP. The other sparticle masses were decoupled,  $\tan\beta = 2$  and  $c\tau_{NLSP} < 0.1 \text{ mm}$ . Also, in the framework of the SPS8 model, a 95% C.L. lower limit was set on the breaking scale  $\Lambda$  of 196 TeV.
- <sup>10</sup> AAD 12W searched in  $1.04 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for the production of squarks and gluinos in events containing jets, missing transverse momentum and no electrons or muons. No excess over the expected SM background is observed. In mSUGRA/CMSSM models with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , squarks and gluinos of equal mass are excluded for masses below 950 GeV at 95% C.L. In a simplified model containing only squarks of the first two generations, a gluino octet and a massless neutralino, squark masses below 875 GeV are excluded at 95% C.L.
- <sup>11</sup> CHATRCHYAN 12 looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $e$  and/or  $\mu$  and/or jets, a large total transverse energy, and  $\cancel{E}_T$ . The event selection is based on the dimensionless razor variable  $R$ , related to the  $\cancel{E}_T$  and  $M_R$ , an indicator of the heavy particle mass scale. No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM  $(m_0, m_{1/2})$  plane for  $\tan\beta = 3, 10$  and 50 (see Fig. 7 and 8). Limits are also obtained for Simplified Model Spectra.
- <sup>12</sup> CHATRCHYAN 12AE searched in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with at least three jets and large missing transverse momentum. No significant excesses over the expected SM backgrounds are observed and 95% C.L. limits on the production cross section of squarks in a scenario where  $\tilde{q} \rightarrow q \tilde{\chi}_1^0$  with a 100% branching ratio, see Fig. 3. For  $m_{\tilde{\chi}_1^0} < 200 \text{ GeV}$ , values of  $m_{\tilde{q}}$  below 760 GeV are excluded at 95% C.L. Also limits in the CMSSM are presented, see Fig. 2.
- <sup>13</sup> CHATRCHYAN 12AL looked in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for anomalous production of events with three or more isolated leptons. Limits on squark and gluino masses are set in  $\overline{R}$  SUSY models with leptonic  $LL\overline{E}$  couplings,  $\lambda_{123} > 0.05$ , and hadronic  $\overline{UDD}$  couplings,  $\lambda_{112}'' > 0.05$ , see their Fig. 5. In the  $\overline{UDD}$  case the leptons arise from supersymmetric cascade decays. A very specific supersymmetric spectrum is assumed. All decays are prompt.
- <sup>14</sup> CHATRCHYAN 12AT searched in  $4.73 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for the production of squarks and gluinos in events containing jets, missing transverse momentum and no electrons or muons. No excess over the expected SM background is observed. In mSUGRA/CMSSM models with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , squarks with masses below 1110 GeV are excluded at 95% C.L. Squarks and gluinos of equal mass are excluded for masses below 1180 GeV at 95% C.L. Exclusions are also derived in various simplified models, see Fig. 6.
- <sup>15</sup> AAD 11B looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with same or opposite charge dileptons ( $e$  or  $\mu$ ) and  $\cancel{E}_T$  from the production of squarks and gluinos with leptonic decays from  $\tilde{\chi}_1^\pm$  or  $\tilde{\chi}_2^0$ . No evidence for an excess over the SM expectation is observed, and limits are derived in the CMSSM  $(m_0, m_{1/2})$  plane (see Fig. 2) and in the  $(m_{\tilde{g}}, m_{\tilde{q}})$  plane under the assumptions  $\tan\beta = 4$ ,  $\mu = 1.5 M$ ,  $m_{\tilde{\chi}_2^0} = M - 100 \text{ GeV}$ ,  $m_{\tilde{\ell}_L} = M/2$ ,  $m_{\tilde{\chi}_1^0} = 100 \text{ GeV}$ , where  $M = \min(m_{\tilde{g}}, m_{\tilde{q}})$  (see Fig. 3). The exclusion limit for a compressed spectrum is 590 GeV for the same charge and 450 GeV for the opposite charge events.
- <sup>16</sup> AAD 11C looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with jets, same flavor opposite charge dileptons ( $e$  or  $\mu$ ) and  $\cancel{E}_T$  from the production of squarks and gluinos with decays  $\tilde{q} \rightarrow q \tilde{\chi}_2^0$  and  $\tilde{\chi}_2^0 \rightarrow \ell^+ \ell^- \tilde{\chi}_1^0$ . No evidence for an excess over the SM expectation is observed, and a limit is derived in the  $(m_{\tilde{g}}, m_{\tilde{q}})$  plane under the assumptions  $\tan\beta = 4$ ,  $\mu = 1.5 M$ ,  $m_{\tilde{\chi}_2^0} = M - 100 \text{ GeV}$ ,  $m_{\tilde{\ell}_L} = M/2$ ,  $m_{\tilde{\chi}_1^0} = 100$

- GeV, where  $M = \min(m_{\tilde{g}}, m_{\tilde{q}})$ . The excluded mass region is shown in a plane of  $(m_{\tilde{g}}, m_{\tilde{q}})$ , see their Fig. 3. The exclusion limit for a compressed spectrum is 503 GeV.
- 17 AAD 11G looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with a single lepton ( $e$  or  $\mu$ ), jets and  $\cancel{E}_T$  from the production of squarks and gluinos. No evidence for an excess over the SM expectation is observed, and a limit is derived in the CMSSM  $(m_0, m_{1/2})$  plane for  $\tan\beta = 3$ , see Fig. 2.
- 18 AAD 11N looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 2$  jets and  $\cancel{E}_T$ . Four signal regions were defined, and the background model was found to be in good agreement with the data. Limits are derived in the  $(m_{\tilde{g}}, m_{\tilde{q}})$  plane (see Fig. 2) for a simplified model where degenerate masses of the squarks of the first two generations are assumed,  $m_{\tilde{\chi}_1^0} = 0$ , and all other masses including third generation squarks are set to 5 TeV. Limits are also derived in the CMSSM  $(m_0, m_{1/2})$  plane (see Fig. 3) for  $\tan\beta = 3$ .
- 19 CHATRCHYAN 11W looked in  $1.14 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 2$  jets, large total jet energy, and  $\cancel{E}_T$ . After combining multi-jet events into two pseudo-jets signal events are selected by a cut on  $\alpha_T = E_T^{j2}/M_T$ , the transverse energy of the less energetic jet over the transverse mass. Given the lack of an excess over the SM backgrounds, limits are derived in the CMSSM  $(m_0, m_{1/2})$  plane (see Fig. 4) for  $\tan\beta = 10$ . The limits are only weakly dependent on  $\tan\beta$  and  $A_0$ .
- 20 AALTONEN 09S searched in  $2 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with at least 2 jets and  $\cancel{E}_T$ . No evidence for a signal is observed. A limit is derived for a mSUGRA scenario in the  $m_{\tilde{q}}$  versus  $m_{\tilde{g}}$  plane, see their Fig. 2. For  $m_{\tilde{g}} < 340 \text{ GeV}$  the bound increases to 400 GeV.
- 21 ABAZOV 08G looked in  $2.1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.96 \text{ TeV}$  for events with acoplanar jets or multijets with large  $\cancel{E}_T$ . No significant excess was found compared to the background expectation. A limit is derived on the masses of squarks and gluinos for specific MSUGRA parameter values, see Figure 3. Similar results would be obtained for a large class of parameter sets. Supersedes the results of ABAZOV 06C.
- 22 ACHARD 04 search for the production of  $\tilde{q}\tilde{q}$  of the first two generations in acoplanar di-jet final states in the 192–209 GeV data. Degeneracy of the squark masses is assumed either for both left and right squarks or for right squarks only, as well as  $B(\tilde{q} \rightarrow q\tilde{\chi}_1^0) = 1$ . See Fig. 7 for the dependence of the limits on  $m_{\tilde{\chi}_1^0}$ . This limit supersedes ACCIARRI 99v.
- 23 ABBOTT 01D looked in  $\sim 108 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8 \text{ TeV}$  for events with  $e\bar{e}$ ,  $\mu\bar{\mu}$ , or  $e\mu$  accompanied by at least 2 jets and  $\cancel{E}_T$ . Excluded regions are obtained in the MSUGRA framework from a scan over the parameters  $0 < m_0 < 300 \text{ GeV}$ ,  $10 < m_{1/2} < 110 \text{ GeV}$ , and  $1.2 < \tan\beta < 10$ .
- 24 BARATE 01 looked for acoplanar dijets +  $\cancel{E}_T$  final states at 189 to 202 GeV. The limit assumes  $B(\tilde{q} \rightarrow q\tilde{\chi}_1^0)=1$ , with  $\Delta m = m_{\tilde{q}} - m_{\tilde{\chi}_1^0}$ . It applies to  $\tan\beta=4$ ,  $\mu=-400 \text{ GeV}$ . See their Fig. 2 for the exclusion in the  $(m_{\tilde{q}}, m_{\tilde{g}})$  plane. These limits include and update the results of BARATE 99Q.
- 25 ABE 96D searched for production of gluinos and five degenerate squarks in final states containing a pair of leptons, two jets, and missing  $E_T$ . The two leptons arise from the semileptonic decays of charginos produced in the cascade decays. The limit is derived for fixed  $\tan\beta = 4.0$ ,  $\mu = -400 \text{ GeV}$ , and  $m_{H^+} = 500 \text{ GeV}$ , and with the cascade decays of the squarks and gluinos calculated within the framework of the Minimal Supergravity scenario.
- 26 CHATRCHYAN 2013AO searched in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with two opposite-sign isolated leptons accompanied by hadronic jets and  $\cancel{E}_T$ . No significant excesses over the expected SM backgrounds are observed and 95% C.L. exclusion limits are derived in the mSUGRA/CMSSM model with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , see Fig. 8.

- 27 CHATRCHYAN 13AV searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for new heavy particle pairs decaying into jets (possibly  $b$ -tagged), leptons and  $\cancel{E}_T$  using the Razor variables. No significant excesses over the expected SM backgrounds are observed and 95% C.L. exclusion limits are derived in the mSUGRA/CMSSM model with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , see Fig. 3. The results are also interpreted in various simplified models, see Fig. 4.
- 28 CHATRCHYAN 13W searched in  $4.93 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with one or more photons, hadronic jets and  $\cancel{E}_T$ . No significant excess above the Standard Model expectations is observed. Limits are set on squark masses in the general gauge-mediated SUSY breaking model (GGM), for both a wino-like and bino-like neutralino NLSP scenario, see Fig. 5.
- 29 DREINER 12A reassesses constraints from CMS (at  $7 \text{ TeV}$ ,  $\sim 4.4 \text{ fb}^{-1}$ ) under the assumption that the first and second generation squarks and the lightest SUSY particle are quasi-degenerate in mass (compressed spectrum).
- 30 DREINER 12A reassesses constraints from CMS (at  $7 \text{ TeV}$ ,  $\sim 4.4 \text{ fb}^{-1}$ ) under the assumption that the first and second generation squarks, the gluino, and the lightest SUSY particle are quasi-degenerate in mass (compressed spectrum).
- 31 AAD 11AE looked in  $34 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 2$  same charge isolated leptons ( $e, \mu$ ) and  $\geq 1$  jet. They are assumed to come from  $\tilde{q}\tilde{q}$  production, where the  $\tilde{q}$  decays to  $\tilde{\chi}_1^\pm$  or  $\tilde{\chi}_2^0$  with equal branching ratios, followed by the decays  $\tilde{\chi}_1^\pm \rightarrow W^\pm \tilde{\chi}_1^0$  and  $\tilde{\chi}_2^0 \rightarrow Z^0 \tilde{\chi}_1^0$ . No evidence for an excess over the expected background is observed. Limits are derived on the cross sections as a function of the masses of the  $\tilde{q}$ ,  $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$  and  $\tilde{\chi}_1^0$  (see Fig. 9 and 10).
- 32 AAD 11AF looked in  $1.34 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with 6 up to 8 jets and  $\cancel{E}_T$ . No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM ( $m_0, m_{1/2}$ ) plane for  $\tan\beta = 10$  (see Fig. 5). The limit improves to  $m_{\tilde{g}} > 680 \text{ GeV}$  for  $m_{\tilde{q}} = 2 m_{\tilde{g}}$ .
- 33 AARON 11 looked in  $255 \text{ pb}^{-1}$  of  $e^+p$  and  $183 \text{ pb}^{-1}$  of  $e^-p$  collisions at  $\sqrt{s} = 319 \text{ GeV}$  for events with at least 1 lepton and jets from  $R_p$  violation with  $LQ\bar{D}$  couplings, assuming dominance of a single  $\lambda'_{ijk}$  coupling. No evidence for an excess over the SM expectation is observed, and limits are derived in the  $(\lambda', m_{\tilde{q}})$  plane for the MSSM with  $\tan\beta = 6$ , see their Figs. 7 and 8. Limits are also derived in a CMSSM-type scenario.
- 34 AARON 11C looked in  $281 \text{ pb}^{-1}$  of  $e^+p$  and  $165 \text{ pb}^{-1}$  of  $e^-p$  collisions at  $\sqrt{s} = 319 \text{ GeV}$  and  $\sqrt{s} = 301 \text{ GeV}$  for contact interactions measured from deviations of the  $d\sigma/dQ^2$  of neutral current events. They are interpreted in the framework of R-parity violation with  $LQ\bar{D}$  couplings. No evidence for an excess over the SM expectation is observed, and limits are derived for  $m_{\tilde{q}}/\lambda'$ , see Table 4.
- 35 CHATRCHYAN 11AC looked in  $36 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 3$  jets, a large total transverse energy, and  $\cancel{E}_T$ . No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM ( $m_0, m_{1/2}$ ) plane and the  $(m_{\tilde{g}}, m_{\tilde{q}})$  plane for  $\tan\beta = 10$  (see Fig. 10). Limits are also obtained for Simplified Model Spectra.
- 36 CHATRCHYAN 11C looked in  $34 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with opposite charge isolated dileptons ( $e$  or  $\mu$ ), jets and  $\cancel{E}_T$  from pair production of  $\tilde{g}$  and  $\tilde{q}$ . No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM ( $m_0, m_{1/2}$ ) plane for  $\tan\beta = 3$  (see Fig. 4).
- 37 CHATRCHYAN 11G looked in  $36 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 2$  isolated photons,  $\geq 1$  jet and  $\cancel{E}_T$ , which may arise in a generalized gauge mediated model from the decay of a  $\tilde{\chi}_1^0$  NLSP. No evidence for an excess over the expected background is observed. Limits are derived in the plane of squark versus gluino mass (see Fig. 4) for several values of  $m_{\tilde{\chi}_1^0}$ .

- 38 CHATRCHYAN 11Q looked in  $36 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with a single isolated lepton ( $e$  or  $\mu$ ),  $\geq 4$  jets and  $\cancel{E}_T$ . No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM  $(m_0, m_{1/2})$  plane for  $\tan\beta = 10$  (see Fig. 7).
- 39 CHATRCHYAN 11V looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 3$  isolated leptons ( $e, \mu$  or  $\tau$ ), with or without jets and  $\cancel{E}_T$ . Multi-lepton final states originate from  $\tilde{q} \rightarrow \tilde{\chi}^0 + X$ , followed by  $\tilde{\chi}^0 \rightarrow \tilde{\ell}^\pm \ell^\mp$  and  $\tilde{\ell} \rightarrow \ell \tilde{G}$ . No evidence for an excess over the expected background is observed. Limits are derived (see Fig. 4) for a GMSB-type scenario with mass-degenerate right-handed sleptons (slepton co-NLSP scenario).
- 40 CHATRCHYAN 11V looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 3$  isolated leptons ( $e, \mu$  or  $\tau$ ), with or without jets and  $\cancel{E}_T$ . No evidence for an excess over the expected background is observed. Limits are derived in the  $\tilde{R}$  framework (see Fig. 4) in the  $(m_{\tilde{g}}, m_{\tilde{q}})$  plane assuming the dominance of a  $\lambda_{122}$  or  $\lambda_{123}$  coupling,  $m_{\tilde{\chi}_1^0} = 300 \text{ GeV}$ ,  $m_{\tilde{\ell}} = 1000 \text{ GeV}$ , and decoupled wino and Higgsino.
- 41 KHACHATRYAN 11I looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 2$  jets and  $\cancel{E}_T$ . After combining multi-jet events into two pseudo-jets signal events are selected by a cut on  $\alpha_T = E_T^{j2}/M_T$ , the transverse energy of the less energetic jet over the transverse mass. No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM  $(m_0, m_{1/2})$  plane (see Fig. 5) for  $\tan\beta = 3$ . Superseded by CHATRCHYAN 11W.
- 42 ABAZOV 09S looked in  $0.96 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with at least 2 jets, a tau decaying hadronically and  $\cancel{E}_T$  from the production  $\tilde{q}_L \tilde{q}_R$ , with the taus originating from the decay of a  $\tilde{\chi}_2^0$  or  $\tilde{\chi}_1^\pm$ . The results were combined with ABAZOV 08G which searched for events with jets and  $\cancel{E}_T$  without requiring taus. No evidence for an excess over the SM expectation is observed. The excluded region is shown for an mSUGRA model in a plane of  $m_{1/2}$  versus  $m_0$  in the "tau corridor," see their Figs. 5 and 6. The largest excluded squark mass in the corridor is 340 GeV for the tau analysis only and 410 GeV for the combined analysis.
- 43 SCHAEEL 07A studied the effect on hadronic cross sections and charge asymmetries of t-channel down-type squark exchange via R-parity violating couplings  $LQ\bar{D}$  at  $\sqrt{s} = 189\text{--}209 \text{ GeV}$ . The limit here refers to the case  $j = 1, 2$  and holds for  $\lambda'_{1jk}$  of electromagnetic strength. The results of this analysis are combined with BARATE 00I.
- 44 CHEKANOV 05A search for lepton flavor violating processes  $e^\pm p \rightarrow \ell X$ , where  $\ell = \mu$  or  $\tau$  with high  $p_T$ , in  $130 \text{ pb}^{-1}$  at 300 and 318 GeV. Such final states may originate from  $LQD$  couplings with simultaneously non-zero  $\lambda'_{1jk}$  and  $\lambda'_{ijk}$  ( $i=2$  or  $3$ ). The quoted mass bounds hold for a  $u$ -type squark, assume a  $\lambda'$  of electromagnetic strength and contributions from only direct squark decays. For  $d$ -type squarks the bounds are strengthened to 278 and 275 GeV for the  $\mu$  and  $\tau$  final states, respectively. Supersedes the results of CHEKANOV 02.
- 45 AKTAS 04D looked in  $77.8 \text{ pb}^{-1}$  of  $e^\pm p$  collisions at  $\sqrt{s} = 319 \text{ GeV}$  for resonant production of  $\tilde{q}$  by R-parity violating  $LQ\bar{D}$  couplings assuming that one of the  $\lambda'$  couplings dominates over all others. They consider final states with or without leptons and/or jets and/or  $\cancel{p}_T$  resulting from direct and indirect decays. They combine the channels to derive limits on  $\lambda'_{1j1}$  and  $\lambda'_{11k}$  as a function of the squark mass, see their Figs. 8 and 9, from a scan over the parameters  $70 < M_2 < 350 \text{ GeV}$ ,  $-300 < \mu < 300 \text{ GeV}$ ,  $\tan\beta = 6$ , for a fixed mass of 90 GeV for degenerate sleptons and an LSP mass  $> 30 \text{ GeV}$ . The quoted limits refer to  $\lambda' = 0.3$ , with  $U=u,c,t$  and  $D=d,s,b$ . Supersedes the results of ADLOFF 01B. Superseded by AARON 11.
- 46 ADLOFF 03 looked for the s-channel production of squarks via  $\tilde{R} LQ\bar{D}$  couplings in  $117.2 \text{ pb}^{-1}$  of  $e^+ p$  data at  $\sqrt{s} = 301$  and  $319 \text{ GeV}$  and of  $e^- p$  data at  $\sqrt{s} = 319$

- GeV. The comparison of the data with the SM differential cross section allows limits to be set on couplings for processes mediated through contact interactions. They obtain lower bounds on the value of  $m_{\tilde{q}}/\lambda'$  of 710 GeV for the process  $e^+ \bar{u} \rightarrow \tilde{d}^k$  (and charge conjugate), mediated by  $\lambda'_{11k}$ , and of 430 GeV for the process  $e^+ d \rightarrow \tilde{u}^j$  (and charge conjugate), mediated by  $\lambda'_{1j1}$ . Superseded by AARON 11C.
- 47 CHEKANOV 03B used  $131.5 \text{ pb}^{-1}$  of  $e^+ p$  and  $e^- p$  data taken at 300 and 318 GeV to look for narrow resonances in the  $eq$  or  $\nu q$  final states. Such final states may originate from  $LQ\bar{D}$  couplings with non-zero  $\lambda'_{1j1}$  (leading to  $\tilde{u}_j$ ) or  $\lambda'_{11k}$  (leading to  $\tilde{d}_k$ ). See their Fig. 8 and explanations in the text for limits. The quoted mass bound assumes that only direct squark decays contribute.
- 48 HEISTER 03G searches for the production of squarks in the case of  $\tilde{R}$  prompt decays with  $\overline{UDD}$  direct couplings at  $\sqrt{s} = 189\text{--}209$  GeV.
- 49 ABAZOV 02F looked in  $77.5 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at 1.8 TeV for events with  $\geq 2\mu + \geq 4$  jets, originating from associated production of squarks followed by an indirect  $\tilde{R}$  decay (of the  $\tilde{\chi}_1^0$ ) via  $LQ\bar{D}$  couplings of the type  $\lambda'_{2jk}$  where  $j=1,2$  and  $k=1,2,3$ . Bounds are obtained in the MSUGRA scenario by a scan in the range  $0 \leq M_0 \leq 400$  GeV,  $60 \leq m_{1/2} \leq 120$  GeV for fixed values  $A_0=0$ ,  $\mu < 0$ , and  $\tan\beta=2$  or 6. The bounds are weaker for  $\tan\beta=6$ . See Figs. 2,3 for the exclusion contours in  $m_{1/2}$  versus  $m_0$  for  $\tan\beta=2$  and 6, respectively.
- 50 ABAZOV 02G search for associated production of gluinos and squarks in  $92.7 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8$  TeV, using events with one electron,  $\geq 4$  jets, and large  $\cancel{E}_T$ . The results are compared to a MSUGRA scenario with  $\mu < 0$ ,  $A_0=0$ , and  $\tan\beta=3$  and allow to exclude a region of the  $(m_0, m_{1/2})$  shown in Fig. 11.
- 51 ABBIENDI 02 looked for events with an electron or neutrino and a jet in  $e^+ e^-$  at 189 GeV. Squarks (or leptoquarks) could originate from a  $LQ\bar{D}$  coupling of an electron with a quark from the fluctuation of a virtual photon. Limits on the couplings  $\lambda'_{1jk}$  as a function of the squark mass are shown in Figs. 8–9, assuming that only direct squark decays contribute.
- 52 ABBIENDI 02B looked for events with an electron or neutrino and a jet in  $e^+ e^-$  at 189–209 GeV. Squarks (or leptoquarks) could originate from a  $LQ\bar{D}$  coupling of an electron with a quark from the fluctuation of a virtual photon. Limits on the couplings  $\lambda'_{1jk}$  as a function of the squark mass are shown in Fig. 4, assuming that only direct squark decays contribute. The quoted limits are read off from Fig. 4. Supersedes the results of ABBIENDI 02.
- 53 ACHARD 02 searches for the production of squarks in the case of  $\tilde{R}$  prompt decays with  $\overline{UDD}$  couplings at  $\sqrt{s}=189\text{--}208$  GeV. The search is performed for direct and indirect decays, assuming one coupling at the time to be nonzero. The limit holds for indirect decays. Stronger limits are reached for  $(\tilde{u}_R, \tilde{d}_R)$  direct (80,56) GeV and  $(\tilde{u}_L, \tilde{d}_L)$  direct or indirect (87,86) GeV decays.
- 54 CHEKANOV 02 search for lepton flavor violating processes  $e^+ p \rightarrow \ell X$ , where  $\ell = \mu$  or  $\tau$  with high  $p_T$ , in  $47.7 \text{ pb}^{-1}$  of  $e^+ p$  collisions at 300 GeV. Such final states may originate from  $LQ\bar{D}$  couplings with simultaneously nonzero  $\lambda'_{1jk}$  and  $\lambda'_{ijk}$  ( $i=2$  or 3). The quoted mass bound assumes that only direct squark decays contribute.
- 55 BARATE 01B searches for the production of squarks in the case of  $\tilde{R}$  prompt decays with  $\overline{LLE}$  indirect or  $\overline{UDD}$  direct couplings at  $\sqrt{s}=189\text{--}202$  GeV. The limit holds for direct decays mediated by  $\tilde{R}$   $\overline{UDD}$  couplings. Limits are also given for  $\overline{LLE}$  indirect decays ( $m_{\tilde{u}_R} > 90$  GeV and  $m_{\tilde{d}_R} > 89$  GeV). Supersedes the results from BARATE 00H.
- 56 BREITWEG 01 searches for squark production in  $47.7 \text{ pb}^{-1}$  of  $e^+ p$  collisions, mediated by  $\tilde{R}$  couplings  $LQ\bar{D}$  and leading to final states with  $\tilde{\nu}$  and  $\geq 1$  jet, complementing the  $e^+ X$  final states of BREITWEG 00E. Limits are derived on  $\lambda' \sqrt{\beta}$ , where  $\beta$  is the

- branching fraction of the squarks into  $e^+ q + \bar{\nu} q$ , as function of the squark mass, see their Fig. 15. The quoted mass limit assumes that only direct squark decays contribute.
- 57 ABBOTT 00C searched in  $\sim 94 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions for events with  $\mu\mu$ +jets, originating from associated production of leptoquarks. The results can be interpreted as limits on production of squarks followed by direct  $\tilde{R}$  decay via  $\lambda'_{2jk} L_2 Q_j d_k^c$  couplings. Bounds are obtained on the cross section for branching ratios of 1 and of 1/2, see their Fig. 4. The former yields the limit on the  $\tilde{u}_L$ . The latter is combined with the bound of ABBOTT 99J from the  $\mu\nu$ +jets channel and of ABBOTT 98E and ABBOTT 98J from the  $\nu\nu$ +jets channel to yield the limit on  $\tilde{d}_R$ .
- 58 ACCIARRI 00P studied the effect on hadronic cross sections of  $t$ -channel down-type squark exchange via  $R$ -parity violating coupling  $\lambda'_{1jk} L_1 Q_j d_k^c$ . The limit here refers to the case  $j=1,2$ , and holds for  $\lambda'_{1jk}=0.3$ . Data collected at  $\sqrt{s}=130\text{--}189$  GeV, superseding the results of ACCIARRI 98J.
- 59 AFFOLDER 00K searched in  $\sim 88 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions for events with 2–3 jets, at least one being  $b$ -tagged, large  $\cancel{E}_T$  and no high  $p_T$  leptons. Such  $\nu\nu$ + $b$ -jets events would originate from associated production of squarks followed by direct  $\tilde{R}$  decay via  $\lambda'_{ij3} L_i Q_j d_3^c$  couplings. Bounds are obtained on the production cross section assuming zero branching ratio to charged leptons.
- 60 BARATE 00I studied the effect on hadronic cross sections and charge asymmetries of  $t$ -channel down-type squark exchange via  $R$ -parity violating coupling  $\lambda'_{1jk} L_1 Q_j d_k^c$ . The limit here refers to the case  $j=1,2$ , and holds for  $\lambda'_{1jk}=0.3$ . A 50 GeV limit is found for up-type squarks with  $k=3$ . Data collected at  $\sqrt{s}=130\text{--}183$  GeV. Superseded by SCHAEEL 07A.
- 61 BREITWEG 00E searches for squark exchange in  $e^+ p$  collisions, mediated by  $\tilde{R}$  couplings  $LQ\bar{D}$  and leading to final states with an identified  $e^+$  and  $\geq 1$  jet. The limit applies to up-type squarks of all generations, and assumes  $B(\tilde{q} \rightarrow qe)=1$ .
- 62 ABBOTT 99 searched for  $\gamma\cancel{E}_T + \geq 2$  jet final states, and set limits on  $\sigma(p\bar{p} \rightarrow \tilde{q}+X) \cdot B(\tilde{q} \rightarrow \gamma\cancel{E}_T X)$ . The quoted limits correspond to  $m_{\tilde{g}} \geq m_{\tilde{q}}$ , with  $B(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \gamma)=1$  and  $B(\tilde{\chi}_1^0 \rightarrow \tilde{G} \gamma)=1$ , respectively. They improve to 310 GeV (360 GeV in the case of  $\gamma\tilde{G}$  decay) for  $m_{\tilde{g}}=m_{\tilde{q}}$ .
- 63 ABBOTT 99K uses events with an electron pair and four jets to search for the decay of the  $\tilde{\chi}_1^0$  LSP via  $\tilde{R} LQ\bar{D}$  couplings. The particle spectrum and decay branching ratios are taken in the framework of minimal supergravity. An excluded region at 95% CL is obtained in the  $(m_0, m_{1/2})$  plane under the assumption that  $A_0=0$ ,  $\mu < 0$ ,  $\tan\beta=2$  and any one of the couplings  $\lambda'_{1jk} > 10^{-3}$  ( $j=1,2$  and  $k=1,2,3$ ) and from which the above limit is computed. For equal mass squarks and gluinos, the corresponding limit is 277 GeV. The results are essentially independent of  $A_0$ , but the limit deteriorates rapidly with increasing  $\tan\beta$  or  $\mu > 0$ .
- 64 ABBOTT 99L consider events with three or more jets and large  $\cancel{E}_T$ . Spectra and decay rates are evaluated in the framework of minimal Supergravity, assuming five flavors of degenerate squarks, and scanning the space of the universal gaugino ( $m_{1/2}$ ) and scalar ( $m_0$ ) masses. See their Figs. 2–3 for the dependence of the limit on the relative value of  $m_{\tilde{q}}$  and  $m_{\tilde{g}}$ .
- 65 ABE 99M looked in  $107 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8$  TeV for events with like sign dielectrons and two or more jets from the sequential decays  $\tilde{q} \rightarrow q\tilde{\chi}_1^0$  and  $\tilde{\chi}_1^0 \rightarrow e q \bar{q}'$ , assuming  $\tilde{R}$  coupling  $L_1 Q_j D_k^c$ , with  $j=2,3$  and  $k=1,2,3$ . They assume five degenerate squark flavors,  $B(\tilde{q} \rightarrow q\tilde{\chi}_1^0)=1$ ,  $B(\tilde{\chi}_1^0 \rightarrow e q \bar{q}')=0.25$  for both  $e^+$  and  $e^-$ , and  $m_{\tilde{g}} \geq$

- 200 GeV. The limit is obtained for  $m_{\tilde{\chi}_1^0} \geq m_{\tilde{q}}/2$  and improves for heavier gluinos or heavier  $\tilde{\chi}_1^0$ .
- 66 ABREU 99G looked for events with an electron or neutrino and a jet in  $e^+e^-$  at 183 GeV. Squarks (or leptoquarks) could originate from a  $LQ\bar{D}$  coupling of an electron with a quark from the fluctuation of a virtual photon. Limits on the couplings  $\lambda'_{1jk}$  as a function of the squark mass are shown in Fig. 4, assuming that only direct squark decays contribute.
- 67 ABBOTT 98E searched in  $\sim 115 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions for events with  $e\nu$ +jets, originating from associated production of squarks followed by direct  $\tilde{R}$  decay via  $\lambda'_{1jk} L_1 Q_j d_k^c$  couplings. Bounds are obtained by combining these results with the previous bound of ABBOTT 97B from the  $ee$ +jets channel and with a reinterpretation of ABACHI 96B  $\nu\nu$ +jets channel.
- 68 ABE 98S looked in  $\sim 110 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8$  TeV for events with  $\mu\mu$ +jets originating from associated production of squarks followed by direct  $\tilde{R}$  decay via  $\lambda'_{2jk} L_2 Q_j d_k^c$  couplings. Bounds are obtained on the production cross section times the square of the branching ratio, see Fig. 2. Mass limits result from the comparison with theoretical cross sections and branching ratio equal to 1 for  $\tilde{u}_L$  and 1/2 for  $\tilde{d}_R$ .
- 69 ACKERSTAFF 98V and ACCIARRI 98J studied the interference of  $t$ -channel squark ( $\tilde{d}_R$ ) exchange via  $R$ -parity violating  $\lambda'_{1jk} L_1 Q_j d_k^c$  coupling in  $e^+e^- \rightarrow q\bar{q}$ . The limit is for  $\lambda'_{1jk}=0.3$ . See paper for related limits on  $\tilde{u}_L$  exchange. Data collected at  $\sqrt{s}=130\text{--}172$  GeV.
- 70 BREITWEG 98 used positron+jet events with missing energy and momentum to look for  $e^+q \rightarrow \tilde{e}\tilde{q}$  via gaugino-like neutralino exchange with decays into  $(e\tilde{\chi}_1^0)(q\tilde{\chi}_1^0)$ . See paper for dependences in  $m_{\tilde{e}}, m_{\tilde{\chi}_1^0}$ .
- 71 DATTA 97 argues that the squark mass bound by ABACHI 95C can be weakened by 10–20 GeV if one relaxes the assumption of the universal scalar mass at the GUT-scale so that the  $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$  in the squark cascade decays have dominant and invisible decays to  $\tilde{\nu}$ .
- 72 DERRICK 97 looked for lepton-number violating final states via  $R$ -parity violating couplings  $\lambda'_{ijk} L_i Q_j d_k$ . When  $\lambda'_{11k} \lambda'_{ijk} \neq 0$ , the process  $eu \rightarrow \tilde{d}_k^* \rightarrow \ell_i u_j$  is possible. When  $\lambda'_{1j1} \lambda'_{ijk} \neq 0$ , the process  $e\bar{d} \rightarrow \tilde{u}_j^* \rightarrow \ell_i \bar{d}_k$  is possible. 100% branching fraction  $\tilde{q} \rightarrow \ell j$  is assumed. The limit quoted here corresponds to  $\tilde{t} \rightarrow \tau q$  decay, with  $\lambda'=0.3$ . For different channels, limits are slightly better. See Table 6 in their paper.
- 73 HEWETT 97 reanalyzed the limits on possible resonances in di-jet mode ( $\tilde{q} \rightarrow q\tilde{g}$ ) from ALITTI 93 quoted in “Limits for Excited  $q$  ( $q^*$ ) from Single Production,” ABE 96 in “SCALE LIMITS for Contact Interactions:  $\Lambda(qqqq)$ ,” and unpublished CDF,  $D\bar{D}$  bounds. The bound applies to the gluino mass of 5 GeV, and improves for lighter gluino. The analysis has gluinos in parton distribution function.
- 74 TEREKHOV 97 improved the analysis of TEREKHOV 96 by including di-jet angular distributions in the analysis.
- 75 AID 96C used positron+jet events with missing energy and momentum to look for  $e^+q \rightarrow \tilde{e}\tilde{q}$  via neutralino exchange with decays into  $(e\tilde{\chi}_1^0)(q\tilde{\chi}_1^0)$ . See the paper for dependences on  $m_{\tilde{e}}, m_{\tilde{\chi}_1^0}$ .
- 76 TEREKHOV 96 reanalyzed the limits on possible resonances in di-jet mode ( $\tilde{u} \rightarrow u\tilde{g}$ ) from ABE 95N quoted in “MASS LIMITS for  $g_A$  (axigluon).” The bound applies only to the case with a light gluino.
- 77 ABACHI 95C assume five degenerate squark flavors with  $m_{\tilde{q}_L} = m_{\tilde{q}_R}$ . Sleptons are assumed to be heavier than squarks. The limits are derived for fixed  $\tan\beta = 2.0$   $\mu =$

– 250 GeV, and  $m_{H^+}=500$  GeV, and with the cascade decays of the squarks and gluinos calculated within the framework of the Minimal Supergravity scenario. The bounds are weakly sensitive to the three fixed parameters for a large fraction of parameter space. No limit is given for  $m_{\text{gluino}} > 547$  GeV.

- 78 ABE 95T looked for a cascade decay of five degenerate squarks into  $\tilde{\chi}_2^0$  which further decays into  $\tilde{\chi}_1^0$  and a photon. No signal is observed. Limits vary widely depending on the choice of parameters. For  $\mu = -40$  GeV,  $\tan\beta = 1.5$ , and heavy gluinos, the range  $50 < m_{\tilde{q}} \text{ (GeV)} < 110$  is excluded at 90% CL. See the paper for details.
- 79 ABE 92L assume five degenerate squark flavors and  $m_{\tilde{q}_L} = m_{\tilde{q}_R}$ . ABE 92L includes the effect of cascade decay, for a particular choice of parameters,  $\mu = -250$  GeV,  $\tan\beta = 2$ . Results are weakly sensitive to these parameters over much of parameter space. No limit for  $m_{\tilde{q}} \leq 50$  GeV (but other experiments rule out that region). Limits are 10–20 GeV higher if  $B(\tilde{q} \rightarrow q\tilde{\gamma}) = 1$ . Limit assumes GUT relations between gaugino masses and the gauge coupling; in particular that for  $|\mu|$  not small,  $m_{\tilde{\chi}_1^0} \approx m_{\tilde{g}}/6$ . This last relation implies that as  $m_{\tilde{g}}$  increases, the mass of  $\tilde{\chi}_1^0$  will eventually exceed  $m_{\tilde{q}}$  so that no decay is possible. Even before that occurs, the signal will disappear; in particular no bounds can be obtained for  $m_{\tilde{g}} > 410$  GeV.  $m_{H^+}=500$  GeV.
- 80 ROY 92 reanalyzed CDF limits on di-lepton events to obtain limits on squark production in  $R$ -parity violating models. The 100% decay  $\tilde{q} \rightarrow q\tilde{\chi}$  where  $\tilde{\chi}$  is the LSP, and the LSP decays either into  $\ell q\bar{d}$  or  $\ell\ell\bar{e}$  is assumed.
- 81 NOJIRI 91 argues that a heavy squark should be nearly degenerate with the gluino in minimal supergravity not to overclose the universe.

## Long-lived $\tilde{q}$ (Squark) MASS LIMIT

The following are bounds on long-lived scalar quarks, assumed to hadronise into hadrons with lifetime long enough to escape the detector prior to a possible decay. Limits may depend on the mixing angle of mass eigenstates:  $\tilde{q}_1 = \tilde{q}_L \cos\theta_q + \tilde{q}_R \sin\theta_q$ . The coupling to the  $Z^0$  boson vanishes for up-type squarks when  $\theta_u=0.98$ , and for down type squarks when  $\theta_d=1.17$ .

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
>683	95	<sup>1</sup> AAD	13AA ATLS	$\tilde{t}$ , $R$ -hadrons, generic interaction model
>612	95	<sup>2</sup> AAD	13AA ATLS	$\tilde{b}$ , $R$ -hadrons, generic interaction model
>344	95	<sup>3</sup> AAD	13BC ATLS	$R$ -hadrons, $\tilde{t} \rightarrow b\tilde{\chi}_1^0$ , Regge model, lifetime between $10^{-5}$ and $10^3$ s, $m_{\tilde{\chi}_1^0} = 100$ GeV
>379	95	<sup>4</sup> AAD	13BC ATLS	$R$ -hadrons, $\tilde{t} \rightarrow t\tilde{\chi}_1^0$ , Regge model, lifetime between $10^{-5}$ and $10^3$ s, $m_{\tilde{\chi}_1^0} = 100$ GeV
>935	95	<sup>5</sup> CHATRCHYAN	13AB CMS	long-lived $\tilde{t}$ forming $R$ -hadrons, cloud interaction model
>249	95	<sup>6</sup> AALTONEN	09Z CDF	$\tilde{t}$
> 95	95	<sup>7</sup> HEISTER	03H ALEP	$\tilde{u}$

> 92	95	<sup>7</sup> HEISTER	03H	ALEP	$\tilde{d}$
none 2–85	95	<sup>8</sup> ABREU	98P	DLPH	$\tilde{u}_L$
none 2–81	95	<sup>8</sup> ABREU	98P	DLPH	$\tilde{u}_R$
none 2–80	95	<sup>8</sup> ABREU	98P	DLPH	$\tilde{u}, \theta_u=0.98$
none 2–83	95	<sup>8</sup> ABREU	98P	DLPH	$\tilde{d}_L$
none 5–40	95	<sup>8</sup> ABREU	98P	DLPH	$\tilde{d}_R$
none 5–38	95	<sup>8</sup> ABREU	98P	DLPH	$\tilde{d}, \theta_d=1.17$

<sup>1</sup> AAD 13AA searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing colored long-lived particles that hadronize forming  $R$ -hadrons. No significant excess above the expected background was found. Long-lived  $R$ -hadrons containing a  $\tilde{t}$  are excluded for masses up to 683 GeV at 95% C.L in a general interaction model. Also, limits independent of the fraction of  $R$ -hadrons that arrive charged in the muon system were derived, see Fig. 6.

<sup>2</sup> AAD 13AA searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing colored long-lived particles that hadronize forming  $R$ -hadrons. No significant excess above the expected background was found. Long-lived  $R$ -hadrons containing a  $\tilde{b}$  are excluded for masses up to 612 GeV at 95% C.L in a general interaction model. Also, limits independent of the fraction of  $R$ -hadrons that arrive charged in the muon system were derived, see Fig. 6.

<sup>3</sup> AAD 13BC searched in  $5.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  and in  $22.9 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  for bottom squark  $R$ -hadrons that have come to rest within the ATLAS calorimeter and decay at some later time to hadronic jets and a neutralino. In absence of an excess of events above the expected backgrounds, limits are set on sbottom masses for the decay  $\tilde{b} \rightarrow b\tilde{\chi}_1^0$ , for different lifetimes, and for a neutralino mass of 100 GeV, see their Table 6 and Fig 10.

<sup>4</sup> AAD 13BC searched in  $5.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  and in  $22.9 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  for bottom squark  $R$ -hadrons that have come to rest within the ATLAS calorimeter and decay at some later time to hadronic jets and a neutralino. In absence of an excess of events above the expected backgrounds, limits are set on stop masses for the decay  $\tilde{t} \rightarrow t\tilde{\chi}_1^0$ , for different lifetimes, and for a neutralino mass of 100 GeV, see their Table 6 and Fig 10.

<sup>5</sup> CHATRCHYAN 13AB looked in  $5.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  and in  $18.8 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  for events with heavy stable particles, identified by their anomalous  $dE/dx$  in the tracker or additionally requiring that it be identified as muon in the muon chambers, from pair production of  $\tilde{t}_1$ 's. No evidence for an excess over the expected background is observed. Limits are derived for pair production of stops as a function of mass in the cloud interaction model (see Fig. 8 and Table 6). In the charge-suppressed model, the limit decreases to 818 GeV.

<sup>6</sup> AALTONEN 09Z searched in  $1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with direct production of a pair of charged massive stable particles identified by their TOF. No excess of events is observed over the expected background. The data are used to set a bound on the production cross section, and the result is compared with the pair production cross section of stable stops as a function of the  $\tilde{t}$  mass, see their Fig. 2.

<sup>7</sup> HEISTER 03H use  $e^+e^-$  data at and around the  $Z^0$  peak to look for hadronizing stable squarks. Combining their results on searches for charged and neutral  $R$ -hadrons with JANOT 03, a lower limit of 15.7 GeV on the mass is obtained. Combining this further with the results of searches for tracks with anomalous ionization in data from 183 to 208 GeV yields the quoted bounds.

<sup>8</sup> ABREU 98P assumes that 40% of the squarks will hadronize into a charged hadron, and 60% into a neutral hadron which deposits most of its energy in hadron calorimeter. Data collected at  $\sqrt{s}=130\text{--}183 \text{ GeV}$ .

## $\tilde{b}$ (Sbottom) MASS LIMIT

Limits in  $e^+e^-$  depend on the mixing angle of the mass eigenstate  $\tilde{b}_1 = \tilde{b}_L \cos\theta_b + \tilde{b}_R \sin\theta_b$ . Coupling to the  $Z$  vanishes for  $\theta_b \sim 1.17$ . As a consequence, no absolute constraint in the mass region  $\lesssim 40$  GeV is available in the literature at this time from  $e^+e^-$  collisions. In the Listings below, we use  $\Delta m = m_{\tilde{b}_1} - m_{\tilde{\chi}_1^0}$ .

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
>600	95	1 CHATRCHYAN 13T	CMS	jets + $\cancel{E}_T$ , $\tilde{b} \rightarrow b\tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} = 0$ GeV
>450	95	2 CHATRCHYAN 13V	CMS	same-sign $\ell^\pm \ell^\pm + \geq 2$ $b$ -jets, $\tilde{b} \rightarrow t\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W^\pm \tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} = 50$ GeV
>390		3 AAD	12AN ATLS	$\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ , simplified model, $m_{\tilde{\chi}_1^0} < 60$ GeV
>410	95	4 CHATRCHYAN 12AI	CMS	$\ell^\pm \ell^\pm + b$ -jets + $\cancel{E}_T$
		5 CHATRCHYAN 12BO	CMS	$\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ , simplified model, $m_{\tilde{\chi}_1^0} = 50$ GeV
>230	95	6 AALTONEN	10R CDF	$\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} < 70$ GeV
>247	95	7 ABAZOV	10L D0	$\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} = 0$ GeV
>220	95	8 ABULENCIA	06I CDF	$\tilde{g} \rightarrow \tilde{b}b, \Delta m > 6$ GeV, $\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0, m_{\tilde{g}} < 270$ GeV
> 95		9 ACHARD	04 L3	$\tilde{b} \rightarrow b\tilde{\chi}_1^0, \theta_b=0, \Delta m > 15-25$ GeV
> 81		9 ACHARD	04 L3	$\tilde{b} \rightarrow b\tilde{\chi}_1^0$ , all $\theta_b, \Delta m > 15-25$ GeV
> 7.5	95	10 JANOT	04 THEO	unstable $\tilde{b}_1, e^+e^- \rightarrow$ hadrons
> 93	95	11 ABDALLAH	03M DLPH	$\tilde{b} \rightarrow b\tilde{\chi}_1^0, \theta_b=0, \Delta m > 7$ GeV
> 76	95	11 ABDALLAH	03M DLPH	$\tilde{b} \rightarrow b\tilde{\chi}_1^0$ , all $\theta_b, \Delta m > 7$ GeV
> 85.1	95	12 ABBIENDI	02H OPAL	$\tilde{b} \rightarrow b\tilde{\chi}_1^0$ , all $\theta_b, \Delta m > 10$ GeV, CDF
> 89	95	13 HEISTER	02K ALEP	$\tilde{b} \rightarrow b\tilde{\chi}_1^0$ , all $\theta_b, \Delta m > 8$ GeV, CDF
none 3.5-4.5	95	14 SAVINOV	01 CLEO	$\tilde{B}$ meson
none 80-145		15 AFFOLDER	00D CDF	$\tilde{b} \rightarrow b\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} < 50$ GeV
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
>620	95	16 AAD	13AU ATLS	2 $b$ -jets + $\cancel{E}_T, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} < 120$ GeV
>550	95	17 CHATRCHYAN 13AT	CMS	jets + $\cancel{E}_T, \tilde{b} \rightarrow b\tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} = 50$ GeV
>294	95	18 AAD	11K ATLS	stable $\tilde{b}$
		19 AAD	11O ATLS	$\tilde{g} \rightarrow \tilde{b}_1 b, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} = 60$ GeV
		20 CHATRCHYAN 11D	CMS	$\tilde{b}, \tilde{t} \rightarrow b$
		21 AALTONEN	09R CDF	$\tilde{g} \rightarrow b\tilde{b}, \tilde{b} \rightarrow b\tilde{\chi}_1^0$
>193	95	22 AALTONEN	07E CDF	$\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} = 40$ GeV
none 35-222	95	23 ABAZOV	06R D0	$\tilde{b} \rightarrow b\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} = 50$ GeV

> 78	95	24	ABDALLAH	04M	DLPH	$\tilde{R}, \tilde{b}_L$ , indirect, $\Delta m > 5$ GeV
none 50–82	95	25	ABDALLAH	03C	DLPH	$\tilde{b} \rightarrow b\tilde{g}$ , stable $\tilde{g}$ , all $\theta_b$ , $\Delta m > 10$ GeV
		26	BERGER	03	THEO	
> 71.5	95	27	HEISTER	03G	ALEP	$\tilde{b}_L, \tilde{R}$ decay
> 27.4	95	28	HEISTER	03H	ALEP	$\tilde{b} \rightarrow b\tilde{g}$ , stable $\tilde{g}$ or $\tilde{b}$
> 48	95	29	ACHARD	02	L3	$\tilde{b}_1, \tilde{R}$ decays
		30	BAEK	02	THEO	
		31	BECHER	02	THEO	
		32	CHEUNG	02B	THEO	
		33	CHO	02	THEO	
		34	BERGER	01	THEO	$p\bar{p} \rightarrow X+b$ -quark
none 52–115	95	35	ABBOTT	99F	D0	$\tilde{b} \rightarrow b\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} < 20$ GeV

<sup>1</sup> CHATRCHYAN 13T searched in  $11.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8$  TeV for events with at least two energetic jets and significant  $\cancel{E}_T$ , using the  $\alpha_T$  variable to discriminate between processes with genuine and misreconstructed  $\cancel{E}_T$ . No significant excess above the Standard Model expectations is observed. Limits are set on sbottom masses in simplified models where the decay  $\tilde{b} \rightarrow b\tilde{\chi}_1^0$  takes place with a branching ratio of 100%, see Fig. 8 and Table 9.

<sup>2</sup> CHATRCHYAN 13V searched in  $10.5 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8$  TeV for events with two isolated same-sign dileptons and at least two  $b$ -jets in the final state. No significant excess above the Standard Model expectations is observed. Limits are set on the bottom mass in a simplified models where the decay  $\tilde{b} \rightarrow t\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W^\pm\tilde{\chi}_1^0$  takes place with a branching ratio of 100%, with varying mass of the  $\tilde{\chi}_1^\pm$ , for  $m_{\tilde{\chi}_1^0} = 50$  GeV, see Fig. 4.

<sup>3</sup> AAD 12AN searched in  $2.05 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for scalar bottom quarks in events with large missing transverse momentum and two  $b$ -jets in the final state. The data are found to be consistent with the Standard Model expectations. Limits are set in an  $R$ -parity conserving minimal supersymmetric scenario, assuming  $B(\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0) = 100\%$ , see their Fig. 2.

<sup>4</sup> CHATRCHYAN 12Al looked in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for events with two same-sign leptons ( $e, \mu$ ), but not necessarily same flavor, at least 2  $b$ -jets and missing transverse energy. No excess beyond the Standard Model expectation is observed. Exclusion limits are derived in a simplified model for sbottom pair production, where the sbottom decays through  $\tilde{b}_1 \rightarrow t\tilde{\chi}_1 W$ , see Fig. 8.

<sup>5</sup> CHATRCHYAN 12B0 searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for scalar bottom quarks in events with large missing transverse momentum and two  $b$ -jets in the final state. The data are found to be consistent with the Standard Model expectations. Limits are set in an  $R$ -parity conserving minimal supersymmetric scenario, assuming  $B(\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0) = 100\%$ , see their Fig. 2.

<sup>6</sup> AALTONEN 10R searched in  $2.65 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV for events with  $\cancel{E}_T$  and exactly two jets, at least one of which is  $b$ -tagged. The results are in agreement with the SM prediction, and a limit on the cross section of 0.1 pb is obtained for the range of masses  $80 < m_{\tilde{b}_1} < 280$  GeV assuming that the sbottom decays exclusively to  $b\tilde{\chi}_1^0$ . The excluded mass region in the framework of conserved  $R_p$  is shown in a plane of  $(m_{\tilde{b}_1}, m_{\tilde{\chi}_1^0})$ , see their Fig.2.

<sup>7</sup> ABAZOV 10L looked in  $5.2 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV for events with at least 2  $b$ -jets and  $\cancel{E}_T$  from the production of  $\tilde{b}_1\tilde{b}_1$ . No evidence for an excess over the SM expectation is observed, and a limit on the cross section is derived under the assumption of 100% branching ratio. The excluded mass region in the framework of

- conserved  $R_p$  is shown in a plane of  $(m_{\tilde{b}_1}, m_{\tilde{\chi}_1^0})$ , see their Fig. 3b. The exclusion also extends to  $m_{\tilde{\chi}_1^0} = 110$  GeV for  $160 < m_{\tilde{b}_1} < 200$  GeV.
- 8 ABULENCIA 06I searched in  $156 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96$  TeV for multijet events with large  $\cancel{E}_T$ . They request at least 2  $b$ -tagged jets and no isolated leptons. They investigate the production of gluinos decaying into  $\tilde{b}_1 b$  followed by  $\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ . Both branching fractions are assumed to be 100% and the LSP mass to be 60 GeV. No significant excess was found compared to the background expectation. Upper limits on the cross-section are extracted and a limit is derived on the masses of sbottom and gluinos, see their Fig.3.
  - 9 ACHARD 04 search for the production of  $\tilde{b}\tilde{b}$  in acoplanar  $b$ -tagged di-jet final states in the 192–209 GeV data. See Fig. 6 for the dependence of the limits on  $m_{\tilde{\chi}_1^0}$ . This limit supersedes ACCIARRI 99V.
  - 10 JANOT 04 reanalyzes  $e^+e^- \rightarrow$  hadrons total cross section data with  $\sqrt{s} = 20\text{--}209$  GeV from PEP, PETRA, TRISTAN, SLC, and LEP and constrains the mass of  $\tilde{b}_1$  assuming it decays quickly to hadrons.
  - 11 ABDALLAH 03M looked for  $\tilde{b}$  pair production in events with acoplanar jets and  $\cancel{E}$  at  $\sqrt{s} = 189\text{--}208$  GeV. The limit improves to 87 (98) GeV for all  $\theta_b$  ( $\theta_b = 0$ ) for  $\Delta m > 10$  GeV. See Fig. 24 and Table 11 for other choices of  $\Delta m$ . These limits include and update the results of ABREU,P 00D.
  - 12 ABBIENDI 02H search for events with two acoplanar jets and  $\cancel{p}_T$  in the 161–209 GeV data. The limit assumes 100% branching ratio and uses the exclusion at large  $\Delta m$  from CDF (AFFOLDER 00D). For  $\theta_b=0$ , the bound improves to  $> 96.9$  GeV. See Fig. 4 and Table 6 for the more general dependence on the limits on  $\Delta m$ . These results supersede ABBIENDI 99M.
  - 13 HEISTER 02K search for bottom squarks in final states with acoplanar jets with  $b$  tagging, using 183–209 GeV data. The mass bound uses the CDF results from AFFOLDER 00D. See Fig. 5 for the more general dependence of the limits on  $\Delta m$ . Updates BARATE 01.
  - 14 SAVINOV 01 use data taken at  $\sqrt{s}=10.52$  GeV, below the  $B\bar{B}$  threshold. They look for events with a pair of leptons with opposite charge and a fully reconstructed hadronic  $D$  or  $D^*$  decay. These could originate from production of a light-sbottom hadron followed by  $\tilde{B} \rightarrow D^{(*)} \ell^- \tilde{\nu}$ , in case the  $\tilde{\nu}$  is the LSP, or  $\tilde{B} \rightarrow D^{(*)} \pi \ell^-$ , in case of  $\cancel{R}$ . The mass range  $3.5 \leq M(\tilde{B}) \leq 4.5$  GeV was explored, assuming 100% branching ratio for either of the decays. In the  $\tilde{\nu}$  LSP scenario, the limit holds only for  $M(\tilde{\nu})$  less than about 1 GeV and for the  $D^*$  decays it is reduced to the range 3.9–4.5 GeV. For the  $\cancel{R}$  decay, the whole range is excluded.
  - 15 AFFOLDER 00D search for final states with 2 or 3 jets and  $\cancel{E}_T$ , one jet with a  $b$  tag. See their Fig. 3 for the mass exclusion in the  $m_{\tilde{t}}, m_{\tilde{\chi}_1^0}$  plane.
  - 16 AAD 13AU searched in  $20.1 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8$  TeV for events containing two jets identified as originating from  $b$ -quarks and large missing transverse momentum. No excess of events above the expected level of Standard Model background was found. Exclusion limits at 95% C.L. are set on the masses of third-generation squarks. Assuming that the decay  $\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$  takes place 100% of the time, a  $\tilde{b}_1$  mass below 620 GeV is excluded for  $m_{\tilde{\chi}_1^0} < 120$  GeV. For more details, see their Fig. 5.
  - 17 CHATRCHYAN 13AT provides interpretations of various searches for supersymmetry by the CMS experiment based on  $4.73\text{--}4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV in the framework of simplified models. Limits are set on the sbottom mass in a simplified models where sbottom quarks are pair-produced and the decay  $\tilde{b} \rightarrow b\tilde{\chi}_1^0$  takes place with a branching ratio of 100%, see Fig. 4.
  - 18 AAD 11K looked in  $34 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for events with heavy stable particles, identified by their anomalous  $dE/dx$  in the tracker or time of flight in the tile calorimeter, from pair production of  $\tilde{b}$ . No evidence for an excess over the SM

- expectation is observed and limits on the mass are derived for pair production of sbottom, see Fig. 4.
- 19 AAD 110 looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with jets, of which at least one is a  $b$ -jet, and  $\cancel{E}_T$ . No excess above the Standard Model was found. Limits are derived in the  $(m_{\tilde{g}}, m_{\tilde{b}_1})$  plane (see Fig. 2) under the assumption of 100% branching ratios and  $\tilde{b}_1$  being the lightest squark. The quoted limit is valid for  $m_{\tilde{b}_1} < 500 \text{ GeV}$ . A similar approach for  $\tilde{t}_1$  as the lightest squark with  $\tilde{g} \rightarrow \tilde{t}_1 t$  and  $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$  with 100% branching ratios leads to a gluino mass limit of  $520 \text{ GeV}$  for  $130 < m_{\tilde{t}_1} < 300 \text{ GeV}$ . Limits are also derived in the CMSSM  $(m_0, m_{1/2})$  plane for  $\tan\beta = 40$ , see Fig. 4, and in scenarios based on the gauge group  $\text{SO}(10)$ .
- 20 CHATRCHYAN 11D looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 2$  jets, at least one of which is  $b$ -tagged, and  $\cancel{E}_T$ , where the  $b$ -jets are decay products of  $\tilde{t}$  or  $\tilde{b}$ . No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM  $(m_0, m_{1/2})$  plane for  $\tan\beta = 50$  (see Fig. 2).
- 21 AALTONEN 09R searched in  $2.5 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with at least 2  $b$ -tagged jets and  $\cancel{E}_T$ , originating from the decay  $\tilde{g} \rightarrow b\tilde{b}$  followed by  $\tilde{b} \rightarrow b\tilde{\chi}_1^0$ . Both decays are assumed to have 100% branching ratio. No significant deviation from the SM prediction is observed. An upper limit on the gluino pair production cross section is calculated as a function of the gluino mass, see their Fig. 2. A limit is derived in the  $m_{\tilde{b}}$  versus  $m_{\tilde{g}}$  plane which improves the results of previous searches, see their Fig. 3.
- 22 AALTONEN 07E searched in  $295 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for multijet events with large  $\cancel{E}_T$ . They request at least one heavy flavor-tagged jet and no identified leptons. The branching ratio  $\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$  is assumed to be 100%. No significant excess was found compared to the background expectation. Upper limits on the cross-section are extracted and a limit is derived on the masses of sbottom versus  $\tilde{\chi}_1^0$ , see their Fig. 5. Superseded by AALTONEN 10R.
- 23 ABAZOV 06R looked in  $310 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with 2 or 3 jets and large  $\cancel{E}_T$  with at least 1  $b$ -tagged jet and a veto against isolated leptons. No excess is observed relative to the SM background expectations. Limits are set on the sbottom pair production cross-section under the assumption that the only decay mode is into  $b\tilde{\chi}_1^0$ . Exclusion contours are derived in the plane of sbottom versus neutralino masses, shown in their Fig. 2. The observed limit is more constraining than the expected one due to a lack of events corresponding to large sbottom masses. Superseded by ABAZOV 10L.
- 24 ABDALLAH 04M use data from  $\sqrt{s} = 192\text{--}208 \text{ GeV}$  to derive limits on sparticle masses under the assumption of  $\cancel{R}$  with  $\overline{UDD}$  couplings. The results are valid for  $\mu = -200 \text{ GeV}$ ,  $\tan\beta = 1.5$ ,  $\Delta m > 5 \text{ GeV}$  and assuming a BR of 1 for the given decay. The limit quoted is for indirect  $\overline{UDD}$  decays using the neutralino constraint of  $38.0 \text{ GeV}$ , also derived in ABDALLAH 04M, and assumes no mixing. For indirect decays it remains at  $78 \text{ GeV}$  when the neutralino constraint is not used. Supersedes the result of ABREU 01D.
- 25 ABDALLAH 03C looked for events of the type  $q\bar{q}R^\pm R^\pm$ ,  $q\bar{q}R^\pm R^0$ , or  $q\bar{q}R^0 R^0$  in  $e^+e^-$  interactions at  $\sqrt{s} = 189\text{--}208 \text{ GeV}$ . The  $R^\pm$  bound states are identified by anomalous  $dE/dx$  in the tracking chambers and the  $R^0$  by missing energy due to their reduced energy loss in the calorimeters. Excluded mass regions in the  $(m(\tilde{b}), m(\tilde{g}))$  plane for  $m(\tilde{g}) > 2 \text{ GeV}$  are obtained for several values of the probability for the gluino to fragment into  $R^\pm$  or  $R^0$ , as shown in their Fig. 19. The limit improves to  $94 \text{ GeV}$  for  $\theta_b=0$ .
- 26 BERGER 03 studies the constraints on a  $\tilde{b}_1$  with mass in the  $2.2\text{--}5.5 \text{ GeV}$  region coming from radiative decays of  $\Upsilon(nS)$  into sbottomonium. The constraints apply only if  $\tilde{b}_1$

- lives long enough to permit formation of the sbottomonium bound state. A small region of mass in the  $m_{\tilde{b}_1} - m_{\tilde{g}}$  plane survives current experimental constraints from CLEO.
- 27 HEISTER 03G searches for the production of  $\tilde{b}$  pairs in the case of  $\tilde{R}$  prompt decays with  $L\bar{L}\bar{E}$ ,  $LQ\bar{D}$  or  $U\bar{D}\bar{D}$  couplings at  $\sqrt{s} = 189\text{--}209$  GeV. The limit holds for indirect decays mediated by  $\tilde{R}$   $U\bar{D}\bar{D}$  couplings. It improves to 90 GeV for indirect decays mediated by  $\tilde{R}$   $L\bar{L}\bar{E}$  couplings and to 80 GeV for indirect decays mediated by  $\tilde{R}$   $LQ\bar{D}$  couplings. Supersedes the results from BARATE 01B.
  - 28 HEISTER 03H use their results on bounds on stable squarks, on stable gluinos and on squarks decaying to a stable gluino from the same paper to derive a mass limit on  $\tilde{b}$ , see their Fig. 13. The limit for a long-lived  $\tilde{b}_1$  is 92 GeV.
  - 29 ACHARD 02 searches for the production of squarks in the case of  $\tilde{R}$  prompt decays with  $U\bar{D}\bar{D}$  couplings at  $\sqrt{s}=189\text{--}208$  GeV. The search is performed for direct and indirect decays, assuming one coupling at the time to be nonzero. The limit is computed for the minimal cross section and holds for indirect decays and reaches 55 GeV for direct decays.
  - 30 BAEK 02 studies the constraints on a  $\tilde{b}_1$  with mass in the 2.2–5.5 GeV region coming from precision measurements of  $Z^0$  decays. It is noted that  $CP$ -violating couplings in the MSSM parameters relax the strong constraints otherwise derived from  $CP$  conservation.
  - 31 BECHER 02 studies the constraints on a  $\tilde{b}_1$  with mass in the 2.2–5.5 GeV region coming from radiative  $B$  meson decays, and sets limits on the off-diagonal flavor-changing couplings  $q\tilde{b}\tilde{g}$  ( $q=d,s$ ).
  - 32 CHEUNG 02B studies the constraints on a  $\tilde{b}_1$  with mass in the 2.2–5.5 GeV region and a gluino in the mass range 12–16 GeV, using precision measurements of  $Z^0$  decays and  $e^+e^-$  annihilations at LEP2. Few detectable events are predicted in the LEP2 data for the model proposed by BERGER 01.
  - 33 CHO 02 studies the constraints on a  $\tilde{b}_1$  with mass in the 2.2–5.5 GeV region coming from precision measurements of  $Z^0$  decays. Strong constraints are obtained for  $CP$ -conserving MSSM couplings.
  - 34 BERGER 01 reanalyzed interpretation of Tevatron data on bottom-quark production. Argues that pair production of light gluinos ( $m \sim 12\text{--}16$  GeV) with subsequent 2-body decay into a light sbottom ( $m \sim 2\text{--}5.5$  GeV) and bottom can reconcile Tevatron data with predictions of perturbative QCD for the bottom production rate. The sbottom must either decay hadronically via a  $R$ -parity- and  $B$ -violating interaction, or be long-lived. Constraints on the mass spectrum are derived from the measurements of time-averaged  $B^0\text{--}\bar{B}^0$  mixing.
  - 35 ABBOTT 99F looked for events with two jets, with or without an associated muon from  $b$  decay, and  $\cancel{E}_T$ . See Fig. 2 for the dependence of the limit on  $m_{\tilde{\chi}_1^0}$ . No limit for  $m_{\tilde{\chi}_1^0} > 47$  GeV. Superseded by ABAZOV 06R.

## $\tilde{t}$ (Stop) MASS LIMIT

Limits depend on the decay mode. In  $e^+e^-$  collisions they also depend on the mixing angle of the mass eigenstate  $\tilde{t}_1 = \tilde{t}_L \cos\theta_t + \tilde{t}_R \sin\theta_t$ . The coupling to the  $Z$  vanishes when  $\theta_t = 0.98$ . In the Listings below, we use  $\Delta m \equiv m_{\tilde{t}_1} - m_{\tilde{\chi}_1^0}$  or  $\Delta m \equiv$

$m_{\tilde{t}_1} - m_{\tilde{\nu}}$ , depending on relevant decay mode. See also bounds in “ $\tilde{q}$  (Squark)

MASS LIMIT.” Limits made obsolete by the most recent analyses of  $e^+e^-$  and  $p\bar{p}$  collisions can be found in previous Editions of this *Review*.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
none 123–167	95	<sup>1</sup> AAD	13T ATLS	1 or 2 $\ell^\pm + b$ -jets + $\cancel{E}_T$ , $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$ , $m_{\tilde{\chi}_1^0} = 55$ GeV, $m_{\tilde{\chi}_1^\pm} = 106$ GeV

> 650	95	2	CHATRCHYAN 13BS CMS	1	$\ell^\pm + \text{jets} + \cancel{E}_T, \tilde{t} \rightarrow t\tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} = 0$ GeV	
> 240	95	3	AAD 12AH ATLS		$Z + \text{jets} + \cancel{E}_T, \text{GMSB}, m_{\tilde{\chi}_1^0} > m_Z$	
none 300	95	4	AAD 12CB ATLS		$\ell^\pm \ell^\mp + \text{jets} + \cancel{E}_T, \tilde{t}_1 \rightarrow$ $t\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} = 0 \text{ GeV}$	
none 370–465	95	5	AAD 12CE ATLS		$\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0, \text{hadronic } t \text{ decays},$ $m_{\tilde{\chi}_1^0} = 0 \text{ GeV}$	
none 230–440	95	6	AAD 12CF ATLS		$\ell^\pm + \text{jets} + \cancel{E}_T, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0,$ $m_{\tilde{\chi}_1^0} = 0 \text{ GeV}$	
> 130	95	7	AAD 12CL ATLS		$\ell^\pm \ell^\mp + \text{jets} + \cancel{E}_T, \tilde{t}_1 \rightarrow$ $b\tilde{\chi}_1^\pm, m_{\tilde{\chi}_1^\pm} = 106 \text{ GeV}$	
		8	AAD 12J ATLS		$pp \rightarrow e\mu + X, R$	
> 180	95	9	AALTONEN 12AO CDF		$\tilde{t}_1 \rightarrow c\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} = 90 \text{ GeV}$	
> 200	95	10	ABAZOV 12H D0		$\tilde{t}_1 \tilde{t}_1 \rightarrow b\bar{b}\mu\tau\tilde{\nu}\tilde{\nu}, m_{\tilde{\nu}} = 45 \text{ GeV}$	
	95	11	ABAZOV 12L D0		long-lived $\tilde{q}$ forming $R$ -hadrons	
> 210	95	12	ABAZOV 11N D0		$\tilde{t}_1 \rightarrow b\ell\tilde{\nu}, m_{\tilde{\nu}} < 110 \text{ GeV},$ $m_{\tilde{\nu}} - m_{\tilde{\nu}'} > 30 \text{ GeV}$	
none 60–180	95	13	AALTONEN 10Y CDF		$\tilde{t}_1 \rightarrow b\ell\tilde{\nu}, m_{\tilde{\nu}} = 45 \text{ GeV}$	
none 95–150	95	14	ABAZOV 08Z D0		$\tilde{t} \rightarrow c\tilde{\chi}_1^0,$ $m_c < \Delta m < m_W + m_b$	
none 80–120	95	15	ABAZOV 04 D0		$\tilde{t} \rightarrow b\ell\nu\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} = 50 \text{ GeV}$	
> 90		16	ACHARD 04 L3		$\tilde{t} \rightarrow c\tilde{\chi}_1^0, \text{all } \theta_t, \Delta m >$ $15\text{--}25 \text{ GeV}$	
> 93		16	ACHARD 04 L3		$\tilde{b} \rightarrow b\ell\tilde{\nu}, \text{all } \theta_t, \Delta m > 15 \text{ GeV}$	
> 88		16	ACHARD 04 L3		$\tilde{b} \rightarrow b\tau\tilde{\nu}, \text{all } \theta_t, \Delta m > 15 \text{ GeV}$	
> 75	95	17	ABDALLAH 03M DLPH		$\tilde{t} \rightarrow c\tilde{\chi}_1^0, \theta_t = 0, \Delta m > 2 \text{ GeV}$	
> 71	95	17	ABDALLAH 03M DLPH		$\tilde{t} \rightarrow c\tilde{\chi}_1^0, \text{all } \theta_t, \Delta m > 2 \text{ GeV}$	
> 96	95	17	ABDALLAH 03M DLPH		$\tilde{t} \rightarrow c\tilde{\chi}_1^0, \theta_t = 0, \Delta m > 10 \text{ GeV}$	
> 92	95	17	ABDALLAH 03M DLPH		$\tilde{t} \rightarrow c\tilde{\chi}_1^0, \text{all } \theta_t, \Delta m > 10 \text{ GeV}$	
> 95.7	95	18	ABBIENDI 02H OPAL		$c\tilde{\chi}_1^0, \text{all } \theta_t, \Delta m > 10 \text{ GeV}$	
> 92.6	95	18	ABBIENDI 02H OPAL		$b\ell\tilde{\nu}, \text{all } \theta_t, \Delta m > 10 \text{ GeV}$	
> 91.5	95	18	ABBIENDI 02H OPAL		$b\tau\tilde{\nu}, \text{all } \theta_t, \Delta m > 10 \text{ GeV}$	
> 63	95	19	HEISTER 02K ALEP		any decay, any lifetime, all $\theta_t$	
> 92	95	19	HEISTER 02K ALEP		$\tilde{t} \rightarrow c\tilde{\chi}_1^0, \text{all } \theta_t, \Delta m > 8 \text{ GeV},$ CDF	
> 97	95	19	HEISTER 02K ALEP		$\tilde{t} \rightarrow b\ell\tilde{\nu}, \text{all } \theta_t, \Delta m > 8 \text{ GeV},$ DØ	
> 78	95	19	HEISTER 02K ALEP		$\tilde{t} \rightarrow b\tilde{\chi}_1^0 W^*, \text{all } \theta_t, \Delta m > 8$ GeV	
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●						
> 580	95	20	AAD 13AU ATLS	2	$b\text{-jets} + \cancel{E}_T, \tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm,$ $m_{\tilde{\chi}_1^\pm} - m_{\tilde{\chi}_1^0} = 5 \text{ GeV}, m_{\tilde{\chi}_1^0}$ $= 100 \text{ GeV}$	

>1020		21	CHATRCHYAN 13BN	CMS	$\geq 3$ leptons + $b$ -jets, $\cancel{R}$ , $LL\bar{E}$ , $\lambda_{123} \neq 0$ , $m_{\tilde{\chi}_1^0} = 200$ GeV
> 820		21	CHATRCHYAN 13BN	CMS	$\geq 3$ leptons + $b$ -jets, $\cancel{R}$ , $LL\bar{E}$ , $\lambda_{233} \neq 0$ , $m_{\tilde{\chi}_1^0} = 200$ GeV
> 525	95	22	CHATRCHYAN 13M	CMS	$\tau$ lepton + $b$ -jet, $\cancel{R}$ , $L_3 Q_3 D_3$ , $\tilde{t}_1 \rightarrow \tau b$ simplified model
> 600	95	23	HAN	13 RVUE	natural SUSY, combination of $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$ and $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$
> 340	95	24	CHATRCHYAN 12AN	CMS	long-lived $\tilde{t} \rightarrow t\tilde{\chi}_1^0$
> 737	95	25	CHATRCHYAN 12L	CMS	long-lived $\tilde{t}$ forming $R$ -hadrons
> 309	95	26	AAD	11K ATLS	stable $\tilde{t}$
> 202	95	27	KHACHATRY...11C	CMS	stable $\tilde{t}_1$
none 128–135	95	28	AALTONEN	100 CDF	$\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm \rightarrow b\ell\tilde{\chi}_1^0\nu$ , $m_{\tilde{\chi}_1^\pm} = 106$ GeV, $m_{\tilde{\chi}_1^0} = 48$ GeV
		29	ABAZOV	09N D0	$\tilde{t} \rightarrow b\tilde{\chi}_1^\pm$
		30	ABAZOV	09O D0	$\tilde{t} \rightarrow b\ell\tilde{\nu}$
> 153	95	31	AALTONEN	08Z CDF	$\cancel{R}$ , $\tilde{t}_1 \rightarrow b\tau$
> 185	95	32	ABAZOV	08 D0	$\tilde{t} \rightarrow b\ell\tilde{\nu}$ , $m_{\tilde{\nu}}=70$ GeV
> 132		33	AALTONEN	07E CDF	$\tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$ , $m_{\tilde{\chi}_1^0}=48$ GeV
none 80–134	95	34	ABAZOV	07B D0	$\tilde{t} \rightarrow c\tilde{\chi}_1^0$ , $m_{\tilde{\chi}_1^0} < 48$ GeV
		35	CHEKANOV	07 ZEUS	$e^+p \rightarrow \tilde{t}_1, \cancel{R}, LQ\bar{D}$
> 77	95	36	ABBIENDI	04F OPAL	$\cancel{R}$ , direct, all $\theta_t$
> 77	95	37	ABDALLAH	04M DLPH	$\cancel{R}$ , indirect, all $\theta_t$ , $\Delta m > 5$ GeV
		38	AKTAS	04B H1	$\cancel{R}, \tilde{t}_1$
> 74.5		39	DAS	04 THEO	$\tilde{t}\tilde{t} \rightarrow b\ell\nu_\ell\chi^0\bar{b}q\bar{q}'\chi^0$ , $m_{\tilde{\chi}_1^0} = 15$ GeV, no $\bar{t} \rightarrow c\chi^0$
none 50–87	95	40	ABDALLAH	03C DLPH	$\tilde{t} \rightarrow c\tilde{g}$ , stable $\tilde{g}$ , all $\theta_t$ , $\Delta M > 10$ GeV
none 80–131	95	41	ACOSTA	03C CDF	$\tilde{t} \rightarrow b\ell\tilde{\nu}$ , $m_{\tilde{\nu}} \leq 63$ GeV
		42	CHAKRAB...	03 THEO	$p\bar{p} \rightarrow \tilde{t}\tilde{t}^*$ , RPV
> 71.5	95	43	HEISTER	03G ALEP	$\tilde{t}_L, \cancel{R}$ decay
> 80	95	44	HEISTER	03H ALEP	$\tilde{t} \rightarrow c\tilde{g}$ , stable $\tilde{g}$ or $\tilde{t}$ , all $\theta_t$ , all $\Delta M$
> 144	95	45	ABAZOV	02C D0	$\tilde{t} \rightarrow b\ell\tilde{\nu}$ , $m_{\tilde{\nu}}=45$ GeV
> 77	95	46	ACHARD	02 L3	$\tilde{t}_1, \cancel{R}$ decays
		47	AFFOLDER	01B CDF	$t \rightarrow \tilde{t}\chi_1^0$
> 61	95	48	ABREU	00I DLPH	$\cancel{R} (LL\bar{E})$ , $\theta_t=0.98, \Delta m > 4$ GeV
none 68–119	95	49	AFFOLDER	00D CDF	$\tilde{t} \rightarrow c\tilde{\chi}_1^0$ , $m_{\tilde{\chi}_1^0} < 40$ GeV
none 84–120	95	50	AFFOLDER	00G CDF	$\tilde{t}_1 \rightarrow b\ell\tilde{\nu}$ , $m_{\tilde{\nu}} < 45$ GeV
> 120	95	51	ABE	99M CDF	$p\bar{p} \rightarrow \tilde{t}_1\tilde{t}_1, \cancel{R}$
none 9–24.4	95	52	AID	96 H1	$e p \rightarrow \tilde{t}\tilde{t}, \cancel{R}$ decays
> 138	95	53	AID	96 H1	$e p \rightarrow \tilde{t}, \cancel{R}, \lambda\cos\theta_t > 0.03$
> 45		54	CHO	96 RVUE	$B^0\text{-}\bar{B}^0$ and $\epsilon$ , $\theta_t=0.98, \tan\beta < 2$
none 11–41	95	55	BUSKULIC	95E ALEP	$\cancel{R} (LL\bar{E})$ , $\theta_t=0.98$
none 6.0–41.2	95		AKERS	94K OPAL	$\tilde{t} \rightarrow c\tilde{\chi}_1^0$ , $\theta_t=0, \Delta m > 2$ GeV

none 5.0–46.0	95	AKERS	94k	OPAL	$\tilde{t} \rightarrow c\tilde{\chi}_1^0, \theta_t=0, \Delta m > 5 \text{ GeV}$
none 11.2–25.5	95	AKERS	94k	OPAL	$\tilde{t} \rightarrow c\tilde{\chi}_1^0, \theta_t=0.98, \Delta m > 2 \text{ GeV}$
none 7.9–41.2	95	AKERS	94k	OPAL	$\tilde{t} \rightarrow c\tilde{\chi}_1^0, \theta_t=0.98, \Delta m > 5 \text{ GeV}$
none 7.6–28.0	95	<sup>56</sup> SHIRAI	94	VNS	$\tilde{t} \rightarrow c\tilde{\chi}_1^0, \text{any } \theta_t, \Delta m > 10 \text{ GeV}$
none 10–20	95	<sup>56</sup> SHIRAI	94	VNS	$\tilde{t} \rightarrow c\tilde{\chi}_1^0, \text{any } \theta_t, \Delta m > 2.5 \text{ GeV}$

<sup>1</sup> AAD 13T searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for pair production of light  $\tilde{t}_1$  squarks with masses similar to, or lighter than, the top quark mass. Final states containing exclusively one or two leptons (electrons or muons), large missing transverse momentum, light jets and  $b$ -jets are used to reconstruct the top squark pair system. The  $\tilde{t}_1$  is assumed to decay through  $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$  with a 100 % branching ratio. The chargino is then assumed to decay through a virtual  $W$  boson,  $\tilde{\chi}_1^\pm \rightarrow W^*\tilde{\chi}_1^0$ . The data are found to be consistent with the Standard Model expectations. The results are interpreted in a simplified model as a function of  $m_{\tilde{t}_1}$  and  $m_{\tilde{\chi}_1^0}$ , for either  $m_{\tilde{\chi}_1^\pm} = 2m_{\tilde{\chi}_1^0}$  or for a fixed choice of  $m_{\tilde{\chi}_1^\pm} = 106 \text{ GeV}$ , see Fig. 2. Assuming  $m_{\tilde{\chi}_1^\pm} = 106 \text{ GeV}$ ,  $\tilde{t}_1$  masses between 123 and 167 GeV are excluded at 95% C.L for  $m_{\tilde{\chi}_1^0} = 55 \text{ GeV}$ .

<sup>2</sup> CHATRCHYAN 13BS searched in  $19.5 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  for events with a single isolated lepton, hadronic jets, large  $\cancel{E}_T$ , and large transverse mass. No significant excess above the Standard Model expectations is observed. Limits are set on stop masses in simplified models where the decay  $\tilde{t} \rightarrow t\tilde{\chi}_1^0$  takes place with a branching ratio which has been varied between 50% and 100%, see Fig. 12, and where the decay  $\tilde{t} \rightarrow b\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W^\pm\tilde{\chi}_1^0$  takes place with a branching ratio of 100%, with varying intermediate mass of the  $\tilde{\chi}_1^\pm$ , see Fig. 11.

<sup>3</sup> AAD 12AH searched in  $2.05 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for pair production of  $\tilde{t}_1$  in events with two same-flavor, opposite-sign leptons ( $e$  or  $\mu$ ) with invariant mass consistent with the  $Z$  boson, large missing transverse momentum and jets in the final state. At least one of the jets is identified as originating from a  $b$ -quark. The data are found to be consistent with the Standard Model expectations. The results are interpreted in a GMSB scenario where the  $\tilde{\chi}_1^0$  is the NLSP and is purely higgsino-like. Other model parameters are  $\tan\beta = 10, m_{\tilde{u}_3} = m_{\tilde{q}_3} = -A_t/2$ . Scalar top masses below 240 GeV are excluded for all values of  $m_{\tilde{\chi}_1^0} > m_Z$ .

<sup>4</sup> AAD 12CB searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for pair production of  $\tilde{t}_1$  in events with two opposite-sign leptons (electrons or muons), jets, and  $\cancel{E}_T$ . The  $\tilde{t}_1$  is assumed to decay through  $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$  with a 100% branching ratio. The data are found to be consistent with the Standard Model expectations. The results are interpreted in a simplified model as a function of  $m_{\tilde{t}_1}$  and  $m_{\tilde{\chi}_1^0}$ , see Fig. 2. Assuming a massless  $\tilde{\chi}_1^0$ , a  $\tilde{t}_1$  with a mass of 300 GeV is excluded at 95% C.L.

<sup>5</sup> AAD 12CE searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for pair production of  $\tilde{t}_1$  where  $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$  with a 100 % branching ratio and where both tops decay hadronically. The data are found to be consistent with the Standard Model expectations. The results are interpreted in a simplified model as a function of  $m_{\tilde{t}_1}$  and  $m_{\tilde{\chi}_1^0}$ , see Fig. 4. For a massless  $\tilde{\chi}_1^0$ , masses of  $\tilde{t}_1$  between 370 GeV and 465 GeV are excluded at 95% C.L. The upper limit deteriorates to 445 GeV for  $m_{\tilde{\chi}_1^0} < 50 \text{ GeV}$ .

- <sup>6</sup> AAD 12CF searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for pair production of  $\tilde{t}_1$  in events with one isolated electron or muon, jets, and  $\cancel{E}_T$ . The  $\tilde{t}_1$  is assumed to decay through  $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$  a 100 % branching ratio. The data are found to be consistent with the Standard Model expectations. The results are interpreted in a simplified model as a function of  $m_{\tilde{t}_1}$  and  $m_{\tilde{\chi}_1^0}$ , see Fig. 2. For a massless  $\tilde{\chi}_1^0$ , masses of  $\tilde{t}_1$  between 230 GeV and 440 GeV are excluded at 95% C.L. The upper limit deteriorates to 400 GeV for  $m_{\tilde{\chi}_1^0} < 125 \text{ GeV}$  and the lower limit is increased to about 330 GeV.
- <sup>7</sup> AAD 12CL searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for pair production of  $\tilde{t}_1$  in events with two opposite-sign leptons (electrons or muons), jets, and  $\cancel{E}_T$ . The  $\tilde{t}_1$  is assumed to decay through  $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$  with a 100 % branching ratio. The chargino is then assumed to decay through a virtual  $W$  boson,  $\tilde{\chi}_1^\pm \rightarrow W^*\tilde{\chi}_1^0$ . The data are found to be consistent with the Standard Model expectations. The results are interpreted in a simplified model as a function of  $m_{\tilde{t}_1}$  and  $m_{\tilde{\chi}_1^0}$  for a fixed choice of  $m_{\tilde{\chi}_1^\pm} = 106 \text{ GeV}$ , see Fig. 2. Assuming  $m_{\tilde{\chi}_1^\pm} = 106 \text{ GeV}$ ,  $\tilde{t}_1$  masses below 130 GeV are excluded at 95% C.L. for  $m_{\tilde{\chi}_1^0} < 70 \text{ GeV}$ .
- <sup>8</sup> AAD 12J looked in  $2.1 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for evidence of lepton flavor violating interactions in the  $e\mu$  continuum due to a  $t$ -channel exchange of an  $R$ -parity violating scalar top quark. No deviations from the SM expectations were found. Limits on  $R$ -parity violating couplings are calculated as a function of the scalar stop mass, see their Fig. 4b.
- <sup>9</sup> AALTONEN 12AO searched in  $2.6 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events containing  $\cancel{E}_T$  and at least two jets, of which at least one is identified as originating from a charm quark. No excess over the expected SM background is observed. Limits are set on the production of  $\tilde{t}_1$  in the assumption that the only decay model is into  $c\tilde{\chi}_1^0$  and for  $m_{\tilde{\chi}_1^0} = 90 \text{ GeV}$ , see Fig. 2. According to Fig. 2 there is an exclusion gap from 100–130 GeV.
- <sup>10</sup> ABAZOV 12H looked in  $7.3 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events containing one muon, one tau decaying hadronically, at least one jet, and missing transverse energy. No evidence for an excess over the SM expectation is observed and 95% C.L. limits are set in the plane  $(m_{\tilde{t}_1}, m_{\tilde{\nu}}$ ), see their Fig. 5 (where  $B(\tilde{t}_1 \rightarrow b\mu\tilde{\nu}) = B(\tilde{t}_1 \rightarrow b\tau\tilde{\nu}) = 1/3$ ) and Fig. 6 (where  $B(\tilde{t}_1 \rightarrow b\mu\tilde{\nu}) = 0.1$  and  $B(\tilde{t}_1 \rightarrow b\tau\tilde{\nu}) = 0.8$ ).
- <sup>11</sup> ABAZOV 12L looked in  $5.2 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for charged massive long-lived particles in events in which one or more particles are reconstructed as muons but have speed and ionization energy loss inconsistent with muons produced in beam collisions. Long-lived stops with mass below 285 GeV are excluded at 95% C.L., using the nominal value of the NLO production cross section. For the latter, a charge survival probability of 38% has been assumed.
- <sup>12</sup> ABAZOV 11N looked in  $5.4 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with exactly one  $e$  and  $\mu$  and  $\cancel{E}_T$  from the production of  $\tilde{t}_1\tilde{t}_1$ . No evidence for an excess over the SM expectation is observed, and a limit is derived in a plane of  $(m_{\tilde{t}_1}, m_{\tilde{\nu}})$ , see their Fig. 4, under the assumption of 100% branching ratio for  $\tilde{t}_1 \rightarrow b\ell\tilde{\nu}$ .
- <sup>13</sup> AALTONEN 10Y searched in  $1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with an oppositely charged lepton pair ( $e$  or  $\mu$ ),  $\cancel{E}_T$  and at least one jet. A limit is derived on the cross section assuming 100% branching ratio of  $\tilde{t}_1 \rightarrow b\ell\tilde{\nu}$  and an invisible  $\tilde{\nu}$ , see their Fig. 10. In Fig. 11, the exclusion contour is shown in the plane of  $(m_{\tilde{t}_1}, m_{\tilde{\nu}})$ .
- <sup>14</sup> ABAZOV 08Z looked in  $995 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with exactly 2 jets, at least one being tagged as heavy quark, and  $\cancel{E}_T$ , originating from stop pair production. Branching ratios are assumed to be 100% for  $\tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$ . No evidence for an excess over the SM expectation is observed. The excluded region is shown in a

- plane of  $m_{\tilde{t}}$  versus  $m_{\tilde{\chi}_1^0}$ , see their Fig. 5. No limit can be obtained for  $m_{\tilde{\chi}_1^0} > 70$  GeV. Supersedes the results of ABAZOV 07B.
- 15 ABAZOV 04 looked at  $108.3pb^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.8$  TeV for events with  $e+\mu+\cancel{E}_T$  as signature for the 3- and 4-body decays of stop into  $b\ell\nu\tilde{\chi}_1^0$  final states. For the  $b\ell\tilde{\nu}$  channel they use the results from ABAZOV 02C. No significant excess is observed compared to the Standard Model expectation and limits are derived on the mass of  $\tilde{t}_1$  for the 3- and 4-body decays in the  $(m_{\tilde{t}}, m_{\tilde{\chi}_1^0})$  plane, see their Figure 4.
- 16 ACHARD 04 search in the 192–209 GeV data for the production of  $\tilde{t}\tilde{t}$  in acoplanar di-jet final states and, in case of  $b\ell\tilde{\nu}$  ( $b\tau\tilde{\nu}$ ) final states, two leptons (taus). The limits for  $\theta_{\tilde{t}}=0$  improve to 95, 96 and 93 GeV, respectively. All limits assume 100% branching ratio for the respective decay modes. See Fig. 6 for the dependence of the limits on  $m_{\tilde{\chi}_1^0}$ . These limits supersede ACCIARRI 99v.
- 17 ABDALLAH 03M looked for  $\tilde{t}$  pair production in events with acoplanar jets and  $\cancel{E}$  at  $\sqrt{s} = 189$ –208 GeV. See Fig. 23 and Table 11 for other choices of  $\Delta m$ . These limits include and update the results of ABREU,P 00D.
- 18 ABBIENDI 02H looked for events with two acoplanar jets,  $\cancel{p}_T$ , and, in the case of  $b\ell\tilde{\nu}$  final states, two leptons, in the 161–209 GeV data. The bound for  $c\tilde{\chi}_1^0$  applies to the region where  $\Delta m < m_W + m_b$ , else the decay  $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^0 W^+$  becomes dominant. The limit for  $b\ell\tilde{\nu}$  assumes equal branching ratios for the three lepton flavors and for  $b\tau\tilde{\nu}$  100% for this channel. For  $\theta_{\tilde{t}}=0$ , the bounds improve to  $> 97.6$  GeV ( $c\tilde{\chi}_1^0$ ),  $> 96.0$  GeV ( $b\ell\tilde{\nu}$ ), and  $> 95.5$  ( $b\tau\tilde{\nu}$ ). See Figs. 5–6 and Table 5 for the more general dependence of the limits on  $\Delta m$ . These results supersede ABBIENDI 99M.
- 19 HEISTER 02K search for top squarks in final states with jets (with/without  $b$  tagging or leptons) or long-lived hadrons, using 183–209 GeV data. The absolute mass bound is obtained by varying the branching ratio of  $\tilde{t} \rightarrow c\tilde{\chi}_1^0$  and the lepton fraction in  $\tilde{t} \rightarrow b\tilde{\chi}_1^0 f\bar{f}'$  decays. The mass bound for  $\tilde{t} \rightarrow c\tilde{\chi}_1^0$  uses the CDF results from AFFOLDER 00D and for  $\tilde{t} \rightarrow b\ell\tilde{\nu}$  the DØ results from ABAZOV 02C. See Figs. 2–5 for the more general dependence of the limits on  $\Delta m$ . Updates BARATE 01 and BARATE 00P.
- 20 AAD 13AU searched in  $20.1\text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8$  TeV for events containing two jets identified as originating from  $b$ -quarks and large missing transverse momentum. No excess of events above the expected level of Standard Model background was found. Exclusion limits at 95% C.L. are set on the masses of third-generation squarks. Assuming that the decay  $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$  takes place 100% of the time, a  $\tilde{t}_1$  mass below 580 GeV (440 GeV) is excluded for  $\Delta m = m_{\tilde{\chi}_1^\pm} - m_{\tilde{\chi}_1^0} = 5$  GeV (20 GeV) and for  $m_{\tilde{\chi}_1^0} = 100$  GeV. For more details, see their Fig. 6.
- 21 CHATRCHYAN 13BN searched in  $19.5\text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8$  TeV for events with three or more isolated leptons and  $b$ -quark jets. No excess above the Standard Model expectations is observed. Limits are set on stop masses in  $\cancel{R}$  SUSY models with leptonic  $LL\bar{E}$  couplings, see Fig. 2. Also limits have been set in a model with  $LQ\bar{D}$  couplings.
- 22 CHATRCHYAN 13M searched in  $4.8\text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for events containing an isolated electron or muon, a hadronically decaying  $\tau$  lepton and two  $b$ -quark jets. No excess above the Standard Model expectations is observed. Limits are set on stop masses in  $\cancel{R}$  SUSY models with  $L_3 Q_3 D_3$  couplings, see Fig. 2. In a simplified model where the decay  $\tilde{t}_1 \rightarrow \tau b$  takes place with a branching ratio of 100%, stop masses below 525 GeV are excluded at 95% C.L.
- 23 HAN 13 used combined ATLAS results based on  $20.1\text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8$  TeV to derive 95% C.L. exclusion limits on the stop mass in the framework of natural SUSY in the MSSM, after considering the constraints from the Higgs mass,  $B$ -physics, and electroweak precision measurements.

- 24 CHATRCHYAN 12AN looked in  $4.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with pair production of long-lived stops. The hadronization of the stops leads to  $R$ -hadrons which may stop inside the detector and later decay via  $\tilde{t} \rightarrow t \tilde{\chi}_1^0$  during gaps between the proton bunches. No significant excess over the expected background is observed. From a counting experiment, a limit at 95% C.L. on the cross section as a function of  $m_{\tilde{t}}$  is derived, see Fig. 4. The mass limit is valid for lifetimes between  $10^{-5}$  and  $10^3$  seconds, for what they call "the daughter top energy  $E_t >$ " 125 GeV and assuming the *cloud* interaction model for  $R$ -hadrons. Supersedes KHACHATRYAN 11.
- 25 CHATRCHYAN 12L looked in  $5.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with heavy stable particles, identified by their anomalous  $dE/dx$  in the tracker or additionally requiring that it be identified as muon in the muon chambers, from pair production of  $\tilde{t}_1$ 's. No evidence for an excess over the expected background is observed. Limits are derived for pair production of stops as a function of mass (see Fig. 3). In the conservative scenario where every hadronic interaction causes it to become neutral, the limit decreases to 626 GeV. Supersedes KHACHATRYAN 11C.
- 26 AAD 11K looked in  $34 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with heavy stable particles, identified by their anomalous  $dE/dx$  in the tracker or time of flight in the tile calorimeter, from pair production of  $\tilde{t}$ . No evidence for an excess over the SM expectation is observed and limits on the mass are derived for pair production of stop, see Fig. 4.
- 27 KHACHATRYAN 11C looked in  $3.1 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with heavy stable particles, identified by their anomalous  $dE/dx$  in the tracker or time of flight in the muon chambers, from pair production of  $\tilde{t}_1$ . No evidence for an excess over the expected background is observed. Limits are derived for pair production of stop as a function of mass, see Fig. 3, and compared to the production cross section in a benchmark scenario.
- 28 AALTONEN 100 searched in  $2.7 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with a charged lepton pair ( $e$  or  $\mu$ ),  $\cancel{E}_T$  and at least two jets. A fit of the data is made to the  $\tilde{t}_1 \tilde{t}_1$  hypothesis. Assuming a 100% branching ratio of  $\tilde{t}_1 \rightarrow b \tilde{\chi}_1^\pm$ , the exclusion is independent of the value of the  $\tilde{\chi}_1^\pm \rightarrow \ell \tilde{\chi}_1^0 \nu$  branching ratio.
- 29 ABAZOV 09N looked in  $0.9 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with  $\geq 3$  jets, at least one being  $b$ -tagged, one electron or muon and  $\cancel{E}_T$  originating from associated production  $\tilde{t}\tilde{t}$ , with one  $\tilde{t}$  decaying leptonically, the other hadronically. The branching ratios for  $\tilde{t}_1 \rightarrow b \tilde{\chi}_1^\pm$  and  $\tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 W^\pm$  are assumed to be 100%. The separation from the dominant  $t\bar{t}$  background is based on a multivariate likelihood discriminant analysis. The tested mass range is  $130 \text{ GeV} \leq m_{\tilde{t}} \leq 190 \text{ GeV}$ ,  $90 \text{ GeV} \leq m_{\tilde{\chi}_1^\pm} \leq 150 \text{ GeV}$  and  $m_{\tilde{\chi}_1^0} = 50 \text{ GeV}$  fixed. The excluded cross section is a factor 2–13 larger than the theoretical expectation in the considered MSSM scenarios, see their Fig. 3.
- 30 ABAZOV 09O looked in  $1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with two electrons or one electron and one muon and  $\cancel{E}_T$  originating from associated production  $\tilde{t}\tilde{t}$ , followed by the three-body decays  $\tilde{t} \rightarrow b\ell\tilde{\nu}$ . No evidence for an excess over the SM expectation is observed. The excluded region is shown in a plane of  $m_{\tilde{\nu}}$  versus  $m_{\tilde{t}}$ , see their Fig. 3. The largest excluded  $\tilde{t}$  mass is 175 GeV for a  $\tilde{\nu}$  mass of 45 GeV, and the largest excluded  $\tilde{\nu}$  mass is 96 GeV for a  $\tilde{t}$  mass of 140 GeV. Superseded by ABAZOV 11N.
- 31 AALTONEN 08Z searched in  $322 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for dijet events with a lepton ( $e$  or  $\mu$ ) and a hadronic  $\tau$  decay produced via  $R$ -parity violating couplings  $LQ\bar{D}$ . No heavy flavour-tagged jets are requested. No significant excess was found compared to the background expectation. Upper limits on the cross-section times the square of the branching ratio  $B(\tilde{t}_1 \rightarrow b\tau)$  are extracted, and a limit is derived on the stop mass assuming  $B(\tilde{t}_1 \rightarrow b\tau) = 1$ , see their Fig. 2. Supersedes the results of ACOSTA 04B.
- 32 ABAZOV 08 looked at approximately  $400 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with  $b\bar{b}\ell\ell'\cancel{E}_T$  with  $\ell\ell' = e^\pm\mu^\mp$  or  $\ell\ell' = \mu^+\mu^-$ , originating from associated

- production  $\tilde{t}\tilde{t}$ . Branching ratios are assumed to be 100% for both  $\tilde{\chi}_1^\pm \rightarrow \ell\tilde{\nu}$  and  $\tilde{\nu} \rightarrow \nu\tilde{\chi}_1^0$ . No evidence for an excess over the SM expectation is observed. The excluded region is shown in a plane of  $m_{\tilde{\nu}}$  versus  $m_{\tilde{t}}$ , see their Fig.3. Superseded by ABAZOV 090.
- 33 AALTONEN 07E searched in  $295 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for multijet events with large  $\cancel{E}_T$ . They request at least one heavy flavor-tagged jet and no identified leptons. The branching ratio  $\tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$  is assumed to be 100%. No significant excess was found compared to the background expectation. Upper limits on the cross-section are extracted and a limit is derived on the masses of stop versus  $\tilde{\chi}_1^0$ , see their Fig. 4.
- 34 ABAZOV 07B looked in  $360 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with a pair of acoplanar heavy-flavor jets with  $\cancel{E}_T$ . No excess is observed relative to the SM background expectations. Limits are set on the production of  $\tilde{t}_1$  under the assumption that the only decay mode is into  $c\tilde{\chi}_1^0$ , see their Fig. 4 for the limit in the  $(m_{\tilde{t}}, m_{\tilde{\chi}_1^0})$  plane. No limit can be obtained for  $m_{\tilde{\chi}_1^0} > 54 \text{ GeV}$ . Supersedes the results of ABAZOV 04B.
- 35 CHEKANOV 07 search for the  $LQ\bar{D}$  R-parity violating process  $e^+p \rightarrow \tilde{t}_1$  in  $65 \text{ pb}^{-1}$  at 318 GeV. Final states may originate from  $LQ\bar{D}$  couplings  $\tilde{t} \rightarrow e^+d$  and from the R-parity conserving decay  $\tilde{t} \rightarrow \tilde{\chi}^+b$ , giving rise to  $e + \text{jet}$ ,  $e + \text{multi-jet}$ , and  $\nu + \text{multi-jet}$ . The excluded region in an MSSM scenario is presented for  $\lambda'_{131}$  as a function of the stop mass in Fig. 6. Other excluded regions in a more restricted mSUGRA model are shown in Fig. 7 and 8.
- 36 ABBIENDI 04F use data from  $\sqrt{s} = 189\text{--}209 \text{ GeV}$ . They derive limits on the stop mass under the assumption of  $\mathcal{R}$  with  $LQ\bar{D}$  or  $\overline{UD\bar{D}}$  couplings. The limit quoted applies to direct decays with  $\overline{UD\bar{D}}$  couplings when the stop decouples from the  $Z^0$  and improves to 88 GeV for  $\theta_t = 0$ . For  $LQ\bar{D}$  couplings, the limit improves to 98 (100) GeV for  $\lambda'_{13k}$  or  $\lambda'_{23k}$  couplings and all  $\theta_t$  ( $\theta_t = 0$ ). For  $\lambda'_{33k}$  couplings it is 96 (98) GeV for all  $\theta_t$  ( $\theta_t = 0$ ). Supersedes the results of ABBIENDI 00.
- 37 ABDALLAH 04M use data from  $\sqrt{s} = 192\text{--}208 \text{ GeV}$  to derive limits on sparticle masses under the assumption of  $\mathcal{R}$  with  $LLE$  or  $\overline{UD\bar{D}}$  couplings. The results are valid for  $\mu = -200 \text{ GeV}$ ,  $\tan\beta = 1.5$ ,  $\Delta m > 5 \text{ GeV}$  and assuming a BR of 1 for the given decay. The limit quoted is for decoupling of the stop from the  $Z^0$  and indirect  $\overline{UD\bar{D}}$  decays using the neutralino constraint of 39.5 GeV for  $LLE$  and of 38.0 GeV for  $\overline{UD\bar{D}}$  couplings, also derived in ABDALLAH 04M. For no mixing (decoupling) and indirect decays via  $LLE$  the limit improves to 92 (87) GeV if the constraint from the neutralino is used and to 88 (81) GeV if it is not used. For indirect decays via  $\overline{UD\bar{D}}$  couplings it improves to 87 GeV for no mixing and using the constraint from the neutralino, whereas it becomes 81 GeV (67) GeV for no mixing (decoupling) if the neutralino constraint is not used. Supersedes the result of ABREU 01D.
- 38 AKTAS 04B looked in  $106 \text{ pb}^{-1}$  of  $e^\pm p$  collisions at  $\sqrt{s} = 319 \text{ GeV}$  and  $301 \text{ GeV}$  for resonant production of  $\tilde{t}_1$  by R-parity violating  $LQ\bar{D}$  couplings couplings with  $\lambda'_{131}$ , others being zero. They consider the decays  $\tilde{t}_1 \rightarrow e^+d$  and  $\tilde{t}_1 \rightarrow W\tilde{b}$  followed by  $\tilde{b} \rightarrow \bar{\nu}_e d$  and assume gauginos too heavy to participate in the decays. They combine the channels  $j\cancel{e}'_T$ ,  $j\mu\cancel{\nu}'_T$ ,  $jj\cancel{\nu}'_T$  to derive limits in the plane  $(m_{\tilde{t}}, \lambda'_{131})$ , see their Fig. 5.
- 39 DAS 04 reanalyzes AFFOLDER 00G data and obtains constraints on  $m_{\tilde{t}_1}$  as a function of  $B(\tilde{t} \rightarrow b\ell\nu\chi^0) \times B(\tilde{t} \rightarrow b\bar{q}q'\chi^0)$ ,  $B(\tilde{t} \rightarrow c\chi^0)$  and  $m_{\chi^0}$ . Bound weakens for larger  $B(\tilde{t} \rightarrow c\chi^0)$  and  $m_{\chi^0}$ .
- 40 ABDALLAH 03C looked for events of the type  $q\bar{q}R^\pm R^\pm$ ,  $q\bar{q}R^\pm R^0$  or  $q\bar{q}R^0 R^0$  in  $e^+e^-$  interactions at  $\sqrt{s} = 189\text{--}208 \text{ GeV}$ . The  $R^\pm$  bound states are identified by

- anomalous  $dE/dx$  in the tracking chambers and the  $R^0$  by missing energy, due to their reduced energy loss in the calorimeters. Excluded mass regions in the  $(m(\tilde{t}), m(\tilde{g}))$  plane for  $m(\tilde{g}) > 2$  GeV are obtained for several values of the probability for the gluino to fragment into  $R^\pm$  or  $R^0$ , as shown in their Fig. 18. The limit improves to 90 GeV for  $\theta_t = 0$ .
- 41 ACOSTA 03C searched in  $107 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8$  TeV for pair production of  $\tilde{t}$  followed by the decay  $\tilde{t} \rightarrow b\ell\tilde{\nu}$ . They looked for events with two isolated leptons ( $e$  or  $\mu$ ), at least one jet and  $\cancel{E}_T$ . The excluded mass range is reduced for larger  $m_{\tilde{\nu}}$ , and no limit is set for  $m_{\tilde{\nu}} > 88.4$  GeV (see Fig. 2). Superseded by AALTONEN 10Y.
- 42 Theoretical analysis of  $e^+e^-+2$  jet final states from the RPV decay of  $\tilde{t}\tilde{t}^*$  pairs produced in  $p\bar{p}$  collisions at  $\sqrt{s}=1.8$  TeV. 95%CL limits of 220 (165) GeV are derived for  $B(\tilde{t} \rightarrow e q)=1$  (0.5).
- 43 HEISTER 03G searches for the production of  $\tilde{t}$  pairs in the case of  $\cancel{R}$  prompt decays with  $LL\bar{E}$ ,  $LQ\bar{D}$  or  $UDD$  couplings at  $\sqrt{s} = 189\text{--}209$  GeV. The limit holds for indirect decays mediated by  $\cancel{R}$   $UDD$  couplings. It improves to 91 GeV for indirect decays mediated by  $\cancel{R}$   $LL\bar{E}$  couplings, to 97 GeV for direct (assuming  $B(\tilde{t}_L \rightarrow q\tau) = 100\%$ ) and to 85 GeV for indirect decays mediated by  $\cancel{R}$   $LQ\bar{D}$  couplings. Supersedes the results from BARATE 01B.
- 44 HEISTER 03H use  $e^+e^-$  data from 183–208 GeV to look for the production of stop decaying into a  $c$  quark and a stable gluino hadronizing into charged or neutral R-hadrons. Combining these results with bounds on stable squarks and on a stable gluino LSP from the same paper yields the quoted limit. See their Fig. 13 for the dependence of the mass limit on the gluino mass and on  $\theta_t$ .
- 45 ABAZOV 02C looked in  $108.3\text{pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8$  TeV for events with  $e\mu\cancel{E}_T$ , originating from associated production  $\tilde{t}\tilde{t}$ . Branching ratios are assumed to be 100%. The bound for the  $b\ell\tilde{\nu}$  decay weakens for large  $\tilde{\nu}$  mass (see Fig. 3), and no limit is set when  $m_{\tilde{\nu}} > 85$  GeV. See Fig. 4 for the limits in case of decays to a real  $\tilde{\chi}_1^\pm$ , followed by  $\tilde{\chi}_1^\pm \rightarrow \ell\tilde{\nu}$ , as a function of  $m_{\tilde{\chi}_1^\pm}$ .
- 46 ACHARD 02 searches for the production of squarks in the case of  $\cancel{R}$  prompt decays with  $UDD$  couplings at  $\sqrt{s}=189\text{--}208$  GeV. The search is performed for direct and indirect decays, assuming one coupling at the time to be nonzero. The limit is computed for the minimal cross section and holds for both direct and indirect decays.
- 47 AFFOLDER 01B searches for decays of the top quark into stop and LSP, in  $t\bar{t}$  events. Limits on the stop mass as a function of the LSP mass and of the decay branching ratio are shown in Fig. 3. They exclude branching ratios in excess of 45% for SLP masses up to 40 GeV.
- 48 ABREU 00I searches for the production of stop in the case of  $R$ -parity violation with  $LL\bar{E}$  couplings, for which only indirect decays are allowed. They investigate topologies with jets plus leptons in data from  $\sqrt{s}=183$  GeV. The lower bound on the stop mass assumes a neutralino mass limit of 27 GeV, also derived in ABREU 00I.
- 49 AFFOLDER 00D search for final states with 2 or 3 jets and  $\cancel{E}_T$ , one jet with a  $c$  tag. See their Fig. 2 for the mass exclusion in the  $(m_{\tilde{t}}, m_{\tilde{\chi}_1^0})$  plane. The maximum excluded  $m_{\tilde{t}}$  value is 119 GeV, for  $m_{\tilde{\chi}_1^0}=40$  GeV.
- 50 AFFOLDER 00G searches for  $\tilde{t}_1\tilde{t}_1^*$  production, with  $\tilde{t}_1 \rightarrow b\ell\tilde{\nu}$ , leading to topologies with  $\geq 1$  isolated lepton ( $e$  or  $\mu$ ),  $\cancel{E}_T$ , and  $\geq 2$  jets with  $\geq 1$  tagged as  $b$  quark by a secondary vertex. See Fig. 4 for the excluded mass range as a function of  $m_{\tilde{\nu}}$ . Cross-section limits for  $\tilde{t}_1\tilde{t}_1^*$ , with  $\tilde{t}_1 \rightarrow b\chi_1^\pm$  ( $\chi_1^\pm \rightarrow \ell^\pm\nu\tilde{\chi}_1^0$ ), are given in Fig. 2. Superseded by AALTONEN 10Y.
- 51 ABE 99M looked in  $107 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8$  TeV for events with like sign dielectrons and two or more jets from the sequential decays  $\tilde{q} \rightarrow q\tilde{\chi}_1^0$  and  $\tilde{\chi}_1^0 \rightarrow e q\tilde{q}'$ , assuming  $\cancel{R}$  coupling  $L_1 Q_j D_k^C$ , with  $j=2,3$  and  $k=1,2,3$ . They assume  $B(\tilde{t}_1 \rightarrow c\tilde{\chi}_1^0)=1$ ,

- $B(\tilde{\chi}_1^0 \rightarrow e q \bar{q}')=0.25$  for both  $e^+$  and  $e^-$ , and  $m_{\tilde{\chi}_1^0} \geq m_{\tilde{t}_1}/2$ . The limit improves for heavier  $\tilde{\chi}_1^0$ .
- 52 AID 96 considers photoproduction of  $\tilde{t}\tilde{t}$  pairs, with 100%  $R$ -parity violating decays of  $\tilde{t}$  to  $e q$ , with  $q=d, s$ , or  $b$  quarks.
- 53 AID 96 considers production and decay of  $\tilde{t}$  via the  $R$ -parity violating coupling  $\lambda' L_1 Q_3 d_1^c$ .
- 54 CHO 96 studied the consistency among the  $B^0-\bar{B}^0$  mixing,  $\epsilon$  in  $K^0-\bar{K}^0$  mixing, and the measurements of  $V_{cb}, V_{ub}/V_{cb}$ . For the range  $25.5 \text{ GeV} < m_{\tilde{t}_1} < m_Z/2$  left by AKERS 94K for  $\theta_t = 0.98$ , and within the allowed range in  $M_2-\mu$  parameter space from chargino, neutralino searches by ACCIARRI 95E, they found the scalar top contribution to  $B^0-\bar{B}^0$  mixing and  $\epsilon$  to be too large if  $\tan\beta < 2$ . For more on their assumptions, see the paper and their reference 10.
- 55 BUSKULIC 95E looked for  $Z \rightarrow \tilde{t}\tilde{t}$ , where  $\tilde{t} \rightarrow c\chi_1^0$  and  $\chi_1^0$  decays via  $R$ -parity violating interactions into two leptons and a neutrino.
- 56 SHIRAI 94 bound assumes the cross section without the  $s$ -channel  $Z$ -exchange and the QCD correction, underestimating the cross section up to 20% and 30%, respectively. They assume  $m_c=1.5 \text{ GeV}$ .

### Heavy $\tilde{g}$ (Gluino) MASS LIMIT

For  $m_{\tilde{g}} > 60-70 \text{ GeV}$ , it is expected that gluinos would undergo a cascade decay via a number of neutralinos and/or charginos rather than undergo a direct decay to photinos as assumed by some papers. Limits obtained when direct decay is assumed are usually higher than limits when cascade decays are included. Limits made obsolete by the most recent analyses of  $p\bar{p}$  collisions can be found in previous Editions of this *Review*.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
>1100	95	1,2 AAD	13AV ATLS	$\geq 7\text{jets} + \cancel{E}_T, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} < 350 \text{ GeV}$
>1150	95	1,3 AAD	13AV ATLS	$\geq 7\text{jets} + \cancel{E}_T, \tilde{g} \rightarrow \tilde{t}\tilde{t}, \tilde{t} \rightarrow t\tilde{\chi}_1^0, m_{\tilde{t}} < 750 \text{ GeV}, m_{\tilde{\chi}_1^0} = 60 \text{ GeV}$
> 900	95	1,4 AAD	13AV ATLS	$\geq 7\text{jets} + \cancel{E}_T, \tilde{g} \rightarrow \tilde{t}\tilde{t}, \tilde{t} \rightarrow \bar{s}\bar{b}$ (RPV), $400 \text{ GeV} < m_{\tilde{t}} < 1000 \text{ GeV}$
>1000	95	1,5 AAD	13AV ATLS	$\geq 7\text{jets} + \cancel{E}_T, \tilde{g} \rightarrow \bar{q}q'\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W^\pm\tilde{\chi}_1^0, (m_{\tilde{\chi}_1^\pm}-m_{\tilde{\chi}_1^0})/(m_{\tilde{g}}-m_{\tilde{\chi}_1^0}) = 0.5, m_{\tilde{\chi}_1^0} < 200 \text{ GeV}$
>1100	95	1,6 AAD	13AV ATLS	$\geq 7\text{jets} + \cancel{E}_T, \text{mSUGRA/CMSSM, large } m_0$
> 700	95	7 CHATRCHYAN 13G	CMS	$0,1,2, \geq 3 \text{ } b\text{-jets} + \cancel{E}_T, \text{CMSSM}$
>1250	95	7 CHATRCHYAN 13G	CMS	$0,1,2, \geq 3 \text{ } b\text{-jets} + \cancel{E}_T, \text{CMSSM, } m_{\tilde{g}} = m_{\tilde{q}}$
>1300	95	8 CHATRCHYAN 13P	CMS	$1 \ell^\pm + \text{jets} + \cancel{E}_T, \text{CMSSM, } m_0 < 800 \text{ GeV}$
>1150	95	9 CHATRCHYAN 13R	CMS	$\geq 1 \tau + \text{jets} + \cancel{E}_T, \text{CMSSM, } m_0 < 440 \text{ GeV}$

>1125	95	10,11	CHATRCHYAN 13T	CMS	jets + $\cancel{E}_T$ , $\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} = 0$ GeV
<b>&gt; 950</b>	95	10,12	CHATRCHYAN 13T	CMS	jets + $\cancel{E}_T$ , $\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} = 0$ GeV
> 950	95	10,13	CHATRCHYAN 13T	CMS	jets + $\cancel{E}_T$ , $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} = 0$ GeV
>1000	95	14	CHATRCHYAN 13V	CMS	same-sign $\ell^\pm \ell^\pm + \geq 2$ $b$ -jets, $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} < 400$ GeV
> 550	95	15	AAD	12AP ATLS	$\ell^\pm \ell^\pm +$ jets + $\cancel{E}_T$ , $\tilde{g} \rightarrow \tilde{t}_1 t$ , CMSSM
> 820	95	16	AAD	12AX ATLS	$\ell +$ jets + $\cancel{E}_T$ , CMSSM, $m_{\tilde{g}} = m_{\tilde{q}}$
> 840	95	17	AAD	12BI ATLS	$\geq 6-9$ jets + $\cancel{E}_T$ , CMSSM, high $m_0$
>1020	95	18	AAD	12BY ATLS	$\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$ , $m_{\tilde{\chi}_1^0} < 400$ GeV
> 940	95	18	AAD	12BY ATLS	$\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ , $m_{\tilde{\chi}_1^0} < 50$ GeV
>1200	95	19	AAD	12CJ ATLS	$\ell^\pm +$ jets + $\cancel{E}_T$ , CMSSM, $m_{\tilde{g}} = m_{\tilde{q}}$
> 666	95	20	AAD	12CU ATLS	$\tilde{g} \rightarrow jjj, \cancel{E}$
<b>&gt; 800</b>	95	21	CHATRCHYAN 12AT	CMS	jets + $\cancel{E}_T$ , CMSSM
>1180	95	21	CHATRCHYAN 12AT	CMS	jets + $\cancel{E}_T$ , CMSSM, $m_{\tilde{g}} = m_{\tilde{q}}$
> 710	95	22	CHATRCHYAN 12U	CMS	$\ell^\pm \ell^\pm +$ jets + $\cancel{E}_T$ , CMSSM
> 700	95	23	AAD	11G ATLS	$\ell +$ jets + $\cancel{E}_T$ , $\tan\beta=3$ , $A_0=0$ , $\mu > 0$ , $m_{\tilde{g}} = m_{\tilde{q}}$
> 500	95	24	AAD	11N ATLS	jets + $\cancel{E}_T$ , degenerate $m_{\tilde{q}}$ of first two generations, $m_{\tilde{\chi}_1^0} = 0$ , all other supersymmetric particles heavy, any $m_{\tilde{q}}$
> 870	95	24	AAD	11N ATLS	jets + $\cancel{E}_T$ , degenerate $m_{\tilde{q}}$ of first two generations, $m_{\tilde{\chi}_1^0} = 0$ , all other supersymmetric particles heavy, $m_{\tilde{q}} = m_{\tilde{g}}$
> 775	95	24	AAD	11N ATLS	jets + $\cancel{E}_T$ , CMSSM, $m_{\tilde{q}} = m_{\tilde{g}}$
> 590	95	25	AAD	11O ATLS	$\tilde{g} \rightarrow \tilde{b}_1 b, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ , $m_{\tilde{\chi}_1^0} = 60$ GeV
> 500	95	26	CHATRCHYAN 11AC	CMS	jets + $\cancel{E}_T$ , CMSSM, $m_{\tilde{q}} < 1000$ GeV
> 280	95	27	AALTONEN	09S CDF	jets + $\cancel{E}_T$ , $\tan\beta=5$ , $\mu < 0$ , $A_0=0$ , any $m_{\tilde{q}}$
> 392	95	27	AALTONEN	09S CDF	jets + $\cancel{E}_T$ , $\tan\beta=5$ , $\mu < 0$ , $A_0=0$ , $m_{\tilde{q}} = m_{\tilde{g}}$
> 308	95	28	ABAZOV	08G D0	jets + $\cancel{E}_T$ , $\tan\beta=3$ , $\mu < 0$ , $A_0=0$ , any $m_{\tilde{q}}$
> 390	95	28	ABAZOV	08G D0	jets + $\cancel{E}_T$ , $\tan\beta=3$ , $\mu < 0$ , $A_0=0$ , $m_{\tilde{q}} = m_{\tilde{g}}$

> 270	95	29	ABULENCIA	06I	CDF	$\tilde{g} \rightarrow \tilde{b}b, \Delta m > 6 \text{ GeV}, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0, m_{\tilde{b}_1} < 220 \text{ GeV}$
> 195	95	30	AFFOLDER	02	CDF	Jets+ $\cancel{E}_T$ , any $m_{\tilde{q}}$
> 300	95	30	AFFOLDER	02	CDF	Jets+ $\cancel{E}_T, m_{\tilde{q}}=m_{\tilde{g}}$
> 129	95	31	ABBOTT	01D	D0	$\ell\ell$ +jets+ $\cancel{E}_T, \tan\beta < 10, m_0 < 300 \text{ GeV}, \mu < 0, A_0=0$
> 175	95	31	ABBOTT	01D	D0	$\ell\ell$ +jets+ $\cancel{E}_T, \tan\beta=2, \text{large } m_0, \mu < 0, A_0=0$
> 255	95	31	ABBOTT	01D	D0	$\ell\ell$ +jets+ $\cancel{E}_T, \tan\beta=2, m_{\tilde{g}}=m_{\tilde{q}}, \mu < 0, A_0=0$
> 168	95	32	AFFOLDER	01J	CDF	$\ell\ell$ +Jets+ $\cancel{E}_T, \tan\beta=2, \mu=-800 \text{ GeV}, m_{\tilde{q}} \gg m_{\tilde{g}}$
> 221	95	32	AFFOLDER	01J	CDF	$\ell\ell$ +Jets+ $\cancel{E}_T, \tan\beta=2, \mu=-800 \text{ GeV}, m_{\tilde{q}}=m_{\tilde{g}}$
> 190	95	33	ABBOTT	99L	D0	Jets+ $\cancel{E}_T, \tan\beta=2, \mu < 0, A=0$
> 260	95	33	ABBOTT	99L	D0	Jets+ $\cancel{E}_T, m_{\tilde{g}}=m_{\tilde{q}}$
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●						
>1360	95	34	AAD	13AB	ATLS	jets + 0,1,2 $\ell^\pm$ , CMSSM
> 900	95	35	AAD	13L	ATLS	jets + $\cancel{E}_T$ , CMSSM, $m_{\tilde{g}} = m_{\tilde{q}}$
> 900	95	36	AAD	13Q	ATLS	$\gamma$ + $b$ + $\cancel{E}_T$ , higgsino-like neutralino, $m_{\tilde{\chi}_1^0} > 220 \text{ GeV}$ , GMSB
>1170	95	37	CHATRCHYAN	13	CMS	$\ell^\pm \ell^\mp$ + jets + $\cancel{E}_T$ , CMSSM
>1020	95	38	CHATRCHYAN	13AK	CMS	$b$ -jets + $\cancel{E}_T, \tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} = 0 \text{ GeV}$
> 870	95	38	CHATRCHYAN	13AK	CMS	$b$ -jets + $\cancel{E}_T, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} = 0 \text{ GeV}$
> 700	95	39,40	CHATRCHYAN	13AM	CMS	$1\ell^\pm$ + $b$ -jets + $\cancel{E}_T, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} = 0 \text{ GeV}$
> 700	95	39,41	CHATRCHYAN	13AM	CMS	$1\ell^\pm$ + $b$ -jets + $\cancel{E}_T$ , CMSSM
>1000	95	42	CHATRCHYAN	13AO	CMS	$\ell^\pm \ell^\mp$ + jets + $\cancel{E}_T$ , CMSSM, $m_0 < 700 \text{ GeV}$
>1350	95	43	CHATRCHYAN	13AT	CMS	jets + $\cancel{E}_T, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$ simplified model, $m_{\tilde{\chi}_1^0} = 50 \text{ GeV}$
> 800	95	44	CHATRCHYAN	13AV	CMS	jets (+ leptons) + $\cancel{E}_T$ , CMSSM, $m_{\tilde{g}} = m_{\tilde{q}}$
>1000	95	45	CHATRCHYAN	13W	CMS	$\geq 1$ photons + jets + $\cancel{E}_T$ , GGM, wino-like NLSP, $m_{\tilde{\chi}_1^0} = 375 \text{ GeV}$
>1070	95	45	CHATRCHYAN	13W	CMS	$\geq 2$ photons + jets + $\cancel{E}_T$ , GGM, bino-like NLSP, $m_{\tilde{\chi}_1^0} = 375 \text{ GeV}$
>1070	95	46	AAD	12BA	CDF	$b$ -jets + $\cancel{E}_T$
> 950	95	47	AAD	12CP	ATLS	$2\gamma$ + $\cancel{E}_T$ , GMSB, bino NLSP, $m_{\tilde{\chi}_1^0} > 50 \text{ GeV}$
> 950	95	48	AAD	12W	ATLS	jets + $\cancel{E}_T$ , CMSSM, $m_{\tilde{g}} = m_{\tilde{q}}$

> 805	95	49 AAD	12X	ATLS	$2\gamma + \cancel{E}_T$ , GMSB, bino NLSP, $m_{\tilde{\chi}_1^0} > 50 \text{ GeV}$
		50 CHATRCHYAN 12		CMS	$e, \mu, \text{jets, razor, CMSSM}$
>1000	95	51 CHATRCHYAN 12AE		CMS	$\text{jets} + \cancel{E}_T, \tilde{g} \rightarrow qq\tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} <$ $200 \text{ GeV}$
		52 CHATRCHYAN 12AH		CMS	$b\text{-jets} + \cancel{E}_T, \tilde{g} \rightarrow bb\tilde{\chi}_1^0$
		52 CHATRCHYAN 12AH		CMS	$b\text{-jets} + \cancel{E}_T, \tilde{g} \rightarrow tt\tilde{\chi}_1^0$
		53 CHATRCHYAN 12AI		CMS	$\ell^\pm \ell^\pm + b\text{-jets} + \cancel{E}_T$
		54 CHATRCHYAN 12AL		CMS	$\geq 3\ell^\pm, \cancel{R}$
none 280–460	95	55 CHATRCHYAN 12BD		CMS	$\tilde{g} \rightarrow jjj, \cancel{R}$
		56 CHATRCHYAN 12Q		CMS	$\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow Z\tilde{\chi}_1^0$
> 500	95	57 DREINER	12A	THEO	$m_{\tilde{g}} \sim m_{\tilde{\chi}_1^0}$
> 650	95	58 DREINER	12A	THEO	$m_{\tilde{g}} = m_{\tilde{q}} \sim m_{\tilde{\chi}_1^0}$
> 520	95	59 AAD	11AF	ATLS	$\geq 6 \text{ jets} + \cancel{E}_T$ , CMSSM
> 560	95	60 AAD	11X	ATLS	$\tilde{g} \rightarrow \tilde{\chi}_1^0 X \rightarrow \gamma \tilde{G} X$
> 155	95	61 AALTONEN	11Q	CDF	$\cancel{R}, \overline{UD}D, m_{\tilde{q}} = m_{\tilde{g}} + 10 \text{ GeV}$
		62 CHATRCHYAN 11AB		CMS	$\ell^\pm \ell^\pm$
		63 CHATRCHYAN 11G		CMS	$\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$
		64 CHATRCHYAN 11Q		CMS	$\ell + \text{jets} + \cancel{E}_T$
>1040	95	65 CHATRCHYAN 11V		CMS	GMSB scenario, $\bar{\ell}$ co-NLSP
		66 CHATRCHYAN 11W		CMS	$\text{jets} + \cancel{E}_T$ , CMSSM
		67 KHACHATRY...11I		CMS	$\text{jets} + \cancel{E}_T$
> 224	95	68 ABAZOV	02F	D0	$\cancel{R} \lambda'_{2jk}$ indirect decays, $\tan\beta=2$ , any $m_{\tilde{q}}$
> 265	95	68 ABAZOV	02F	D0	$\cancel{R} \lambda'_{2jk}$ indirect decays, $\tan\beta=2$ , $m_{\tilde{q}} = m_{\tilde{g}}$
		69 ABAZOV	02G	D0	$p\bar{p} \rightarrow \tilde{g}\tilde{g}, \tilde{g}\tilde{q}$
		70 CHEUNG	02B	THEO	
		71 BERGER	01	THEO	$p\bar{p} \rightarrow X + b\text{-quark}$
> 240	95	72 ABBOTT	99	D0	$\tilde{g} \rightarrow \tilde{\chi}_2^0 X \rightarrow \tilde{\chi}_1^0 \gamma X, m_{\tilde{\chi}_2^0} -$ $m_{\tilde{\chi}_1^0} > 20 \text{ GeV}$
> 320	95	72 ABBOTT	99	D0	$\tilde{g} \rightarrow \tilde{\chi}_1^0 X \rightarrow \tilde{G} \gamma X$
> 227	95	73 ABBOTT	99K	D0	any $m_{\tilde{q}}, \cancel{R}, \tan\beta=2, \mu < 0$
> 212	95	74 ABACHI	95C	D0	$m_{\tilde{g}} \geq m_{\tilde{q}}$ ; with cascade decays
> 144	95	74 ABACHI	95C	D0	Any $m_{\tilde{q}}$ ; with cascade decays
		75 ABE	95T	CDF	$\tilde{g} \rightarrow \tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \gamma$
		76 HEBBEKER	93	RVUE	$e^+e^-$ jet analyses
> 218	90	77 ABE	92L	CDF	$m_{\tilde{q}} \leq m_{\tilde{g}}$ ; with cascade decay
> 100		78 ROY	92	RVUE	$p\bar{p} \rightarrow \tilde{g}\tilde{g}; \cancel{R}$
		79 NOJIRI	91	COSM	
none 4–53	90	80 ALBAJAR	87D	UA1	Any $m_{\tilde{q}} > m_{\tilde{g}}$
none 4–75	90	80 ALBAJAR	87D	UA1	$m_{\tilde{q}} = m_{\tilde{g}}$
none 16–58	90	81 ANSARI	87D	UA2	$m_{\tilde{q}} \lesssim 100 \text{ GeV}$

- <sup>1</sup> AAD 13AV searched in  $20.3 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  for events containing large number of jets (7 or more), with missing transverse momentum and no isolated electrons or muons. The sensitivity of the search is enhanced by considering the number of  $b$ -tagged jets and the scalar sum of masses of large-radius jets in an event. No evidence was found for excesses above the expected level of Standard Model background.
- <sup>2</sup> Exclusion limits at 95% C.L. are set on the gluino mass assuming the gluino decays exclusively via an off-shell top squark,  $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ , see their Fig. 9.
- <sup>3</sup> Exclusion limits at 95% C.L. are set on the gluino mass assuming the gluino decays exclusively via an on-shell top quark,  $\tilde{g} \rightarrow \tilde{t}\bar{t}$ , with consecutively  $\tilde{t} \rightarrow t\tilde{\chi}_1^0$ , assuming  $m_{\tilde{\chi}_1^0} = 60 \text{ GeV}$ , see their Fig. 10.
- <sup>4</sup> Exclusion limits at 95% C.L. are set on the gluino mass assuming the gluino decays exclusively via an on-shell top quark,  $\tilde{g} \rightarrow \tilde{t}\bar{t}$ , with the stop consecutively decaying via the R-parity- and baryon-number-violating decay  $\tilde{t} \rightarrow \bar{s}b$ , see Fig. 14.
- <sup>5</sup> Exclusion limits at 95% C.L. are set on the gluino mass assuming the gluino decays exclusively via an on-shell quark and a chargino,  $\tilde{g} \rightarrow \bar{q}q'\tilde{\chi}_1^\pm$ , with consecutively  $\tilde{\chi}_1^\pm \rightarrow W^\pm\tilde{\chi}_1^0$ , see their Fig. 11. An alternative interpretation in the case where the gluino can decay via  $\tilde{\chi}_1^\pm$  or  $\tilde{\chi}_2^0$  is given in Fig. 12.
- <sup>6</sup> Exclusion limits at the 95% C.L. are derived in the mSUGRA/CMSSM model with parameters  $\tan\beta = 30$ ,  $A_0 = -2 m_0$  and  $\mu > 0$ , see their Fig. 13. For large universal scalar masses  $m_0$ , gluino masses smaller than 1.1 TeV are excluded at 95% C.L.
- <sup>7</sup> CHATRCHYAN 13G searched in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for the production of squarks and gluinos in events containing 0,1,2,  $\geq 3$   $b$ -jets, missing transverse momentum and no electrons or muons. No excess over the expected SM background is observed. In mSUGRA/CMSSM models with  $\tan\beta = 10$ ,  $A_0 = 0$ , and  $\mu > 0$ , gluinos with masses below 700 GeV are excluded at 95% C.L. Squarks and gluinos of equal mass are excluded for masses below 1250 GeV at 95% C.L. Exclusions are also derived in various simplified models, see Fig. 7.
- <sup>8</sup> CHATRCHYAN 13P searched in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing a single isolated electron or muon, energetic jets and large  $\cancel{E}_T$ . No significant excesses over the expected SM backgrounds are observed and 95% C.L. exclusion limits are derived in the mSUGRA/CMSSM model with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , see Fig. 17. The results are also interpreted in a simplified model, see Fig. 19.
- <sup>9</sup> CHATRCHYAN 13R searched in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing one or more hadronically decaying  $\tau$  leptons, energetic jets and large  $\cancel{E}_T$ . No significant excesses over the expected SM backgrounds are observed and 95% C.L. exclusion limits are derived in the mSUGRA/CMSSM model with  $\tan\beta=10$ ,  $A_0=0$  and  $\mu > 0$ , see Fig. 7. The results are also interpreted in various simplified models, see Fig. 9.
- <sup>10</sup> CHATRCHYAN 13T searched in  $11.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  for events with at least two energetic jets and significant  $\cancel{E}_T$ , using the  $\alpha_T$  variable to discriminate between processes with genuine and misreconstructed  $\cancel{E}_T$ . No significant excess above the Standard Model expectations is observed.
- <sup>11</sup> CHATRCHYAN 13T limits are set on gluino masses in simplified models where the decay  $\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$  takes place with a branching ratio of 100%, see Fig. 8 and Table 9.
- <sup>12</sup> CHATRCHYAN 13T limits are set on gluino masses in simplified models where the decay  $\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$  takes place with a branching ratio of 100%, see Fig. 8 and Table 9.
- <sup>13</sup> CHATRCHYAN 13T limits are set on gluino masses in simplified models where the decay  $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$  takes place with a branching ratio of 100%, see Fig. 8 and Table 9.

- 14 CHATRCHYAN 13V searched in  $10.5 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  for events with two isolated same-sign dileptons and at least two  $b$ -jets in the final state. No significant excess above the Standard Model expectations is observed. Limits are set on the gluino mass in simplified models where the decay  $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$  takes place with a branching ratio of 100%, or where the decay  $\tilde{g} \rightarrow \tilde{t}t, \tilde{t} \rightarrow t\tilde{\chi}_1^0$  takes place with a branching ratio of 100%, with varying mass of the  $\tilde{\chi}_1^0$ , or where the decay  $\tilde{g} \rightarrow \tilde{b}b, \tilde{b} \rightarrow t\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W^\pm\tilde{\chi}_1^0$  takes place with a branching ratio of 100%, with varying mass of the  $\tilde{\chi}_1^\pm$ , see Fig. 4.
- 15 AAD 12AP searched in  $2.05 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for gluinos decaying via the scalar partner of the top quark into events with two same-sign leptons, jets and missing transverse energy. No excess beyond the Standard Model expectation is observed. Exclusion limits are derived in the mSUGRA/CMSSM model with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , see Fig. 4, and in simplified models, see Figs. 2 and 3.
- 16 AAD 12AX searched in  $1.04 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for supersymmetry in events containing jets, missing transverse momentum and one isolated electron or muon. No excess over the expected SM background is observed and model-independent limits are set on the cross section of new physics contributions to the signal regions. In mSUGRA/CMSSM models with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , squarks and gluinos of equal mass are excluded for masses below 820 GeV at 95% C.L. Limits are also set on simplified models for gluino production and decay via an intermediate chargino and on supersymmetric models with bilinear  $R$ -parity violation. Supersedes AAD 11G.
- 17 AAD 12BI looked in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 6$  to  $\geq 9$  jets plus  $\cancel{E}_T$ . No excess over the expected background is observed. Limits are derived in the CMSSM ( $m_0, m_{1/2}$ ) plane for  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , see their Fig. 7. Limits are also set in the  $(m_{\tilde{g}}, m_{\tilde{\chi}_1^0})$  plane in a simplified supersymmetric model with four tops +  $\cancel{E}_T$  in the final state. Supersedes AAD 11AF.
- 18 AAD 12BY searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with large missing transverse momentum and at least three  $b$ -jets in the final state. The data are found to be consistent with the Standard Model expectations. In a simplified supersymmetric scenario where  $\tilde{g} \rightarrow \tilde{b}_1 b$  and  $\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ , with branching ratios of 100% for both decays, a 95% C.L. limit on the gluino mass of 1000 GeV is set for  $m_{\tilde{b}_1} < 870 \text{ GeV}$  and  $m_{\tilde{\chi}_1^0} = 60 \text{ GeV}$ . In a scenario where the sbottom is heavier than the gluino and the gluino decays through a three-body decay into bottom quarks 100% of the time,  $\tilde{g} \rightarrow b b \tilde{\chi}_1^0$ , the limit on the gluino mass becomes 1020 GeV, provided  $m_{\tilde{\chi}_1^0} < 400 \text{ GeV}$ . In a scenario where  $\tilde{g} \rightarrow \tilde{t}_1 t$  and  $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$ , with branching ratios of 100% for both decays, a 95% C.L. limit on the gluino mass of 820 GeV is set for  $m_{\tilde{t}_1} < 640 \text{ GeV}$  and  $m_{\tilde{\chi}_1^0} = 60 \text{ GeV}$ . In a scenario where the stop is heavier than the gluino and the gluino decays through a three-body decay into top quarks 100% of the time,  $\tilde{g} \rightarrow t t \tilde{\chi}_1^0$ , the limit on the gluino mass becomes 940 GeV, provided  $m_{\tilde{\chi}_1^0} < 50 \text{ GeV}$ .
- 19 AAD 12CJ searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing one or more isolated leptons (electrons or muons), jets and  $\cancel{E}_T$ . The observations are in good agreement with the SM expectations and exclusion limits have been set in number of SUSY models. In the mSUGRA/CMSSM model with  $\tan\beta = 10$ ,  $A_0 = 0$ , and  $\mu > 0$ , 95% C.L. exclusion limits have been derived for  $m_{\tilde{g}} < 1200 \text{ GeV}$ , assuming equal squark and gluino masses. In minimal GMSB, values of the effective SUSY breaking scale  $\Lambda < 50 \text{ TeV}$  are excluded at 95% C.L. for  $\tan\beta < 45$ . Also exclusion limits in a number of simplified models have been presented, see Figs. 10 and 11.

- 20 AAD 12CU searched in  $4.6 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for pair production of gluinos decaying into six-quark final states in an  $R$ -parity violating supersymmetric model. The data are found to be consistent with the Standard Model expectations. Based on an analysis where all six jets in the final state are resolved, a 95% C.L. limit of 666 GeV is placed on the gluino mass. The gluino decay is assumed to be prompt.
- 21 CHATRCHYAN 12AT searched in  $4.73 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for the production of squarks and gluinos in events containing jets, missing transverse momentum and no electrons or muons. No excess over the expected SM background is observed. In mSUGRA/CMSSM models with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , gluinos with masses below 800 GeV are excluded at 95% C.L. Squarks and gluinos of equal mass are excluded for masses below 1180 GeV at 95% C.L. Exclusions are also derived in various simplified models, see Fig. 6.
- 22 CHATRCHYAN 12U looked in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with two same-sign leptons ( $e, \mu, \tau$ ) not necessarily the same flavor, jets and missing transverse energy. No excess beyond the Standard Model expectation is observed. Exclusion limits are derived in the mSUGRA/CMSSM model with  $\tan\beta = 10$ ,  $A_0 = 0$ , and  $\mu > 0$ , see Fig. 3. The limit is independent of the squark masses. The exclusion includes a  $-1 \sigma_{th}$  reduction to account for the theory uncertainty on the cross section.
- 23 AAD 11G looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with a single lepton ( $e$  or  $\mu$ ), jets and  $\cancel{E}_T$  from the production of squarks and gluinos. No evidence for an excess over the SM expectation is observed, and a limit is derived in the CMSSM ( $m_0, m_{1/2}$ ) plane for  $\tan\beta = 3$ , see Fig. 2.
- 24 AAD 11N looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 2$  jets and  $\cancel{E}_T$ . Four signal regions were defined, and the background model was found to be in good agreement with the data. Limits are derived in the  $(m_{\tilde{g}}, m_{\tilde{q}})$  plane (see Fig. 2) for a simplified model where degenerate masses of the squarks of the first two generations are assumed,  $m_{\tilde{\chi}_1^0} = 0$ , and all other masses including third generation squarks are set to 5 TeV. Limits are also derived in the CMSSM ( $m_0, m_{1/2}$ ) plane (see Fig. 3) for  $\tan\beta = 3$ .
- 25 AAD 11O looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with jets, of which at least one is a  $b$ -jet, and  $\cancel{E}_T$ . No excess above the Standard Model was found. Limits are derived in the  $(m_{\tilde{g}}, m_{\tilde{b}_1})$  plane (see Fig. 2) under the assumption of 100% branching ratios and  $\tilde{b}_1$  being the lightest squark. The quoted limit is valid for  $m_{\tilde{b}_1} < 500 \text{ GeV}$ . A similar approach for  $\tilde{t}_1$  as the lightest squark with  $\tilde{g} \rightarrow \tilde{t}_1 t$  and  $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$  with 100% branching ratios leads to a gluino mass limit of 520 GeV for  $130 < m_{\tilde{t}_1} < 300 \text{ GeV}$ . Limits are also derived in the CMSSM ( $m_0, m_{1/2}$ ) plane for  $\tan\beta = 40$ , see Fig. 4, and in scenarios based on the gauge group SO(10).
- 26 CHATRCHYAN 11AC looked in  $36 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 3$  jets, a large total transverse energy, and  $\cancel{E}_T$ . No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM ( $m_0, m_{1/2}$ ) plane and the  $(m_{\tilde{g}}, m_{\tilde{q}})$  plane for  $\tan\beta = 10$  (see Fig. 10). Limits are also obtained for Simplified Model Spectra.
- 27 AALTONEN 09S searched in  $2 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with at least 2 jets and  $\cancel{E}_T$ . No evidence for a signal is observed. A limit is derived for a mSUGRA scenario in the  $m_{\tilde{q}}$  versus  $m_{\tilde{g}}$  plane, see their Fig. 2.
- 28 ABAZOV 08G looked in  $2.1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.96 \text{ TeV}$  for events with acoplanar jets or multijets with large  $\cancel{E}_T$ . No significant excess was found compared to the background expectation. A limit is derived on the masses of squarks and gluinos for specific MSUGRA parameter values, see Figure 3. Similar results would be obtained for a large class of parameter sets. Supersedes the results of ABAZOV 06C.
- 29 ABULENCIA 06I searched in  $156 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for multijet events with large  $\cancel{E}_T$ . They request at least 2  $b$ -tagged jets and no isolated leptons.

- They investigate the production of gluinos decaying into  $\tilde{b}_1 b$  followed by  $\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ . Both branching fractions are assumed to be 100% and the LSP mass to be 60 GeV. No significant excess was found compared to the background expectation. Upper limits on the cross-section are extracted and a limit is derived on the masses of sbottom and gluinos, see their Fig.3.
- 30 AFFOLDER 02 searched in  $\sim 84 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions for events with  $\geq 3$  jets and  $\cancel{E}_T$ , arising from the production of gluinos and/or squarks. Limits are derived by scanning the parameter space, for  $m_{\tilde{q}} \geq m_{\tilde{g}}$  in the framework of minimal Supergravity, assuming five flavors of degenerate squarks, and for  $m_{\tilde{q}} < m_{\tilde{g}}$  in the framework of constrained MSSM, assuming conservatively four flavors of degenerate squarks. See Fig. 3 for the variation of the limit as function of the squark mass. Supersedes the results of ABE 97K.
- 31 ABBOTT 01D looked in  $\sim 108 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8$  TeV for events with  $e e$ ,  $\mu\mu$ , or  $e\mu$  accompanied by at least 2 jets and  $\cancel{E}_T$ . Excluded regions are obtained in the MSUGRA framework from a scan over the parameters  $0 < m_0 < 300$  GeV,  $10 < m_{1/2} < 110$  GeV, and  $1.2 < \tan\beta < 10$ .
- 32 AFFOLDER 01J searched in  $\sim 106 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions for events with 2 like-sign leptons ( $e$  or  $\mu$ ),  $\geq 2$  jets and  $\cancel{E}_T$ , expected to arise from the production of gluinos and/or squarks with cascade decays into  $\tilde{\chi}^\pm$  or  $\tilde{\chi}_2^0$ . Spectra and decay rates are evaluated in the framework of minimal Supergravity, assuming five flavors of degenerate squarks and a pseudoscalar Higgs mass  $m_A=500$  GeV. The limits are derived for  $\tan\beta=2$ ,  $\mu=-800$  GeV, and scanning over  $m_{\tilde{g}}$  and  $m_{\tilde{q}}$ . See Fig. 2 for the variation of the limit as function of the squark mass. These limits supersede the results of ABE 96D.
- 33 ABBOTT 99L consider events with three or more jets and large  $\cancel{E}_T$ . Spectra and decay rates are evaluated in the framework of minimal Supergravity, assuming five flavors of degenerate squarks, and scanning the space of the universal gaugino ( $m_{1/2}$ ) and scalar ( $m_0$ ) masses See their Figs. 2–3 for the dependence of the limit on the relative value of  $m_{\tilde{q}}$  and  $m_{\tilde{g}}$ .
- 34 AAD 13AB searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for the production of squarks and gluinos in events containing 0,1 or 2 high- $p_T$  leptons and with or without jets identified as originating from  $b$ -quarks. No excess over the expected SM background is observed. Limits are derived in mSUGRA/CMSSM models with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , see their Fig. 12. Also, exclusion limits in simplified models containing gluinos, squarks, charginos, and stops are set, see their Figures 10 and 11.
- 35 AAD 13L searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for the production of squarks and gluinos in events containing jets, missing transverse momentum and no high- $p_T$  electrons or muons. No excess over the expected SM background is observed. In mSUGRA/CMSSM models with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , squarks and gluinos of equal mass are excluded for masses below 1360 GeV at 95% C.L. In a simplified model containing only squarks of the first two generations, a gluino octet and a massless neutralino, gluino masses below 860 GeV are excluded at 95% C.L. for squark masses below 2 TeV. See their Figures 10–15 for more precise bounds.
- 36 AAD 13Q searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for events containing a high- $p_T$  isolated photon, at least one jet identified as originating from a bottom quark, and high missing transverse momentum. Such signatures may originate from supersymmetric models with gauge-mediated supersymmetry breaking in events in which one of a pair of higgsino-like neutralinos decays into a photon and a gravitino while the other decays into a Higgs boson and a gravitino. No significant excess above the expected background was found and limits were set on the gluino mass as a function of the neutralino mass in a generalized GMSB model (GGM) with a higgsino-like neutralino NLSP, see their Fig. 4. For neutralino masses greater than 220 GeV, gluino masses below 900 GeV are excluded at 95% C.L.
- 37 CHATRCHYAN 13 looked in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7$  TeV for events with two opposite-sign leptons ( $e$ ,  $\mu$ ,  $\tau$ ), jets and missing transverse energy. No excess beyond the Standard Model expectation is observed. Exclusion limits are derived in the mSUGRA/CMSSM model with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , see Fig. 6.

- 38 CHATRCHYAN 13AK searched in  $19.4 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  for events with large  $\cancel{E}_T$ , no isolated electron or muon, and at least three jets with one or more identified as a  $b$ -quark jet. No significant excesses over the expected SM backgrounds are observed and 95% C.L. limits on the production cross section of gluinos are set as a function of the gluino and neutralino mass in a scenario where  $\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$  with a 100% branching ratio, see Fig. 7 left, and in a scenario where  $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$  with a 100% branching ratio, see Fig. 7, right. Supersedes CHATRCHYAN 12AH.
- 39 CHATRCHYAN 13AM searched in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with large  $\cancel{E}_T$ , a single isolated electron or muon, and multiple jets including some identified as a  $b$ -quark jet. No significant excesses over the expected SM backgrounds are observed and 95% C.L.
- 40 CHATRCHYAN 13AM limits on the production cross section of gluinos are set as a function of the gluino and neutralino mass in a scenario where  $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$  with a 100% branching ratio, see Fig. 10.
- 41 CHATRCHYAN 13AM exclusion limits are derived in the mSUGRA/CMSSM model with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , see Fig. 8.
- 42 CHATRCHYAN 2013AO searched in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with two opposite-sign isolated leptons accompanied by hadronic jets and  $\cancel{E}_T$ . No significant excesses over the expected SM backgrounds are observed and 95% C.L. exclusion limits are derived in the mSUGRA/CMSSM model with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , see Fig. 8. The results are also interpreted in an asymmetric simplified model of gluino pair production, see Fig. 7.
- 43 CHATRCHYAN 13AT provides interpretations of various searches for supersymmetry by the CMS experiment based on  $4.73\text{--}4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  in the framework of simplified models. Limits are set on the gluino mass in a simplified models where gluinos are pair-produced and the decay  $\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$  takes place with a branching ratio of 100%, see Fig. 4.
- 44 CHATRCHYAN 13AV searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for new heavy particle pairs decaying into jets (possibly  $b$ -tagged), leptons and  $\cancel{E}_T$  using the Razor variables. No significant excesses over the expected SM backgrounds are observed and 95% C.L. exclusion limits are derived in the mSUGRA/CMSSM model with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , see Fig. 3. The results are also interpreted in various simplified models, see Fig. 4.
- 45 CHATRCHYAN 13W searched in  $4.93 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with one or more photons, hadronic jets and  $\cancel{E}_T$ . No significant excess above the Standard Model expectations is observed. Limits are set on gluino masses in the general gauge-mediated SUSY breaking model (GGM), for both a wino-like and bino-like neutralino NLSP scenario, see Fig. 5.
- 46 AAD 12BA searched in  $2.05 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with heavy flavor jets and large  $\cancel{E}_T$  due to  $\tilde{g} \rightarrow \tilde{t}_1 b$  or  $\tilde{g} \rightarrow \tilde{t}_1 t$  decays. No significant excess above the expected background was found and limits were set on the gluino mass in simplified  $R$ -parity conserving models in which only scalar bottoms and tops appear in the gluino decay and in an  $\text{SO}(10)$  model framework.
- 47 AAD 12CP searched in  $4.8 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with two photons and large  $\cancel{E}_T$  due to  $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$  decays in a GMSB framework. No significant excess above the expected background was found and limits were set on the gluino mass as a function of the neutralino mass in a generalized GMSB model (GGM) with a bino-like neutralino NLSP. The other sparticle masses were decoupled,  $\tan\beta = 2$  and  $c\tau_{NLSP} < 0.1 \text{ mm}$ . Also, in the framework of the SPS8 model, a 95% C.L. lower limit was set on the breaking scale  $\Lambda$  of 196 TeV.

- 48 AAD 12W searched in  $1.04 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for the production of squarks and gluinos in events containing jets, missing transverse momentum and no electrons or muons. No excess over the expected SM background is observed. In mSUGRA/CMSSM models with  $\tan\beta = 10$ ,  $A_0 = 0$  and  $\mu > 0$ , squarks and gluinos of equal mass are excluded for masses below 950 GeV at 95% C.L. In a simplified model containing only squarks of the first two generations, a gluino octet and a massless neutralino, gluino masses below 700 GeV are excluded at 95% C.L.
- 49 AAD 12X searched in  $1.07 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with two photons and large  $\cancel{E}_T$  due to  $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$  decays in a GMSB framework. No significant excess above the expected background was found and limits were set on the gluino mass as a function of the neutralino mass in a generalized GMSB model (GGM) with a bino-like neutralino NLSP. The other sparticle masses were set to 1.5 TeV,  $\tan\beta = 2$  and  $c\tau_{NLSP} < 0.1 \text{ mm}$ . Also, in the framework of the SPS8 model, a 95% C.L. lower limit was set on the breaking scale  $\Lambda$  of 145 TeV. Superseded by AAD 12CP.
- 50 CHATRCHYAN 12 looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $e$  and/or  $\mu$  and/or jets, a large total transverse energy, and  $\cancel{E}_T$ . The event selection is based on the dimensionless razor variable  $R$ , related to the  $\cancel{E}_T$  and  $M_R$ , an indicator of the heavy particle mass scale. No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM  $(m_0, m_{1/2})$  plane for  $\tan\beta = 3, 10$  and 50 (see Fig. 7 and 8). Limits are also obtained for Simplified Model Spectra.
- 51 CHATRCHYAN 12AE searched in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with at least three jets and large missing transverse momentum. No significant excesses over the expected SM backgrounds are observed and 95% C.L. limits on the production cross section of gluinos in a scenario where  $\tilde{g} \rightarrow qq\tilde{\chi}_1^0$  with a 100% branching ratio, see Fig. 3. For  $m_{\tilde{\chi}_1^0} < 200 \text{ GeV}$ , values of  $m_{\tilde{g}}$  below 1000 GeV are excluded at 95% C.L. Also limits in the CMSSM are presented, see Fig. 2.
- 52 CHATRCHYAN 12AH searched in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with large  $\cancel{E}_T$ , at least three jets, and at least one, two or three  $b$ -quark jets. No significant excesses over the expected SM backgrounds are observed and 95% C.L. limits on the production cross section of gluinos are set as a function of the gluino and neutralino mass in a scenario where  $\tilde{g} \rightarrow bb\tilde{\chi}_1^0$  with a 100% branching ratio, see Fig. 14, and in a scenario where  $\tilde{g} \rightarrow tt\tilde{\chi}_1^0$  with a 100% branching ratio, see Fig. 15.
- 53 CHATRCHYAN 12AI looked in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with two same-sign leptons ( $e, \mu$ ), but not necessarily same flavor, at least 2  $b$ -jets and missing transverse energy. No excess beyond the Standard Model expectation is observed. Exclusion limits are derived in simplified models where gluinos are pair produced and decay through  $\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$  (intermediate stop, real or virtual), see Fig. 6, or through  $\tilde{g} \rightarrow b\bar{b}W^+\tilde{\chi}_1^0$  (intermediate sbottom), see Fig. 8.
- 54 CHATRCHYAN 12AL looked in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for anomalous production of events with three or more isolated leptons. Limits on squark and gluino masses are set in  $\overline{R}$  SUSY models with leptonic  $LL\bar{E}$  couplings,  $\lambda_{123} > 0.05$ , and hadronic  $\overline{UDD}$  couplings,  $\lambda''_{112} > 0.05$ , see their Fig. 5. In the  $\overline{UDD}$  case the leptons arise from supersymmetric cascade decays. A very specific supersymmetric spectrum is assumed. All decays are prompt.
- 55 CHATRCHYAN 12BD searched in  $5.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for three-jet resonances produced in the decay of a gluino in  $R$ -parity violating supersymmetric models. No excess over the expected SM background is observed. Assuming a branching ratio for gluino decay into three jets of 100%, limits are set on the cross section of gluino pair production, see Fig. 4. Gluino masses between 280 GeV and 460 GeV are excluded at 95% C.L.

- 56 CHATRCHYAN 12Q looked in  $4.98 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for anomalous production of events with a  $Z$ -boson, jets and significant  $\cancel{E}_T$ . No evidence for an excess over the expected background is observed. Limits are set in a simplified supersymmetric model where the  $\tilde{\chi}_2^0 \rightarrow Z\tilde{\chi}_1^0$  decay is dominant, see Figs. 5 and 6.
- 57 DREINER 12A reassesses constraints from CMS (at  $7 \text{ TeV}$ ,  $\sim 4.4 \text{ fb}^{-1}$ ) under the assumption that the gluino and the lightest SUSY particle are quasi-degenerate in mass (compressed spectrum).
- 58 DREINER 12A reassesses constraints from CMS (at  $7 \text{ TeV}$ ,  $\sim 4.4 \text{ fb}^{-1}$ ) under the assumption that the first and second generation squarks, the gluino, and the lightest SUSY particle are quasi-degenerate in mass (compressed spectrum).
- 59 AAD 11AF looked in  $1.34 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with 6 up to 8 jets and  $\cancel{E}_T$ . No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM ( $m_0, m_{1/2}$ ) plane for  $\tan\beta = 10$  (see Fig. 5). The limit improves to  $m_{\tilde{g}} > 680 \text{ GeV}$  for  $m_{\tilde{q}} = 2 m_{\tilde{g}}$ .
- 60 AAD 11X looked in  $36 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 2$  photons and  $\cancel{E}_T$  from the pair production of gluinos with cascade decays to  $\tilde{\chi}_1^0$  followed by  $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$  prompt decay. No evidence for an excess over the SM expectation is observed, and a limit on the number of new physics events is set. Limits are derived in a Generalized Gauge Mediated model in the ( $m_{\tilde{g}}, m_{\tilde{\chi}_1^0}$ ) plane (see Fig. 5) under the assumptions  $\tan\beta = 2$  and all sparticle masses at  $1.5 \text{ TeV}$ , except the  $\tilde{g}$ ,  $\tilde{\chi}_1^0$ , and  $\tilde{G}$ . Superseded by AAD 12X.
- 61 AALTONEN 11Q searched in  $3.2 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with at least 6 jets from the pair production of gluinos and squarks with the subsequent decays  $\tilde{g} \rightarrow 3 \text{ jets}$  in the MSSM framework with  $\cancel{E}_T$ . No statistically significant bumps in the 3-jet systems are observed over the SM background. Limits on the cross section times branching ratio are derived as a function of the gluino mass, displayed in Fig. 3. For decoupled squarks in the range  $0.5 < m_{\tilde{q}} < 0.7 \text{ TeV}$  gluinos are excluded below  $144 \text{ GeV}$ . The quoted limit is for near degeneracy of squark and gluino masses.
- 62 CHATRCHYAN 11AB looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 2$  same charge isolated leptons ( $e, \mu$  or  $\tau$ ), jets and  $\cancel{E}_T$ . Such events might be produced from  $\tilde{g}\tilde{g}$  or  $\tilde{g}\tilde{q}$  decaying via charginos into leptons. No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM ( $m_0, m_{1/2}$ ) plane for  $\tan\beta = 3$  (see Fig. 10).
- 63 CHATRCHYAN 11G looked in  $36 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 2$  isolated photons,  $\geq 1$  jet and  $\cancel{E}_T$ , which may arise in a generalized gauge mediated model from the decay of a  $\tilde{\chi}_1^0$  NLSP. No evidence for an excess over the expected background is observed. Limits are derived in the plane of squark versus gluino mass (see Fig. 4) for several values of  $m_{\tilde{\chi}_1^0}$ .
- 64 CHATRCHYAN 11Q looked in  $36 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with a single isolated lepton ( $e$  or  $\mu$ ),  $\geq 4$  jets and  $\cancel{E}_T$ . No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM ( $m_0, m_{1/2}$ ) plane for  $\tan\beta = 10$  (see Fig. 7).
- 65 CHATRCHYAN 11V looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 3$  isolated leptons ( $e, \mu$  or  $\tau$ ), with or without jets and  $\cancel{E}_T$ . Multi-lepton final states originate from  $\tilde{q} \rightarrow \tilde{\chi}^0 + X$ , followed by  $\tilde{\chi}^0 \rightarrow \tilde{\ell}^\pm \ell^\mp$  and  $\tilde{\ell} \rightarrow \ell\tilde{G}$ . No evidence for an excess over the expected background is observed. Limits are derived (see Fig. 4) for a GMSB-type scenario with mass-degenerate right-handed sleptons (slepton co-NLSP scenario).
- 66 CHATRCHYAN 11W looked in  $1.14 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 2$  jets, large total jet energy, and  $\cancel{E}_T$ . After combining multi-jet events into two pseudo-jets signal events are selected by a cut on  $\alpha_T = E_T^{j2}/M_T$ , the transverse energy

- of the less energetic jet over the transverse mass. Given the lack of an excess over the SM backgrounds, limits are derived in the CMSSM  $(m_0, m_{1/2})$  plane (see Fig. 4) for  $\tan\beta = 10$ . The limits are only weakly dependent on  $\tan\beta$  and  $A_0$ .
- 67 KHACHATRYAN 11I looked in  $35 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 2$  jets and  $\cancel{E}_T$ . After combining multi-jet events into two pseudo-jets signal events are selected by a cut on  $\alpha_T = E_T^{j2}/M_T$ , the transverse energy of the less energetic jet over the transverse mass. No evidence for an excess over the expected background is observed. Limits are derived in the CMSSM  $(m_0, m_{1/2})$  plane (see Fig. 5) for  $\tan\beta = 3$ . Superseded by CHATRCHYAN 11W.
- 68 ABAZOV 02F looked in  $77.5 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at 1.8 TeV for events with  $\geq 2\mu + \geq 4$  jets, originating from associated production of squarks followed by an indirect  $\cancel{R}$  decay (of the  $\tilde{\chi}_1^0$ ) via  $LQ\bar{D}$  couplings of the type  $\lambda'_{2jk}$  where  $j=1,2$  and  $k=1,2,3$ . Bounds are obtained in the MSUGRA scenario by a scan in the range  $0 \leq M_0 \leq 400 \text{ GeV}$ ,  $60 \leq m_{1/2} \leq 120 \text{ GeV}$  for fixed values  $A_0=0$ ,  $\mu < 0$ , and  $\tan\beta=2$  or 6. The bounds are weaker for  $\tan\beta=6$ . See Figs. 2,3 for the exclusion contours in  $m_{1/2}$  versus  $m_0$  for  $\tan\beta=2$  and 6, respectively.
- 69 ABAZOV 02G search for associated production of gluinos and squarks in  $92.7 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8 \text{ TeV}$ , using events with one electron,  $\geq 4$  jets, and large  $\cancel{E}_T$ . The results are compared to a MSUGRA scenario with  $\mu < 0$ ,  $A_0=0$ , and  $\tan\beta=3$  and allow to exclude a region of the  $(m_0, m_{1/2})$  shown in Fig. 11.
- 70 CHEUNG 02B studies the constraints on a  $\tilde{b}_1$  with mass in the 2.2–5.5 GeV region and a gluino in the mass range 12–16 GeV, using precision measurements of  $Z^0$  decays and  $e^+e^-$  annihilations at LEP2. Few detectable events are predicted in the LEP2 data for the model proposed by BERGER 01.
- 71 BERGER 01 reanalyzed interpretation of Tevatron data on bottom-quark production. Argues that pair production of light gluinos ( $m \sim 12\text{--}16 \text{ GeV}$ ) with subsequent 2-body decay into a light sbottom ( $m \sim 2\text{--}5.5 \text{ GeV}$ ) and bottom can reconcile Tevatron data with predictions of perturbative QCD for the bottom production rate. The sbottom must either decay hadronically via a  $R$ -parity- and  $B$ -violating interaction, or be long-lived.
- 72 ABBOTT 99 searched for  $\gamma\cancel{E}_T + \geq 2$  jet final states, and set limits on  $\sigma(p\bar{p} \rightarrow \tilde{g} + X) \cdot B(\tilde{g} \rightarrow \gamma\cancel{E}_T X)$ . The quoted limits correspond to  $m_{\tilde{g}} \geq m_{\tilde{q}}$ , with  $B(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \gamma) = 1$  and  $B(\tilde{\chi}_1^0 \rightarrow \tilde{G} \gamma) = 1$ , respectively. They improve to 310 GeV (360 GeV in the case of  $\gamma\tilde{G}$  decay) for  $m_{\tilde{g}} = m_{\tilde{q}}$ .
- 73 ABBOTT 99K uses events with an electron pair and four jets to search for the decay of the  $\tilde{\chi}_1^0$  LSP via  $\cancel{R} LQ\bar{D}$  couplings. The particle spectrum and decay branching ratios are taken in the framework of minimal supergravity. An excluded region at 95% CL is obtained in the  $(m_0, m_{1/2})$  plane under the assumption that  $A_0=0$ ,  $\mu < 0$ ,  $\tan\beta=2$  and any one of the couplings  $\lambda'_{1jk} > 10^{-3}$  ( $j=1,2$  and  $k=1,2,3$ ) and from which the above limit is computed. For equal mass squarks and gluinos, the corresponding limit is 277 GeV. The results are essentially independent of  $A_0$ , but the limit deteriorates rapidly with increasing  $\tan\beta$  or  $\mu > 0$ .
- 74 ABACHI 95C assume five degenerate squark flavors with  $m_{\tilde{q}_L} = m_{\tilde{q}_R}$ . Sleptons are assumed to be heavier than squarks. The limits are derived for fixed  $\tan\beta = 2.0$ ,  $\mu = -250 \text{ GeV}$ , and  $m_{H^+} = 500 \text{ GeV}$ , and with the cascade decays of the squarks and gluinos calculated within the framework of the Minimal Supergravity scenario. The bounds are weakly sensitive to the three fixed parameters for a large fraction of parameter space.
- 75 ABE 95T looked for a cascade decay of gluino into  $\tilde{\chi}_2^0$  which further decays into  $\tilde{\chi}_1^0$  and a photon. No signal is observed. Limits vary widely depending on the choice of parameters. For  $\mu = -40 \text{ GeV}$ ,  $\tan\beta = 1.5$ , and heavy squarks, the range  $50 < m_{\tilde{g}} (\text{GeV}) < 140$  is excluded at 90% CL. See the paper for details.

- 76 HEBBEKER 93 combined jet analyses at various  $e^+e^-$  colliders. The 4-jet analyses at TRISTAN/LEP and the measured  $\alpha_s$  at PEP/PETRA/TRISTAN/LEP are used. A constraint on effective number of quarks  $N=6.3 \pm 1.1$  is obtained, which is compared to that with a light gluino,  $N=8$ .
- 77 ABE 92L bounds are based on similar assumptions as ABACHI 95C. Not sensitive to  $m_{\text{gluino}} < 40$  GeV (but other experiments rule out that region).
- 78 ROY 92 reanalyzed CDF limits on di-lepton events to obtain limits on gluino production in  $R$ -parity violating models. The 100% decay  $\tilde{g} \rightarrow q\bar{q}\tilde{\chi}$  where  $\tilde{\chi}$  is the LSP, and the LSP decays either into  $\ell q\bar{d}$  or  $\ell\ell\bar{e}$  is assumed.
- 79 NOJIRI 91 argues that a heavy gluino should be nearly degenerate with squarks in minimal supergravity not to overclose the universe.
- 80 The limits of ALBAJAR 87D are from  $p\bar{p} \rightarrow \tilde{g}\tilde{g}X$  ( $\tilde{g} \rightarrow q\bar{q}\tilde{\gamma}$ ) and assume  $m_{\tilde{q}} > m_{\tilde{g}}$ . These limits apply for  $m_{\tilde{\gamma}} \lesssim 20$  GeV and  $\tau(\tilde{g}) < 10^{-10}$  s.
- 81 The limit of ANSARI 87D assumes  $m_{\tilde{q}} > m_{\tilde{g}}$  and  $m_{\tilde{\gamma}} \approx 0$ .

### Long-lived/light $\tilde{g}$ (Gluino) MASS LIMIT

Limits on light gluinos ( $m_{\tilde{g}} < 5$  GeV), or gluinos which leave the detector before decaying.

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
> 985	95	1 AAD	13AA ATLS	$\tilde{g}$ , $R$ -hadrons, generic interaction model
> 832	95	2 AAD	13BC ATLS	$R$ -hadrons, $\tilde{g} \rightarrow g/q\bar{q}\tilde{\chi}_1^0$ , generic $R$ -hadron model, lifetime between $10^{-5}$ and $10^3$ s, $m_{\tilde{\chi}_1^0} = 100$ GeV
>1322	95	3 CHATRCHYAN 13AB	CMS	long-lived $\tilde{g}$ forming $R$ -hadrons, $f = 0.1$ , cloud interaction model
none 200–341	95	4 AAD	12P ATLS	long-lived $\tilde{g} \rightarrow g\tilde{\chi}_1^0$ , $m_{\tilde{\chi}_1^0} = 100$ GeV
> 640	95	5 CHATRCHYAN 12AN	CMS	long-lived $\tilde{g} \rightarrow g\tilde{\chi}_1^0$
>1098	95	6 CHATRCHYAN 12L	CMS	long-lived $\tilde{g}$ forming $R$ -hadrons, $f = 0.1$
> 586	95	7 AAD	11K ATLS	stable $\tilde{g}$
> 544	95	8 AAD	11P ATLS	stable $\tilde{g}$ , GMSB scenario, $\tan\beta=5$
> 370	95	9 KHACHATRY...11	CMS	long lived $\tilde{g}$
> 398	95	10 KHACHATRY...11C	CMS	stable $\tilde{g}$
> 15	90	11 BERGER	10 THEO	hadron scattering data, $\alpha_s$
> 51	95	12 KAPLAN	08 THEO	event shapes at LEP
		13 ABAZOV	07L D0	long-lived $\tilde{g}$
> 12		14 BERGER	05 THEO	hadron scattering data
none 2–18	95	15 ABDALLAH	03C DLPH	$e^+e^- \rightarrow q\bar{q}\tilde{g}\tilde{g}$ , stable $\tilde{g}$
> 5		16 ABDALLAH	03G DLPH	QCD beta function
		17 HEISTER	03 ALEP	Color factors
> 26.9	95	18 HEISTER	03H ALEP	$e^+e^- \rightarrow q\bar{q}\tilde{g}\tilde{g}$
> 6.3		19 JANOT	03 RVUE	$\Delta\Gamma_{had} < 3.9$ MeV
		20 MAFI	00 THEO	$pp \rightarrow \text{jets} + \cancel{p}_T$
		21 ALAVI-HARATI 199E	KTEV	$pN \rightarrow R^0$ , with $R^0 \rightarrow \rho^0\tilde{\gamma}$ and $R^0 \rightarrow \pi^0\tilde{\gamma}$

		22	BAER	99	RVUE	Stable $\tilde{g}$ hadrons	
		23	FANTI	99	NA48	$p\text{Be} \rightarrow R^0 \rightarrow \eta\tilde{\gamma}$	
		24	ACKERSTAFF	98V	OPAL	$e^+e^- \rightarrow \tilde{\chi}_1^+\tilde{\chi}_1^-$	
		25	ADAMS	97B	KTEV	$pN \rightarrow R^0 \rightarrow \rho^0\tilde{\gamma}$	
		26	ALBUQUERQ..	97	E761	$R^+(uud\tilde{g}) \rightarrow S^0(uds\tilde{g})\pi^+$ , $X^-(ssd\tilde{g}) \rightarrow S^0\pi^-$	
>	6.3	95	27	BARATE	97L	ALEP	Color factors
>	5	99	28	CSIKOR	97	RVUE	$\beta$ function, $Z \rightarrow \text{jets}$
>	1.5	90	29	DEGOUVEA	97	THEO	$Z \rightarrow jjjj$
			30	FARRAR	96	RVUE	$R^0 \rightarrow \pi^0\tilde{\gamma}$
none	1.9–13.6	95	31	AKERS	95R	OPAL	$Z$ decay into a long-lived $(\tilde{g}q\bar{q})^\pm$
<	0.7		32	CLAVELLI	95	RVUE	quarkonia
none	1.5–3.5		33	CAKIR	94	RVUE	$\Upsilon(1S) \rightarrow \gamma + \text{gluonium}$
not	3–5		34	LOPEZ	93C	RVUE	LEP
$\approx$	4		35	CLAVELLI	92	RVUE	$\alpha_s$ running
			36	ANTONIADIS	91	RVUE	$\alpha_s$ running
>	1		37	ANTONIADIS	91	RVUE	$pN \rightarrow \text{missing energy}$
			38	NAKAMURA	89	SPEC	$R\text{-}\Delta^{++}$
>	3.8	90	39	ARNOLD	87	EMUL	$\pi^-$ (350 GeV). $\sigma \simeq A^1$
>	3.2	90	39	ARNOLD	87	EMUL	$\pi^-$ (350 GeV). $\sigma \simeq A^{0.72}$
none	0.6–2.2	90	40	TUTS	87	CUSB	$\Upsilon(1S) \rightarrow \gamma + \text{gluonium}$
none	1–4.5	90	41	ALBRECHT	86C	ARG	$1 \times 10^{-11} \lesssim \tau \lesssim 1 \times 10^{-9} \text{s}$
none	1–4	90	42	BADIER	86	BDMP	$1 \times 10^{-10} < \tau < 1 \times 10^{-7} \text{s}$
none	3–5		43	BARNETT	86	RVUE	$p\bar{p} \rightarrow \text{gluino gluino gluon}$
none			44	VOLOSHIN	86	RVUE	If (quasi) stable; $\tilde{g}uud$
none	0.5–2		45	COOPER-...	85B	BDMP	For $m_{\tilde{q}}=300$ GeV
none	0.5–4		45	COOPER-...	85B	BDMP	For $m_{\tilde{q}} < 65$ GeV
none	0.5–3		45	COOPER-...	85B	BDMP	For $m_{\tilde{q}}=150$ GeV
none	2–4		46	DAWSON	85	RVUE	$\tau > 10^{-7} \text{s}$
none	1–2.5		46	DAWSON	85	RVUE	For $m_{\tilde{q}}=100$ GeV
none	0.5–4.1	90	47	FARRAR	85	RVUE	FNAL beam dump
>	1		48	GOLDMAN	85	RVUE	Gluonium
>	1–2		49	HABER	85	RVUE	
			50	BALL	84	CALO	
			51	BRICK	84	RVUE	
			52	FARRAR	84	RVUE	
>	2		53	BERGSMA	83C	RVUE	For $m_{\tilde{q}} < 100$ GeV
>	2–3		54	CHANOWITZ	83	RVUE	$\tilde{g}u\bar{d}, \tilde{g}uud$
>	1.5–2		55	KANE	82	RVUE	Beam dump
				FARRAR	78	RVUE	$R$ -hadron

<sup>1</sup>AAD 13AA searched in  $4.7 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events containing colored long-lived particles that hadronize forming  $R$ -hadrons. No significant excess above the expected background was found. Long-lived  $R$ -hadrons containing a  $\tilde{g}$  are excluded for masses up to 985 GeV at 95% C.L in a general interaction model. Also, limits independent of the fraction of  $R$ -hadrons that arrive charged in the muon system were derived, see Fig. 6.

- <sup>2</sup> AAD 13BC searched in  $5.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  and in  $22.9 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  for bottom squark R-hadrons that have come to rest within the ATLAS calorimeter and decay at some later time to hadronic jets and a neutralino. In absence of an excess of events above the expected backgrounds, limits are set on gluino masses for different decays, lifetimes, and neutralino masses, see their Table 6 and Fig. 10.
- <sup>3</sup> CHATRCHYAN 13AB looked in  $5.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  and in  $18.8 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  for events with heavy stable particles, identified by their anomalous  $dE/dx$  in the tracker or additionally requiring that it be identified as muon in the muon chambers, from pair production of  $\tilde{g}$ 's. No evidence for an excess over the expected background is observed. Limits are derived for pair production of gluinos as a function of mass (see Fig. 8 and Table 5), depending on the fraction,  $f$ , of formation of  $\tilde{g}-g$  (R-gluonball) states. The quoted limit is for  $f = 0.1$ , while for  $f = 0.5$  it degrades to 1276 GeV. In the conservative scenario where every hadronic interaction causes it to become neutral, the limit decreases to 928 GeV for  $f = 0.1$ .
- <sup>4</sup> AAD 12P looked in  $31 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with pair production of long-lived gluinos. The hadronization of the gluinos leads to  $R$ -hadrons which may stop inside the detector and later decay via  $\tilde{g} \rightarrow g\tilde{\chi}_1^0$  during gaps between the proton bunches. No significant excess over the expected background is observed. From a counting experiment, a limit at 95% C.L. on the cross section as a function of  $m_{\tilde{g}}$  is derived for  $m_{\tilde{\chi}_1^0} = 100 \text{ GeV}$ , see Fig. 4. The limit is valid for lifetimes between  $10^{-5}$  and  $10^3$  seconds and assumes the *Generic* matter interaction model for the production cross section.
- <sup>5</sup> CHATRCHYAN 12AN looked in  $4.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with pair production of long-lived gluinos. The hadronization of the gluinos leads to  $R$ -hadrons which may stop inside the detector and later decay via  $\tilde{g} \rightarrow g\tilde{\chi}_1^0$  during gaps between the proton bunches. No significant excess over the expected background is observed. From a counting experiment, a limit at 95% C.L. on the cross section as a function of  $m_{\tilde{g}}$  is derived, see Fig. 3. The mass limit is valid for lifetimes between  $10^{-5}$  and  $10^3$  seconds, for what they call "the daughter gluon energy  $E_g > 100 \text{ GeV}$  and assuming the *cloud* interaction model for  $R$ -hadrons. Supersedes KHACHATRYAN 11.
- <sup>6</sup> CHATRCHYAN 12L looked in  $5.0 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with heavy stable particles, identified by their anomalous  $dE/dx$  in the tracker or additionally requiring that it be identified as muon in the muon chambers, from pair production of  $\tilde{g}$ 's. No evidence for an excess over the expected background is observed. Limits are derived for pair production of gluinos as a function of mass (see Fig. 3), depending on the fraction,  $f$ , of formation of  $\tilde{g}-g$  ( $R$ -glueball) states. The quoted limit is for  $f = 0.1$ , while for  $f = 0.5$  it degrades to 1046 GeV. In the conservative scenario where every hadronic interaction causes it to become neutral, the limit decreases to 928 GeV for  $f=0.1$ . Supersedes KHACHATRYAN 11C.
- <sup>7</sup> AAD 11K looked in  $34 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with heavy stable particles, identified by their anomalous  $dE/dx$  in the tracker or time of flight in the tile calorimeter, from pair production of  $\tilde{g}$ . No evidence for an excess over the SM expectation is observed. Limits are derived for pair production of gluinos as a function of mass (see Fig. 4), for a fraction,  $f = 10\%$ , of formation of  $\tilde{g} - g$  ( $R$ -gluonball). If instead of a phase space driven approach for the hadronic scattering of the  $R$ -hadrons, a triple-Regge model or a bag-model is used, the limit degrades to 566 and 562 GeV, respectively.
- <sup>8</sup> AAD 11P looked in  $37 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with heavy stable particles, reconstructed and identified by their time of flight in the Muon System. There is no requirement on their observation in the tracker to increase the sensitivity to cases where gluinos have a large fraction,  $f$ , of formation of neutral  $\tilde{g} - g$  ( $R$ -gluonball). No evidence for an excess over the SM expectation is observed. Limits are derived as a function of mass (see Fig. 4), for  $f=0.1$ . For fractions  $f = 0.5$  and  $1.0$  the limit degrades to 537 and 530 GeV, respectively.

- <sup>9</sup> KHACHATRYAN 11 looked in  $10 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with pair production of long-lived gluinos. The hadronization of the gluinos leads to R-hadrons which may stop inside the detector and later decay via  $\tilde{g} \rightarrow g\tilde{\chi}_1^0$  during gaps between the proton bunches. No significant excess over the expected background is observed. From a counting experiment, a limit at 95% C.L. on the cross section times branching ratio is derived for  $m_{\tilde{g}} - m_{\tilde{\chi}_1^0} > 100 \text{ GeV}$ , see their Fig. 2. Assuming 100% branching ratio, lifetimes between 75 ns and  $3 \times 10^5 \text{ s}$  are excluded for  $m_{\tilde{g}} = 300 \text{ GeV}$ . The  $\tilde{g}$  mass exclusion is obtained with the same assumptions for lifetimes between 10  $\mu\text{s}$  and 1000 s, but shows some dependence on the model for R-hadron interactions with matter, illustrated in Fig. 3. From a time-profile analysis, the mass exclusion is 382 GeV for a lifetime of 10  $\mu\text{s}$  under the same assumptions as above.
- <sup>10</sup> KHACHATRYAN 11C looked in  $3.1 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with heavy stable particles, identified by their anomalous  $dE/dx$  in the tracker or additionally requiring that it be identified as muon in the muon chambers, from pair production of  $\tilde{g}$ . No evidence for an excess over the expected background is observed. Limits are derived for pair production of gluinos as a function of mass (see Fig. 3), depending on the fraction,  $f$ , of formation of  $\tilde{g} - g$  (R-gluonball). The quoted limit is for  $f=0.1$ , while for  $f=0.5$  it degrades to 357 GeV. In the conservative scenario where every hadronic interaction causes it to become neutral, the limit decreases to 311 GeV for  $f=0.1$ .
- <sup>11</sup> BERGER 10 updated the results of BERGER 05. They fit parton distribution functions including the effects of a light gluino as an extra parton. Different data on  $\alpha_s$  is also included. A fit for  $\alpha_s(M_Z)$  is performed as a function of the gluino mass. The bound is determined by comparing the quality of the fit to the CT10 fit, and the CT10 tolerance criterion is used to define the significance. The lower bound is 25 GeV for fixed  $\alpha_s(M_Z) = 0.118$ .
- <sup>12</sup> KAPLAN 08 reanalysed jet event shape data from LEP 1 and LEP 2 using soft collinear effective theory methods. These data are sensitive to the effects of new degrees of freedoms, including a relatively light gluino, at different energy scales, roughly between 5 and 50 GeV. The analysis relies on theoretical modeling of and approximations for non-perturbative effects and matching between different scales.
- <sup>13</sup> ABAZOV 07L looked in approximately  $410 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events with a long-lived gluino from split supersymmetry, decaying after stopping in the detector into  $g\tilde{\chi}_1^0$  with lifetimes from 30  $\mu\text{s}$  to 100 h. The signal signature is a largely empty event with a single large transverse energy deposit in the calorimeter. The main background is due to cosmic muons interacting in the calorimeter. The data agree with the estimated background and allow the authors to estimate a limit on the rate of an out-of-time monojet signal of a given energy. Assuming the branching ratios  $\tilde{g} \rightarrow g\tilde{\chi}_1^0$  to be 100% the results can be translated to limits on the gluino cross section versus the gluino mass for fixed  $\tilde{\chi}_1^0$  mass. After comparing to the expected gluino cross sections, the excluded region of gluino masses can be obtained, see examples in their Fig. 3.
- <sup>14</sup> BERGER 05 include the light gluino in proton PDF and perform global analysis of hadronic data. Effects on the running of  $\alpha_s$  also included. Strong dependency on  $\alpha_s(m_Z)$ . Bound quoted for  $\alpha_s(m_Z) = 0.118$ . Superseded by BERGER 10.
- <sup>15</sup> ABDALLAH 03C looked for events of the type  $q\bar{q}R^\pm R^\pm$ ,  $q\bar{q}R^\pm R^0$  or  $q\bar{q}R^0 R^0$  in  $e^+e^-$  interactions at 91.2 GeV collected in 1994. The  $R^\pm$  bound states are identified by anomalous  $dE/dx$  in the tracking chambers and the  $R^0$  by missing energy, due to their reduced energy loss in the calorimeters. The upper value of the excluded range depends on the probability for the gluino to fragment into  $R^\pm$  or  $R^0$ , see their Fig. 17. It improves to 23 GeV for 100% fragmentation to  $R^\pm$ .
- <sup>16</sup> ABDALLAH 03G used  $e^+e^-$  data at and around the  $Z^0$  peak, above the  $Z^0$  up to  $\sqrt{s} = 202 \text{ GeV}$  and events from radiative return to cover the low energy region. They perform a direct measurement of the QCD beta-function from the means of fully inclusive event observables. Compared to the energy range, gluinos below 5 GeV can be considered massless and are firmly excluded by the measurement.

- 17 HEISTER 03 use  $e^+e^-$  data from 1994 and 1995 at and around the  $Z^0$  peak to measure the 4-jet rate and angular correlations. The comparison with QCD NLO calculations allow  $\alpha_S(M_Z)$  and the color factor ratios to be extracted and the results are in agreement with the expectations from QCD. The inclusion of a massless gluino in the beta functions yields  $T_R / C_F = 0.15 \pm 0.06 \pm 0.06$  (expectation is  $T_R / C_F = 3/8$ ), excluding a massless gluino at more than 95% CL. As no NLO calculations are available for massive gluinos, the earlier LO results from BARATE 97L for massive gluinos remain valid.
- 18 HEISTER 03H use  $e^+e^-$  data at and around the  $Z^0$  peak to look for stable gluinos hadronizing into charged or neutral  $R$ -hadrons with arbitrary branching ratios. Combining these results with bounds on the  $Z^0$  hadronic width from electroweak measurements (JANOT 03) to cover the low mass region the quoted lower limit on the mass of a long-lived gluino is obtained.
- 19 JANOT 03 excludes a light gluino from the upper limit on an additional contribution to the  $Z$  hadronic width. At higher confidence levels,  $m_{\tilde{g}} > 5.3(4.2)$  GeV at  $3\sigma(5\sigma)$  level.
- 20 MAFI 00 reanalyzed CDF data assuming a stable heavy gluino as the LSP, with model for  $R$ -hadron-nucleon scattering. Gluino masses between 35 GeV and 115 GeV are excluded based on the CDF Run I data. Combined with the analysis of BAER 99, this allows a LSP gluino mass between 25 and 35 GeV if the probability of fragmentation into charged  $R$ -hadron  $P > 1/2$ . The cosmological exclusion of such a gluino LSP are assumed to be avoided as in BAER 99. Gluino could be NLSP with  $\tau_{\tilde{g}} \sim 100$  yrs, and decay to gluon gravitino.
- 21 ALAVI-HARATI 99E looked for  $R^0$  bound states, yielding  $\pi^+\pi^-$  or  $\pi^0$  in the final state. The experiment is sensitive to values of  $\Delta m = m_{R^0} - m_{\tilde{\gamma}}$  larger than 280 MeV and 140 MeV for the two decay modes, respectively, and to  $R^0$  mass and lifetime in the ranges 0.8–5 GeV and  $10^{-10}$ – $10^{-3}$  s. The limits obtained depend on  $B(R^0 \rightarrow \pi^+\pi^- \text{ photino})$  and  $B(R^0 \rightarrow \pi^0 \text{ photino})$  on the value of  $m_{R^0}/m_{\tilde{\gamma}}$ , and on the ratio of production rates  $\sigma(R^0)/\sigma(K_L^0)$ . See Figures in the paper for the excluded  $R^0$  production rates as a function of  $\Delta m$ ,  $R^0$  mass and lifetime. Using the production rates expected from perturbative QCD, and assuming dominance of the above decay channels over the suitable phase space,  $R^0$  masses in the range 0.8–5 GeV are excluded at 90%CL for a large fraction of the sensitive lifetime region. ALAVI-HARATI 99E updates and supersedes the results of ADAMS 97B.
- 22 BAER 99 set constraints on the existence of stable  $\tilde{g}$  hadrons, in the mass range  $m_{\tilde{g}} > 3$  GeV. They argue that strong-interaction effects in the low-energy annihilation rates could leave small enough relic densities to evade cosmological constraints up to  $m_{\tilde{g}} < 10$  TeV. They consider  $\text{jet} + \cancel{E}_T$  as well as heavy-ionizing charged-particle signatures from production of stable  $\tilde{g}$  hadrons at LEP and Tevatron, developing modes for the energy loss of  $\tilde{g}$  hadrons inside the detectors. Results are obtained as a function of the fragmentation probability  $P$  of the  $\tilde{g}$  into a charged hadron. For  $P < 1/2$ , and for various energy-loss models, OPAL and CDF data exclude gluinos in the  $3 < m_{\tilde{g}}(\text{GeV}) < 130$  mass range. For  $P > 1/2$ , gluinos are excluded in the mass ranges  $3 < m_{\tilde{g}}(\text{GeV}) < 23$  and  $50 < m_{\tilde{g}}(\text{GeV}) < 200$ .
- 23 FANTI 99 looked for  $R^0$  bound states yielding high  $P_T \eta \rightarrow 3\pi^0$  decays. The experiment is sensitive to a region of  $R^0$  mass and lifetime in the ranges of 1–5 GeV and  $10^{-10}$ – $10^{-3}$  s. The limits obtained depend on  $B(R^0 \rightarrow \eta\tilde{\gamma})$ , on the value of  $m_{R^0}/m_{\tilde{\gamma}}$ , and on the ratio of production rates  $\sigma(R^0)/\sigma(K_L^0)$ . See Fig. 6–7 for the excluded production rates as a function of  $R^0$  mass and lifetime.
- 24 ACKERSTAFF 98V excludes the light gluino with universal gaugino mass where charginos, neutralinos decay as  $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0 \rightarrow q\bar{q}\tilde{g}$  from total hadronic cross sections at  $\sqrt{s}=130$ –172 GeV. See paper for the case of nonuniversal gaugino mass.
- 25 ADAMS 97B looked for  $\rho^0 \rightarrow \pi^+\pi^-$  as a signature of  $R^0 = (\tilde{g}g)$  bound states. The experiment is sensitive to an  $R^0$  mass range of 1.2–4.5 GeV and to a lifetime range of

- $10^{-10}$ – $10^{-3}$  sec. Precise limits depend on the assumed value of  $m_{R^0}/m_{\tilde{\gamma}}$ . See Fig. 7 for the excluded mass and lifetime region.
- 26 ALBUQUERQUE 97 looked for weakly decaying baryon-like states which contain a light gluino, following the suggestions in FARRAR 96. See their Table 1 for limits on the production fraction. These limits exclude gluino masses in the range 100–600 MeV for the predicted lifetimes (FARRAR 96) and production rates, which are assumed to be comparable to those of strange or charmed baryons.
  - 27 BARATE 97L studied the QCD color factors from four-jet angular correlations and the differential two-jet rate in  $Z$  decay. Limit obtained from the determination of  $n_f = 4.24 \pm 0.29 \pm 1.15$ , assuming  $T_F/C_F=3/8$  and  $C_A/C_F=9/4$ .
  - 28 CSIKOR 97 combined the  $\alpha_s$  from  $\sigma(e^+e^- \rightarrow \text{hadron})$ ,  $\tau$  decay, and jet analysis in  $Z$  decay. They exclude a light gluino below 5 GeV at more than 99.7%CL.
  - 29 DEGOUVEA 97 reanalyzed AKERS 95A data on  $Z$  decay into four jets to place constraints on a light stable gluino. The mass limit corresponds to the pole mass of 2.8 GeV. The analysis, however, is limited to the leading-order QCD calculation.
  - 30 FARRAR 96 studied the possible  $R^0=(\tilde{g}g)$  component in Fermilab E799 experiment and used its bound  $B(K_L^0 \rightarrow \pi^0 \nu \bar{\nu}) \leq 5.8 \times 10^{-5}$  to place constraints on the combination of  $R^0$  production cross section and its lifetime.
  - 31 AKERS 95R looked for  $Z$  decay into  $q\bar{q}\tilde{g}\tilde{g}$ , by searching for charged particles with  $dE/dx$  consistent with  $\tilde{g}$  fragmentation into a state  $(\tilde{g}q\bar{q})^\pm$  with lifetime  $\tau > 10^{-7}$  sec. The fragmentation probability into a charged state is assumed to be 25%.
  - 32 CLAVELLI 95 updates the analysis of CLAVELLI 93, based on a comparison of the hadronic widths of charmonium and bottomonium  $S$ -wave states. The analysis includes a parametrization of relativistic corrections. Claims that the presence of a light gluino improves agreement with the data by slowing down the running of  $\alpha_s$ .
  - 33 CAKIR 94 reanalyzed TUTS 87 and later unpublished data from CUSB to exclude pseudo-scalar gluinonium  $\eta_{\tilde{g}}(\tilde{g}\tilde{g})$  of mass below 7 GeV. It was argued, however, that the perturbative QCD calculation of the branching fraction  $\mathcal{T} \rightarrow \eta_{\tilde{g}}\gamma$  is unreliable for  $m_{\eta_{\tilde{g}}} < 3$  GeV. The gluino mass is defined by  $m_{\tilde{g}}=(m_{\eta_{\tilde{g}}})/2$ . The limit holds for any gluino lifetime.
  - 34 LOPEZ 93C uses combined restraint from the radiative symmetry breaking scenario within the minimal supergravity model, and the LEP bounds on the  $(M_2, \mu)$  plane. Claims that the light gluino window is strongly disfavored.
  - 35 CLAVELLI 92 claims that a light gluino mass around 4 GeV should exist to explain the discrepancy between  $\alpha_s$  at LEP and at quarkonia ( $\mathcal{T}$ ), since a light gluino slows the running of the QCD coupling.
  - 36 ANTONIADIS 91 argue that possible light gluinos ( $< 5$  GeV) contradict the observed running of  $\alpha_s$  between 5 GeV and  $m_Z$ . The significance is less than 2 s.d.
  - 37 ANTONIADIS 91 interpret the search for missing energy events in 450 GeV/c  $pN$  collisions, AKESSON 91, in terms of light gluinos.
  - 38 NAKAMURA 89 searched for a long-lived ( $\tau \gtrsim 10^{-7}$  s) charge-( $\pm 2$ ) particle with mass  $\lesssim 1.6$  GeV in proton-Pt interactions at 12 GeV and found that the yield is less than  $10^{-8}$  times that of the pion. This excludes  $R\text{-}\Delta^{++}$  (a  $\tilde{g}uuu$  state) lighter than 1.6 GeV.
  - 39 The limits assume  $m_{\tilde{q}} = 100$  GeV. See their figure 3 for limits vs.  $m_{\tilde{q}}$ .
  - 40 The gluino mass is defined by half the bound  $\tilde{g}\tilde{g}$  mass. If zero gluino mass gives a  $\tilde{g}\tilde{g}$  of mass about 1 GeV as suggested by various glueball mass estimates, then the low-mass bound can be replaced by zero. The high-mass bound is obtained by comparing the data with nonrelativistic potential-model estimates.
  - 41 ALBRECHT 86C search for secondary decay vertices from  $\chi_{b1}(1P) \rightarrow \tilde{g}\tilde{g}g$  where  $\tilde{g}$ 's make long-lived hadrons. See their figure 4 for excluded region in the  $m_{\tilde{g}} - m_{\tilde{q}}$  and  $m_{\tilde{g}} - m_{\tilde{q}}$  plane. The lower  $m_{\tilde{g}}$  region below  $\sim 2$  GeV may be sensitive to fragmentation

- effects. Remark that the  $\tilde{g}$ -hadron mass is expected to be  $\sim 1$  GeV (glueball mass) in the zero  $\tilde{g}$  mass limit.
- 42 BADIÉR 86 looked for secondary decay vertices from long-lived  $\tilde{g}$ -hadrons produced at 300 GeV  $\pi^-$  beam dump. The quoted bound assumes  $\tilde{g}$ -hadron nucleon total cross section of  $10\mu\text{b}$ . See their figure 7 for excluded region in the  $m_{\tilde{g}} - m_{\tilde{q}}$  plane for several assumed total cross-section values.
- 43 BARNETT 86 rule out light gluinos ( $m = 3\text{--}5$  GeV) by calculating the monojet rate from gluino gluino gluon events (and from gluino gluon events) and by using UA1 data from  $p\bar{p}$  collisions at CERN.
- 44 VOLOSHIN 86 rules out stable gluino based on the cosmological argument that predicts too much hydrogen consisting of the charged stable hadron  $\tilde{g}uud$ . Quasi-stable ( $\tau > 1. \times 10^{-7}$  s) light gluino of  $m_{\tilde{g}} < 3$  GeV is also ruled out by nonobservation of the stable charged particles,  $\tilde{g}uud$ , in high energy hadron collisions.
- 45 COOPER-SARKAR 85B is BEBC beam-dump. Gluinos decaying in dump would yield  $\tilde{\gamma}$ 's in the detector giving neutral-current-like interactions. For  $m_{\tilde{q}} > 330$  GeV, no limit is set.
- 46 DAWSON 85 first limit from neutral particle search. Second limit based on FNAL beam dump experiment.
- 47 FARRAR 85 points out that BALL 84 analysis applies only if the  $\tilde{g}$ 's decay before interacting, i.e.  $m_{\tilde{q}} < 80m_{\tilde{g}}^{1.5}$ . FARRAR 85 finds  $m_{\tilde{g}} < 0.5$  not excluded for  $m_{\tilde{q}} = 30\text{--}1000$  GeV and  $m_{\tilde{g}} < 1.0$  not excluded for  $m_{\tilde{q}} = 100\text{--}500$  GeV by BALL 84 experiment.
- 48 GOLDMAN 85 use nonobservation of a pseudoscalar  $\tilde{g}\text{--}\tilde{g}$  bound state in radiative  $\psi$  decay.
- 49 HABER 85 is based on survey of all previous searches sensitive to low mass  $\tilde{g}$ 's. Limit makes assumptions regarding the lifetime and electric charge of the lightest supersymmetric particle.
- 50 BALL 84 is FNAL beam dump experiment. Observed no interactions of  $\tilde{\gamma}$  in the calorimeter, where  $\tilde{\gamma}$ 's are expected to come from pair-produced  $\tilde{g}$ 's. Search for long-lived  $\tilde{\gamma}$  interacting in calorimeter 56m from target. Limit is for  $m_{\tilde{q}} = 40$  GeV and production cross section proportional to  $A^{0.72}$ . BALL 84 find no  $\tilde{g}$  allowed below 4.1 GeV at CL = 90%. Their figure 1 shows dependence on  $m_{\tilde{q}}$  and A. See also KANE 82.
- 51 BRICK 84 reanalyzed FNAL 147 GeV HBC data for  $R\text{--}\Delta(1232)^{++}$  with  $\tau > 10^{-9}$  s and  $p_{\text{lab}} > 2$  GeV. Set CL = 90% upper limits 6.1, 4.4, and 29 microbarns in  $pp$ ,  $\pi^+p$ ,  $K^+p$  collisions respectively.  $R\text{--}\Delta^{++}$  is defined as being  $\tilde{g}$  and 3 up quarks. If mass = 1.2–1.5 GeV, then limits may be lower than theory predictions.
- 52 FARRAR 84 argues that  $m_{\tilde{g}} < 100$  MeV is not ruled out if the lightest R-hadrons are long-lived. A long lifetime would occur if R-hadrons are lighter than  $\tilde{\gamma}$ 's or if  $m_{\tilde{q}} > 100$  GeV.
- 53 BERGSMAN 83C is reanalysis of CERN-SPS beam-dump data. See their figure 1.
- 54 CHANOWITZ 83 find in bag-model that charged  $s$ -hadron exists which is stable against strong decay if  $m_{\tilde{g}} < 1$  GeV. This is important since tracks from decay of neutral  $s$ -hadron cannot be reconstructed to primary vertex because of missed  $\tilde{\gamma}$ . Charged  $s$ -hadron leaves track from vertex.
- 55 KANE 82 inferred above  $\tilde{g}$  mass limit from retroactive analysis of hadronic collision and beam dump experiments. Limits valid if  $\tilde{g}$  decays inside detector.
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**LIGHT  $\tilde{G}$  (Gravitino) MASS LIMITS FROM COLLIDER EXPERIMENTS**

The following are bounds on light ( $\ll 1$  eV) gravitino indirectly inferred from its coupling to matter suppressed by the gravitino decay constant.

Unless otherwise stated, all limits assume that other supersymmetric particles besides the gravitino are too heavy to be produced. The gravitino is assumed to be undetected and to give rise to a missing energy ( $\cancel{E}$ ) signature.

VALUE (eV)	CL%	DOCUMENT ID	TECN	COMMENT
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●				
$> 1.09 \times 10^{-5}$	95	1 ABDALLAH	05B DLPH	$e^+ e^- \rightarrow \tilde{G} \tilde{G} \gamma$
$> 1.35 \times 10^{-5}$	95	2 ACHARD	04E L3	$e^+ e^- \rightarrow \tilde{G} \tilde{G} \gamma$
$> 1.3 \times 10^{-5}$		3 HEISTER	03C ALEP	$e^+ e^- \rightarrow \tilde{G} \tilde{G} \gamma$
$> 11.7 \times 10^{-6}$	95	4 ACOSTA	02H CDF	$p\bar{p} \rightarrow \tilde{G} \tilde{G} \gamma$
$> 8.7 \times 10^{-6}$	95	5 ABBIENDI,G	00D OPAL	$e^+ e^- \rightarrow \tilde{G} \tilde{G} \gamma$
$> 10.0 \times 10^{-6}$	95	6 ABREU	00Z DLPH	$e^+ e^- \rightarrow \tilde{G} \tilde{G} \gamma$
$> 11 \times 10^{-6}$	95	7 AFFOLDER	00J CDF	$p\bar{p} \rightarrow \tilde{G} \tilde{G} + \text{jet}$
$> 8.9 \times 10^{-6}$	95	8 ACCIARRI	99R L3	$e^+ e^- \rightarrow \tilde{G} \tilde{G} \gamma$
$> 7.9 \times 10^{-6}$	95	9 ACCIARRI	98V L3	$e^+ e^- \rightarrow \tilde{G} \tilde{G} \gamma$
$> 8.3 \times 10^{-6}$	95	9 BARATE	98J ALEP	$e^+ e^- \rightarrow \tilde{G} \tilde{G} \gamma$

<sup>1</sup> ABDALLAH 05B use data from  $\sqrt{s} = 180\text{--}208$  GeV. They look for events with a single photon +  $\cancel{E}$  final states from which a cross section limit of  $\sigma < 0.18$  pb at 208 GeV is obtained, allowing a limit on the mass to be set. Supersedes the results of ABREU 00Z.

<sup>2</sup> ACHARD 04E use data from  $\sqrt{s} = 189\text{--}209$  GeV. They look for events with a single photon +  $\cancel{E}$  final states from which a limit on the Gravitino mass is set corresponding to  $\sqrt{F} > 238$  GeV. Supersedes the results of ACCIARRI 99R.

<sup>3</sup> HEISTER 03C use the data from  $\sqrt{s} = 189\text{--}209$  GeV to search for  $\gamma \cancel{E}_T$  final states.

<sup>4</sup> ACOSTA 02H looked in  $87$  pb<sup>-1</sup> of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8$  TeV for events with a high- $E_T$  photon and  $\cancel{E}_T$ . They compared the data with a GMSB model where the final state could arise from  $q\bar{q} \rightarrow \tilde{G} \tilde{G} \gamma$ . Since the cross section for this process scales as  $1/|F|^4$ , a limit at 95% CL is derived on  $|F|^{1/2} > 221$  GeV. A model independent limit for the above topology is also given in the paper.

<sup>5</sup> ABBIENDI,G 00D searches for  $\gamma \cancel{E}$  final states from  $\sqrt{s}=189$  GeV.

<sup>6</sup> ABREU 00Z search for  $\gamma \cancel{E}$  final states using data from  $\sqrt{s}=189$  GeV. Superseded by ABDALLAH 05B.

<sup>7</sup> AFFOLDER 00J searches for final states with an energetic jet (from quark or gluon) and large  $\cancel{E}_T$  from undetected gravitinos.

<sup>8</sup> ACCIARRI 99R search for  $\gamma \cancel{E}$  final states using data from  $\sqrt{s}=189$  GeV. Superseded by ACHARD 04E.

<sup>9</sup> Searches for  $\gamma \cancel{E}$  final states at  $\sqrt{s}=183$  GeV.

## Supersymmetry Miscellaneous Results

Results that do not appear under other headings or that make nonminimal assumptions.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
• • •		We do not use the following data for averages, fits, limits, etc. • • •		
none 100–185	95	<sup>1</sup> AAD	13P ATLS	dark $\gamma$ , hidden valley
		<sup>2</sup> AALTONEN	12AB CDF	hidden-valley Higgs
		<sup>3</sup> AAD	11AA ATLS	scalar gluons
		<sup>4</sup> CHATRCHYAN	11E CMS	$\mu\mu$ resonances
		<sup>5</sup> ABAZOV	10N D0	$\gamma_D$ , hidden valley
		<sup>6</sup> LOVE	08A CLEO	$R, Y \rightarrow \mu\tau$
		<sup>7</sup> ABULENCIA	06P CDF	$l\gamma\cancel{E}_T, ll\gamma, \text{GMSB}$
		<sup>8</sup> ACOSTA	04E CDF	
		<sup>9</sup> TCHIKILEV	04 ISTR	$K^- \rightarrow \pi^- \pi^0 P$
		<sup>10</sup> AFFOLDER	02D CDF	$p\bar{p} \rightarrow \gamma b (\cancel{E}_T)$
		<sup>11</sup> AFFOLDER	01H CDF	$p\bar{p} \rightarrow \gamma\gamma X$
		<sup>12</sup> ABBOTT	00G D0	$p\bar{p} \rightarrow 3l + \cancel{E}_T, R, LL\bar{E}$
		<sup>13</sup> ABREU,P	00C DLPH	$e^+e^- \rightarrow \gamma + S/P$
		<sup>14</sup> ABACHI	97 D0	$\gamma\gamma X$
		<sup>15</sup> BARBER	84B RVUE	
		<sup>16</sup> HOFFMAN	83 CNTR	$\pi p \rightarrow n(e^+e^-)$

<sup>1</sup> AAD 13P searched in  $5 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for single lepton-jets with at least four muons; pairs of lepton-jets, each with two or more muons; and pairs of lepton-jets with two or more electrons. All of these could be signatures of Hidden Valley supersymmetric models. No statistically significant deviations from the Standard Model expectations are found. 95% C.L. limits are placed on the production cross section times branching ratio of dark photons for several parameter sets of a Hidden Valley model.

<sup>2</sup> AALTONEN 12AB looked in  $5.1 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for anomalous production of multiple low-energy leptons in association with a  $W$  or  $Z$  boson. Such events may occur in hidden valley models in which a supersymmetric Higgs boson is produced in association with a  $W$  or  $Z$  boson, with  $H \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0$  pair and with the  $\tilde{\chi}_1^0$  further decaying into a dark photon ( $\gamma_D$ ) and the unobservable lightest SUSY particle of the hidden sector. As the  $\gamma_D$  is expected to be light, it may decay into a lepton pair. No significant excess over the SM expectation is observed and a limit at 95% C.L. is set on the cross section for a benchmark model of supersymmetric hidden-valley Higgs production.

<sup>3</sup> AAD 11AA looked in  $34 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with  $\geq 4$  jets originating from pair production of scalar gluons, each decaying to two gluons. No two-jet resonances are observed over the SM background. Limits are derived on the cross section times branching ratio (see Fig. 3). Assuming 100% branching ratio for the decay to two gluons, the quoted exclusion range is obtained, except for a 5 GeV mass window around 140 GeV.

<sup>4</sup> CHATRCHYAN 11E looked in  $35 \text{ pb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 7 \text{ TeV}$  for events with collimated  $\mu$  pairs (leptonic jets) from the decay of hidden sector states. No evidence for new resonance production is found. Limits are derived and compared to various SUSY models (see Fig. 4) where the LSP, either the  $\tilde{\chi}_1^0$  or a  $\tilde{q}$ , decays to dark sector particles.

<sup>5</sup> ABAZOV 10N looked in  $5.8 \text{ fb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for events from hidden valley models in which a  $\tilde{\chi}_1^0$  decays into a dark photon,  $\gamma_D$ , and the unobservable lightest SUSY particle of the hidden sector. As the  $\gamma_D$  is expected to be light, it may decay into a tightly collimated lepton pair, called lepton jet. They searched for events with  $\cancel{E}_T$  and two isolated lepton jets observable by an opposite charged lepton pair  $e\bar{e}$ ,  $e\mu$  or  $\mu\mu$ . No significant excess over the SM expectation is observed, and a limit at 95% C.L. on the cross section times branching ratio is derived, see their Table I. They also

- examined the invariant mass of the lepton jets for a narrow resonance, see their Fig. 4, but found no evidence for a signal.
- <sup>6</sup> LOVE 08A searched for decays of  $Y(nS)$  with  $n = 1, 2, 3$  into  $\mu\tau$  in 1.1, 1.3, 1.4  $\text{fb}^{-1}$ , respectively, in the CLEO III detector at CESR. The signature is a muon with  $\approx 97\%$  of the beam energy and an electron from the decay of  $\tau$ . No evidence for lepton flavour violation is found and 95% CL limits on the branching ratio are estimated to be 6.0, 14.4 and  $20.3 \times 10^{-6}$  for  $n = 1, 2, 3$ , respectively.
  - <sup>7</sup> ABULENCIA 06P searched in  $305 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.96 \text{ TeV}$  for an excess of events with  $\ell\gamma\cancel{E}_T$  and  $\ell\ell\gamma$  ( $\ell = e, \mu$ ). No significant excess was found compared to the background expectation. No events are found such as the  $e e \gamma \gamma \cancel{E}_T$  event observed in ABE 99I.
  - <sup>8</sup> ACOSTA 04E looked in  $107 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s} = 1.8 \text{ TeV}$  for events with two same sign leptons without selection of other objects nor  $\cancel{E}_T$ . No significant excess is observed compared to the Standard Model expectation and constraints are derived on the parameter space of MSUGRA models, see Figure 4.
  - <sup>9</sup> Looked for the scalar partner of a goldstino in decays  $K^- \rightarrow \pi^- \pi^0 P$  from a 25 GeV  $K^-$  beam produced at the IHEP 70 GeV proton synchrotron. The sgoldstino is assumed to be sufficiently long-lived to be invisible. A 90% CL upper limit on the decay branching ratio is set at  $\sim 9.0 \times 10^{-6}$  for a sgoldstino mass range from 0 to 200 MeV, excluding the interval near  $m(\pi^0)$ , where the limit is  $\sim 3.5 \times 10^{-5}$ .
  - <sup>10</sup> AFFOLDER 02D looked in  $85 \text{ pb}^{-1}$  of  $p\bar{p}$  collisions at  $\sqrt{s}=1.8 \text{ TeV}$  for events with a high- $E_T$  photon, and a  $b$ -tagged jet with or without  $\cancel{E}_T$ . They compared the data with models where the final state could arise from cascade decays of gluinos and/or squarks into  $\tilde{\chi}^\pm$  and  $\tilde{\chi}_2^0$  or direct associated production of  $\tilde{\chi}_2^0 \tilde{\chi}_2^\pm$ , followed by  $\tilde{\chi}_2^0 \rightarrow \gamma \tilde{\chi}_1^0$  or a GMSB model where  $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$ . It is concluded that the experimental sensitivity is insufficient to detect the associated production or the GMSB model, but some sensitivity may exist to the cascade decays. A model independent limit for the above topology is also given in the paper.
  - <sup>11</sup> AFFOLDER 01H searches for  $p\bar{p} \rightarrow \gamma\gamma X$  events, where the di-photon system originates from sgoldstino production, in  $100 \text{ pb}^{-1}$  of data. Upper limits on the cross section times branching ratio are shown as function of the di-photon mass  $>70 \text{ GeV}$  in Fig. 5. Excluded regions are derived in the plane of the sgoldstino mass versus the supersymmetry breaking scale for two representative sets of parameter values, as shown in Figs. 6 and 7.
  - <sup>12</sup> ABBOTT 00G searches for trilepton final states ( $\ell=e,\mu$ ) with  $\cancel{E}_T$  from the indirect decay of gauginos via  $LL\bar{E}$  couplings. Efficiencies are computed for all possible production and decay modes of SUSY particles in the framework of the Minimal Supergravity scenario. See Figs. 1–4 for excluded regions in the  $m_{1/2}$  versus  $m_0$  plane.
  - <sup>13</sup> ABREU,P 00C look for the  $CP$ -even ( $S$ ) and  $CP$ -odd ( $P$ ) scalar partners of the goldstino, expected to be produced in association with a photon. The  $S/P$  decay into two photons or into two gluons and both the tri-photon and the photon + two jets topologies are investigated. Upper limits on the production cross section are shown in Fig. 5 and the excluded regions in Fig. 6. Data collected at  $\sqrt{s}= 189\text{--}202 \text{ GeV}$ .
  - <sup>14</sup> ABACHI 97 searched for  $p\bar{p} \rightarrow \gamma\gamma \cancel{E}_T + X$  as supersymmetry signature. It can be caused by selectron, sneutrino, or neutralino production with a radiative decay of their decay products. They placed limits on cross sections.
  - <sup>15</sup> BARBER 84B consider that  $\tilde{\mu}$  and  $\tilde{e}$  may mix leading to  $\mu \rightarrow e\tilde{\gamma}\tilde{\gamma}$ . They discuss mass-mixing limits from decay dist. asym. in LBL-TRIUMF data and  $e^+$  polarization in SIN data.
  - <sup>16</sup> HOFFMAN 83 set CL = 90% limit  $d\sigma/dt B(e^+e^-) < 3.5 \times 10^{-32} \text{ cm}^2/\text{GeV}^2$  for spin-1 partner of Goldstone fermions with  $140 < m < 160 \text{ MeV}$  decaying  $\rightarrow e^+e^-$  pair.

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CHATRCHYAN	12Q	PL B716 260	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
CHATRCHYAN	12U	PRL 109 071803	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
DAW	12	ASP 35 397	E. Daw <i>et al.</i>	(DRIFT-1Id Collab)
DREINER	12A	EPL 99 61001	H.K. Dreiner, M. Kramer, J. Tattersall	(BONN+)
ELLIS	12B	EPJ C72 2005	J. Ellis, K. Olive	
FELIZARDO	12	PRL 108 201302	M. Felizardo <i>et al.</i>	(SIMPLE Collab.)
FENG	12B	PR D85 075007	J. Feng, K. Matchev, D. Sanford	
FOWLIE	12A	PR D86 075010	A. Fowlie <i>et al.</i>	
KADASTIK	12	JHEP 1205 061	M. Kadastik <i>et al.</i>	
KIM	12	PRL 108 181301	S.C. Kim <i>et al.</i>	(KIMS Collab.)
STREGE	12	JCAP 1203 030	C. Strege <i>et al.</i>	(LOIC, AMST, MADU, GRAN+)
AAD	11AA	EPJ C71 1828	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11AE	JHEP 1110 107	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11AF	JHEP 1111 099	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11B	EPJ C71 1682	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11C	EPJ C71 1647	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11G	PRL 106 131802	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11H	PRL 106 251801	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11K	PL B701 1	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11N	PL B701 186	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11O	PL B701 398	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11P	PL B703 428	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11X	EPJ C71 1744	G. Aad <i>et al.</i>	(ATLAS Collab.)
AAD	11Z	EPJ C71 1809	G. Aad <i>et al.</i>	(ATLAS Collab.)

AALTONEN	11Q	PRL 107 042001	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AARON	11	EPJ C71 1572	E.D. Aaron <i>et al.</i>	(H1 Collab.)
AARON	11C	PL B705 52	F. D. Aaron <i>et al.</i>	(H1 Collab.)
ABAZOV	11N	PL B696 321	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABRAMOWSKI	11	PRL 106 161301	A. Abramowski <i>et al.</i>	(H.E.S.S. Collab.)
AHMED	11A	PR D84 011102	Z. Ahmed <i>et al.</i>	(CDMS and EDELWEISS Collabs.)
APRILE	11B	PRL 107 131302	E. Aprile <i>et al.</i>	(XENON100 Collab.)
ARMENGAUD	11	PL B702 329	E. Armengaud <i>et al.</i>	(EDELWEISS II Collab.)
BUCHMUEL...	11	EPJ C71 1583	O. Buchmueller <i>et al.</i>	
BUCHMUEL...	11B	EPJ C71 1722	O. Buchmueller <i>et al.</i>	
CHATRCHYAN	11AB	JHEP 1106 077	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
CHATRCHYAN	11AC	JHEP 1108 155	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
CHATRCHYAN	11B	JHEP 1106 093	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
CHATRCHYAN	11C	JHEP 1106 026	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
CHATRCHYAN	11D	JHEP 1107 113	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
CHATRCHYAN	11E	JHEP 1107 098	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
CHATRCHYAN	11G	PRL 106 211802	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
CHATRCHYAN	11Q	JHEP 1108 156	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
CHATRCHYAN	11V	PL B704 411	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
CHATRCHYAN	11W	PRL 107 221804	S. Chatrchyan <i>et al.</i>	(CMS Collab.)
KHACHATRY...	11	PRL 106 011801	V. Khachatryan <i>et al.</i>	(CMS Collab.)
KHACHATRY...	11C	JHEP 1103 024	V. Khachatryan <i>et al.</i>	(CMS Collab.)
KHACHATRY...	11I	PL B698 196	V. Khachatryan <i>et al.</i>	(CMS Collab.)
ROSZKOWSKI	11	PR D83 015014	L. Roszkowski <i>et al.</i>	
AALTONEN	10	PRL 104 011801	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	10O	PRL 104 251801	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	10R	PRL 105 081802	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	10Y	PR D82 092001	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	10Z	PRL 105 191801	T. Aaltonen <i>et al.</i>	(CDF Collab.)
ABAZOV	10L	PL B693 95	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	10M	PRL 105 191802	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	10N	PRL 105 211802	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	10P	PRL 105 221802	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABDO	10	JCAP 1004 014	A.A. Abdo <i>et al.</i>	(Fermi-LAT Collab.)
ACKERMANN	10	JCAP 1005 025	M. Ackermann	(Fermi-LAT Collab.)
AHMED	10	SCI 327 1619	Z. Ahmed <i>et al.</i>	(CDMS II Collab.)
ARMENGAUD	10	PL B687 294	E. Armengaud <i>et al.</i>	(EDELWEISS II Collab.)
BERGER	10	PR D82 114023	E.L. Berger <i>et al.</i>	
ELLIS	10	EPJ C69 201	J. Ellis, A. Mustafayev, K. Olive	
AALTONEN	09G	PR D79 052004	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	09R	PRL 102 221801	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	09S	PRL 102 121801	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	09V	PRL 102 091805	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	09Z	PRL 103 021802	T. Aaltonen <i>et al.</i>	(CDF Collab.)
ABAZOV	09M	PRL 102 161802	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	09N	PL B674 4	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	09O	PL B675 289	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	09S	PL B680 24	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	09T	PL B680 34	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABBASI	09B	PRL 102 201302	R. Abbasi <i>et al.</i>	(IceCube Collab.)
AHMED	09	PRL 102 011301	Z. Ahmed <i>et al.</i>	(CDMS Collab.)
ANGLOHER	09	ASP 31 270	G. Angloher <i>et al.</i>	(CRESST Collab.)
ARCHAMBAU...	09	PL B682 185	S. Archambault <i>et al.</i>	(PICASSO Collab.)
BUCHMUEL...	09	EPJ C64 391	O. Buchmueller <i>et al.</i>	(LOIC, FNAL, CERN+)
DREINER	09	EPJ C62 547	H. Dreiner <i>et al.</i>	
LEBEDENKO	09	PR D80 052010	V.N. Lebedenko <i>et al.</i>	(ZEPLIN-III Collab.)
LEBEDENKO	09A	PRL 103 151302	V.N. Lebedenko <i>et al.</i>	(ZEPLIN-III Collab.)
SORENSEN	09	NIM A601 339	P. Sorensen <i>et al.</i>	(XENON10 Collab.)
AALTONEN	08AE	PRL 101 251801	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	08L	PR D77 052002	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	08U	PR D78 032015	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	08Z	PRL 101 071802	T. Aaltonen <i>et al.</i>	(CDF Collab.)
ABAZOV	08	PL B659 500	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	08F	PL B659 856	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	08G	PL B660 449	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	08Q	PRL 100 241803	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	08X	PRL 101 111802	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	08Z	PL B665 1	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ANGLE	08	PRL 100 021303	J. Angle <i>et al.</i>	(XENON10 Collab.)
ANGLE	08A	PRL 101 091301	J. Angle <i>et al.</i>	(XENON10 Collab.)

BEDNYAKOV	08	PAN 71 111	V.A. Bednyakov, H.P. Klapdor-Kleingrothaus, I.V. Krivosheina	
		Translated from YAF 71 112.		
BEHNKE	08	SCI 319 933	E. Behnke	(COUPP Collab.)
BENETTI	08	ASP 28 495	P. Benetti <i>et al.</i>	(WARP Collab.)
BUCHMUELLER...	08	JHEP 0809 117	O. Buchmueller <i>et al.</i>	
ELLIS	08	PR D78 075012	J. Ellis, K. Olive, P. Sandick	(CERN, MINN)
KAPLAN	08	PRL 101 022002	D.E. Kaplan, M.D. Schwartz	
LOVE	08A	PRL 101 201601	W. Love <i>et al.</i>	(CLEO Collab.)
AALTONEN	07E	PR D76 072010	T. Aaltonen <i>et al.</i>	(CDF Collab.)
AALTONEN	07J	PRL 99 191806	T. Aaltonen <i>et al.</i>	(CDF Collab.)
ABAZOV	07B	PL B645 119	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	07L	PRL 99 131801	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABULENCIA	07H	PRL 98 131804	A. Abulencia <i>et al.</i>	(CDF Collab.)
ABULENCIA	07N	PRL 98 221803	A. Abulencia <i>et al.</i>	(CDF Collab.)
ABULENCIA	07P	PRL 99 121801	A. Abulencia <i>et al.</i>	(CDF Collab.)
ALNER	07A	ASP 28 287	G.J. Alner <i>et al.</i>	(ZEPLIN-II Collab.)
CALIBBI	07	JHEP 0709 081	L. Calibbi <i>et al.</i>	
CHEKANOV	07	EPJ C50 269	S. Chekanov <i>et al.</i>	(ZEUS Collab.)
ELLIS	07	JHEP 0706 079	J. Ellis, K. Olive, P. Sandick	(CERN, MINN)
LEE	07A	PRL 99 091301	H.S. Lee <i>et al.</i>	(KIMS Collab.)
SCHAEEL	07A	EPJ C49 411	S. Schaeel <i>et al.</i>	(ALEPH Collab.)
ABAZOV	06C	PL B638 119	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	06D	PL B638 441	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	06I	PRL 97 111801	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	06P	PRL 97 161802	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	06R	PRL 97 171806	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABBIENDI	06B	EPJ C46 307	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABDALLAH	06C	EPJ C45 589	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABULENCIA	06I	PRL 96 171802	A. Abulencia <i>et al.</i>	(CDF Collab.)
ABULENCIA	06M	PRL 96 211802	A. Abulencia <i>et al.</i>	(CDF Collab.)
ABULENCIA	06P	PRL 97 031801	A. Abulencia <i>et al.</i>	(CDF Collab.)
ACHTERBERG	06	ASP 26 129	A. Achterberg <i>et al.</i>	(AMANDA Collab.)
ACKERMANN	06	ASP 24 459	M. Ackermann <i>et al.</i>	(AMANDA Collab.)
AKERIB	06	PR D73 011102	D.S. Akerib <i>et al.</i>	(CDMS Collab.)
AKERIB	06A	PRL 96 011302	D.S. Akerib <i>et al.</i>	(CDMS Collab.)
ALLANACH	06	PR D73 015013	B.C. Allanach <i>et al.</i>	
BENOIT	06	PL B637 156	A. Benoit <i>et al.</i>	
DE-AUSTRI	06	JHEP 0605 002	R.R. de Austri, R. Trotta, L. Roszkowski	
DEBOER	06	PL B636 13	W. de Boer <i>et al.</i>	
LEP-SLC	06	PRPL 427 257	ALEPH, DELPHI, L3, OPAL, SLD and working groups	
SHIMIZU	06A	PL B633 195	Y. Shimizu <i>et al.</i>	
SMITH	06	PL B642 567	N.J.T. Smith, A.S. Murphy, T.J. Summer	
ABAZOV	05A	PRL 94 041801	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	05U	PRL 95 151805	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABDALLAH	05B	EPJ C38 395	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABULENCIA	05A	PRL 95 252001	A. Abulencia <i>et al.</i>	(CDF Collab.)
ACOSTA	05E	PR D71 031104	D. Acosta <i>et al.</i>	(CDF Collab.)
ACOSTA	05R	PRL 95 131801	D. Acosta <i>et al.</i>	(CDF Collab.)
AKERIB	05	PR D72 052009	D.S. Akerib <i>et al.</i>	(CDMS Collab.)
AKTAS	05	PL B616 31	A. Aktas <i>et al.</i>	(H1 Collab.)
ALNER	05	PL B616 17	G.J. Alner <i>et al.</i>	(UK Dark Matter Collab.)
ALNER	05A	ASP 23 444	G.J. Alner <i>et al.</i>	(UK Dark Matter Collab.)
ANGLOHER	05	ASP 23 325	G. Angloher <i>et al.</i>	(CRESST-II Collab.)
BAER	05	JHEP 0507 065	H. Baer <i>et al.</i>	(FSU, MSU, HAWA)
BARNABE-HEIDER...	05	PL B624 186	M. Barnabe-Heider <i>et al.</i>	(PICASSO Collab.)
BERGER	05	PR D71 014007	E.L. Berger <i>et al.</i>	
CHEKANOV	05A	EPJ C44 463	S. Chekanov <i>et al.</i>	(ZEUS Collab.)
ELLIS	05	PR D71 095007	J. Ellis <i>et al.</i>	
SANGLARD	05	PR D71 122002	V. Sanglard <i>et al.</i>	(EDELWEISS Collab.)
ABAZOV	04	PL B581 147	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	04B	PRL 93 011801	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABBIENDI	04	EPJ C32 453	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	04F	EPJ C33 149	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	04H	EPJ C35 1	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	04N	PL B602 167	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABDALLAH	04H	EPJ C34 145	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABDALLAH	04M	EPJ C36 1	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
Also		EPJ C37 129 (errata)	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ACHARD	04	PL B580 37	P. Achard <i>et al.</i>	(L3 Collab.)
ACHARD	04E	PL B587 16	P. Achard <i>et al.</i>	(L3 Collab.)
ACOSTA	04B	PRL 92 051803	D. Acosta <i>et al.</i>	(CDF Collab.)

ACOSTA	04E	PRL 93 061802	D. Acosta <i>et al.</i>	(CDF Collab.)
AKERIB	04	PRL 93 211301	D. Akerib <i>et al.</i>	(CDMSII Collab.)
AKTAS	04B	PL B599 159	A. Aktas <i>et al.</i>	(H1 Collab.)
AKTAS	04D	EPJ C36 425	A. Aktas <i>et al.</i>	(H1 Collab.)
BALTZ	04	JHEP 0410 052	E. Baltz, P. Gondolo	
BELANGER	04	JHEP 0403 012	G. Belanger <i>et al.</i>	
BOTTINO	04	PR D69 037302	A. Bottino <i>et al.</i>	
DAS	04	PL B596 293	S.P. Das, A. Datta, M. Maity	
DESAI	04	PR D70 083523	S. Desai <i>et al.</i>	(Super-Kamiokande Collab.)
ELLIS	04	PR D69 015005	J. Ellis <i>et al.</i>	
ELLIS	04B	PR D70 055005	J. Ellis <i>et al.</i>	
HEISTER	04	PL B583 247	A. Heister <i>et al.</i>	(ALEPH Collab.)
JANOT	04	PL B594 23	P. Janot	
PIERCE	04A	PR D70 075006	A. Pierce	
TCHIKILEV	04	PL B602 149	O.G. Tchikilev <i>et al.</i>	(ISTRA+ Collab.)
ABBIENDI	03H	EPJ C29 479	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	03L	PL B572 8	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABDALLAH	03C	EPJ C26 505	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABDALLAH	03D	EPJ C27 153	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABDALLAH	03F	EPJ C28 15	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABDALLAH	03G	EPJ C29 285	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ABDALLAH	03M	EPJ C31 421	J. Abdallah <i>et al.</i>	(DELPHI Collab.)
ACOSTA	03C	PRL 90 251801	D. Acosta <i>et al.</i>	(CDF Collab.)
ACOSTA	03E	PRL 91 171602	D. Acosta <i>et al.</i>	(CDF Collab.)
ADLOFF	03	PL B568 35	C. Adloff <i>et al.</i>	(H1 Collab.)
AHMED	03	ASP 19 691	B. Ahmed <i>et al.</i>	(UK Dark Matter Collab.)
AKERIB	03	PR D68 082002	D. Akerib <i>et al.</i>	(CDMS Collab.)
BAER	03	JCAP 0305 006	H. Baer, C. Balazs	
BAER	03A	JCAP 0309 007	H. Baer <i>et al.</i>	
BERGER	03	PL B552 223	E. Berger <i>et al.</i>	
BOTTINO	03	PR D68 043506	A. Bottino <i>et al.</i>	
BOTTINO	03A	PR D67 063519	A. Bottino, N. Fornengo, S. Scopel	
CHAKRAB...	03	PR D68 015005	S. Chakrabarti, M. Guhait, N.K. Mondal	
CHATTOPAD...	03	PR D68 035005	U. Chattopadhyay, A. Corsetti, P. Nath	
CHEKANOV	03B	PR D68 052004	S. Chekanov <i>et al.</i>	(ZEUS Collab.)
ELLIS	03	ASP 18 395	J. Ellis, K.A. Olive, Y. Santoso	
ELLIS	03B	NP B652 259	J. Ellis <i>et al.</i>	
ELLIS	03C	PL B565 176	J. Ellis <i>et al.</i>	
ELLIS	03D	PL B573 162	J. Ellis <i>et al.</i>	
ELLIS	03E	PR D67 123502	J. Ellis <i>et al.</i>	
HEISTER	03	EPJ C27 1	A. Heister <i>et al.</i>	(ALEPH)
HEISTER	03C	EPJ C28 1	A. Heister <i>et al.</i>	(ALEPH Collab.)
HEISTER	03G	EPJ C31 1	A. Heister <i>et al.</i>	(ALEPH Collab.)
HEISTER	03H	EPJ C31 327	A. Heister <i>et al.</i>	(ALEPH Collab.)
JANOT	03	PL B564 183	P. Janot	
KLAPDOR-K...	03	ASP 18 525	H.V. Klapdor-Kleingrothaus <i>et al.</i>	
LAHANAS	03	PL B568 55	A. Lahanas, D. Nanopoulos	
TAKEDA	03	PL B572 145	A. Takeda <i>et al.</i>	
ABAZOV	02C	PRL 88 171802	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	02F	PRL 89 171801	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	02G	PR D66 112001	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABAZOV	02H	PRL 89 261801	V.M. Abazov <i>et al.</i>	(D0 Collab.)
ABBIENDI	02	EPJ C23 1	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	02B	PL B526 233	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	02H	PL B545 272	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
Also		PL B548 258 (errata)	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABRAMS	02	PR D66 122003	D. Abrams <i>et al.</i>	(CDMS Collab.)
ACHARD	02	PL B524 65	P. Achard <i>et al.</i>	(L3 Collab.)
ACOSTA	02H	PRL 89 281801	D. Acosta <i>et al.</i>	(CDF Collab.)
AFFOLDER	02	PRL 88 041801	T. Affolder <i>et al.</i>	(CDF Collab.)
AFFOLDER	02D	PR D65 052006	T. Affolder <i>et al.</i>	(CDF Collab.)
ANGLOHER	02	ASP 18 43	G. Angloher <i>et al.</i>	(CRESST Collab.)
ARNOWITT	02	hep-ph/0211417	R. Arnowitt, B. Dutta	
BAEK	02	PL B541 161	S. Baek	
BECHER	02	PL B540 278	T. Becher <i>et al.</i>	
BENOIT	02	PL B545 43	A. Benoit <i>et al.</i>	(EDELWEISS Collab.)
CHEKANOV	02	PR D65 092004	S. Chekanov <i>et al.</i>	(ZEUS Collab.)
CHEUNG	02B	PRL 89 221801	K. Cheung, W.-Y. Keung	
CHO	02	PRL 89 091801	G.-C. Cho	
ELLIS	02B	PL B532 318	J. Ellis, A. Ferstl, K.A. Olive	
GHODBANE	02	NP B647 190	N. Ghodbane <i>et al.</i>	

HEISTER	02	PL B526 191	A. Heister <i>et al.</i>	(ALEPH Collab.)
HEISTER	02E	PL B526 206	A. Heister <i>et al.</i>	(ALEPH Collab.)
HEISTER	02F	EPJ C25 1	A. Heister <i>et al.</i>	(ALEPH Collab.)
HEISTER	02J	PL B533 223	A. Heister <i>et al.</i>	(ALEPH Collab.)
HEISTER	02K	PL B537 5	A. Heister <i>et al.</i>	(ALEPH Collab.)
HEISTER	02N	PL B544 73	A. Heister <i>et al.</i>	(ALEPH Collab.)
HEISTER	02R	EPJ C25 339	A. Heister <i>et al.</i>	(ALEPH Collab.)
KIM	02	PL B527 18	H.B. Kim <i>et al.</i>	
KIM	02B	JHEP 0212 034	Y.G. Kim <i>et al.</i>	
LAHANAS	02	EPJ C23 185	A. Lahanas, V.C. Spanos	
MORALES	02B	ASP 16 325	A. Morales <i>et al.</i>	(COSME Collab.)
MORALES	02C	PL B532 8	A. Morales <i>et al.</i>	(IGEX Collab.)
ABBIENDI	01	PL B501 12	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBOTT	01D	PR D63 091102	B. Abbott <i>et al.</i>	(D0 Collab.)
ABREU	01	EPJ C19 29	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	01B	EPJ C19 201	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	01C	PL B502 24	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	01D	PL B500 22	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	01G	PL B503 34	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACCIARRI	01	EPJ C19 397	M. Acciarri <i>et al.</i>	(L3 Collab.)
ADAMS	01	PRL 87 041801	T. Adams <i>et al.</i>	(NuTeV Collab.)
ADLOFF	01B	EPJ C20 639	C. Adloff <i>et al.</i>	(H1 Collab.)
AFFOLDER	01B	PR D63 091101	T. Affolder <i>et al.</i>	(CDF Collab.)
AFFOLDER	01H	PR D64 092002	T. Affolder <i>et al.</i>	(CDF Collab.)
AFFOLDER	01J	PRL 87 251803	T. Affolder <i>et al.</i>	(CDF Collab.)
BALTZ	01	PRL 86 5004	E. Baltz, P. Gondolo	
BARATE	01	PL B499 67	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	01B	EPJ C19 415	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARGER	01C	PL B518 117	V. Barger, C. Kao	
BAUDIS	01	PR D63 022001	L. Baudis <i>et al.</i>	(Heidelberg-Moscow Collab.)
BENOIT	01	PL B513 15	A. Benoit <i>et al.</i>	(EDELWEISS Collab.)
BERGER	01	PRL 86 4231	E. Berger <i>et al.</i>	
BERNABEI	01	PL B509 197	R. Bernabei <i>et al.</i>	(DAMA Collab.)
BOTTINO	01	PR D63 125003	A. Bottino <i>et al.</i>	
BREITWEG	01	PR D63 052002	J. Breitweg <i>et al.</i>	(ZEUS Collab.)
CORSETTI	01	PR D64 125010	A. Corsetti, P. Nath	
ELLIS	01B	PL B510 236	J. Ellis <i>et al.</i>	
ELLIS	01C	PR D63 065016	J. Ellis, A. Ferstl, K.A. Olive	
GOMEZ	01	PL B512 252	M.E. Gomez, J.D. Vergados	
LAHANAS	01	PL B518 94	A. Lahanas, D.V. Nanopoulos, V. Spanos	
SAVINOV	01	PR D63 051101	V. Savinov <i>et al.</i>	(CLEO Collab.)
ABBIENDI	00	EPJ C12 1	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	00G	EPJ C14 51	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	00H	EPJ C14 187	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
Also		EPJ C16 707 (errata)	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	00J	EPJ C12 551	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	00R	EPJ C13 553	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI,G	00D	EPJ C18 253	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBOTT	00C	PRL 84 2088	B. Abbott <i>et al.</i>	(D0 Collab.)
ABBOTT	00G	PR D62 071701	B. Abbott <i>et al.</i>	(D0 Collab.)
ABREU	00I	EPJ C13 591	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00J	PL B479 129	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00Q	PL B478 65	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00S	PL B485 45	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00T	PL B485 95	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00U	PL B487 36	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00V	EPJ C16 211	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00W	PL B489 38	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	00Z	EPJ C17 53	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU,P	00C	PL B494 203	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU,P	00D	PL B496 59	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABUSAIDI	00	PRL 84 5699	R. Abusaidi <i>et al.</i>	(CDMS Collab.)
ACCIARRI	00C	EPJ C16 1	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	00D	PL B472 420	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	00K	PL B482 31	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	00P	PL B489 81	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCOMANDO	00	NP B585 124	E. Accomando <i>et al.</i>	
AFFOLDER	00D	PRL 84 5704	T. Affolder <i>et al.</i>	(CDF Collab.)
AFFOLDER	00G	PRL 84 5273	T. Affolder <i>et al.</i>	(CDF Collab.)
AFFOLDER	00J	PRL 85 1378	T. Affolder <i>et al.</i>	(CDF Collab.)
AFFOLDER	00K	PRL 85 2056	T. Affolder <i>et al.</i>	(CDF Collab.)

BARATE	00G	EPJ C16 71	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	00H	EPJ C13 29	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	00I	EPJ C12 183	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	00P	PL B488 234	R. Barate <i>et al.</i>	(ALEPH Collab.)
BERNABEI	00	PL B480 23	R. Bernabei <i>et al.</i>	(DAMA Collab.)
BERNABEI	00C	EPJ C18 283	R. Bernabei <i>et al.</i>	(DAMA Collab.)
BERNABEI	00D	NJP 2 15	R. Bernabei <i>et al.</i>	(DAMA Collab.)
BOEHM	00B	PR D62 035012	C. Boehm, A. Djouadi, M. Drees	
BREITWEG	00E	EPJ C16 253	J. Breitweg <i>et al.</i>	(ZEUS Collab.)
CHO	00B	NP B574 623	G.-C. Cho, K. Hagiwara	
ELLIS	00	PR D62 075010	J. Ellis <i>et al.</i>	
FENG	00	PL B482 388	J.L. Feng, K.T. Matchev, F. Wilczek	
LEP	00	CERN-EP-2000-016	LEP Collabs.	(ALEPH, DELPHI, L3, OPAL, SLD+)
MAFI	00	PR D62 035003	A. Mafi, S. Raby	
MALTONI	00	PL B476 107	M. Maltoni <i>et al.</i>	
MORALES	00	PL B489 268	A. Morales <i>et al.</i>	(IGEX Collab.)
PDG	00	EPJ C15 1	D.E. Groom <i>et al.</i>	(PDG Collab.)
SPOONER	00	PL B473 330	N.J.C. Spooner <i>et al.</i>	(UK Dark Matter Col.)
ABBIENDI	99	EPJ C6 1	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	99F	EPJ C8 23	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	99M	PL B456 95	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBIENDI	99T	EPJ C11 619	G. Abbiendi <i>et al.</i>	(OPAL Collab.)
ABBOTT	99	PRL 82 29	B. Abbott <i>et al.</i>	(D0 Collab.)
ABBOTT	99F	PR D60 031101	B. Abbott <i>et al.</i>	(D0 Collab.)
ABBOTT	99J	PRL 83 2896	B. Abbott <i>et al.</i>	(D0 Collab.)
ABBOTT	99K	PRL 83 4476	B. Abbott <i>et al.</i>	(D0 Collab.)
ABBOTT	99L	PRL 83 4937	B. Abbott <i>et al.</i>	(D0 Collab.)
ABE	99I	PR D59 092002	F. Abe <i>et al.</i>	(CDF Collab.)
ABE	99M	PRL 83 2133	F. Abe <i>et al.</i>	(CDF Collab.)
ABREU	99A	EPJ C11 383	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	99C	EPJ C6 385	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	99F	EPJ C7 595	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	99G	PL B446 62	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACCIARRI	99H	PL B456 283	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	99I	PL B459 354	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	99L	PL B462 354	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	99R	PL B470 268	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	99V	PL B471 308	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	99W	PL B471 280	M. Acciarri <i>et al.</i>	(L3 Collab.)
ALAVI-HARATI	99E	PRL 83 2128	A. Alavi-Harati <i>et al.</i>	(FNAL KTeV Collab.)
AMBROSIO	99	PR D60 082002	M. Ambrosio <i>et al.</i>	(Macro Collab.)
BAER	99	PR D59 075002	H. Baer, K. Cheung, J.F. Gunion	
BARATE	99E	EPJ C7 383	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	99Q	PL B469 303	R. Barate <i>et al.</i>	(ALEPH Collab.)
BAUDIS	99	PR D59 022001	L. Baudis <i>et al.</i>	(Heidelberg-Moscow Collab.)
BELLI	99C	NP B563 97	P. Belli <i>et al.</i>	(DAMA Collab.)
BERNABEI	99	PL B450 448	R. Bernabei <i>et al.</i>	(DAMA Collab.)
FANTI	99	PL B446 117	V. Fanti <i>et al.</i>	(CERN NA48 Collab.)
MALTONI	99B	PL B463 230	M. Maltoni, M.I. Vysotsky	
OOTANI	99	PL B461 371	W. Ootani <i>et al.</i>	
ABBOTT	98	PRL 80 442	B. Abbott <i>et al.</i>	(D0 Collab.)
ABBOTT	98C	PRL 80 1591	B. Abbott <i>et al.</i>	(D0 Collab.)
ABBOTT	98E	PRL 80 2051	B. Abbott <i>et al.</i>	(D0 Collab.)
ABBOTT	98J	PRL 81 38	B. Abbott <i>et al.</i>	(D0 Collab.)
ABE	98J	PRL 80 5275	F. Abe <i>et al.</i>	(CDF Collab.)
ABE	98S	PRL 81 4806	F. Abe <i>et al.</i>	(CDF Collab.)
ABREU	98	EPJ C1 1	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ABREU	98P	PL B444 491	P. Abreu <i>et al.</i>	(DELPHI Collab.)
ACCIARRI	98F	EPJ C4 207	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	98J	PL B433 163	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACCIARRI	98V	PL B444 503	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACKERSTAFF	98K	EPJ C4 47	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
ACKERSTAFF	98L	EPJ C2 213	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
ACKERSTAFF	98P	PL B433 195	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
ACKERSTAFF	98V	EPJ C2 441	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
BARATE	98H	PL B420 127	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	98J	PL B429 201	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	98K	PL B433 176	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	98S	EPJ C4 433	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	98X	EPJ C2 417	R. Barate <i>et al.</i>	(ALEPH Collab.)
BERNABEI	98	PL B424 195	R. Bernabei <i>et al.</i>	(DAMA Collab.)

BERNABEI	98C	PL B436 379	R. Bernabei <i>et al.</i>	(DAMA Collab.)
BREITWEG	98	PL B434 214	J. Breitweg <i>et al.</i>	(ZEUS Collab.)
ELLIS	98	PR D58 095002	J. Ellis <i>et al.</i>	
ELLIS	98B	PL B444 367	J. Ellis, T. Falk, K. Olive	
PDG	98	EPJ C3 1	C. Caso <i>et al.</i>	(PDG Collab.)
ABACHI	97	PRL 78 2070	S. Abachi <i>et al.</i>	(D0 Collab.)
ABBOTT	97B	PRL 79 4321	B. Abbott <i>et al.</i>	(D0 Collab.)
ABE	97K	PR D56 R1357	F. Abe <i>et al.</i>	(CDF Collab.)
ACCIARRI	97U	PL B414 373	M. Acciarri <i>et al.</i>	(L3 Collab.)
ACKERSTAFF	97H	PL B396 301	K. Ackerstaff <i>et al.</i>	(OPAL Collab.)
ADAMS	97B	PRL 79 4083	J. Adams <i>et al.</i>	(FNAL KTeV Collab.)
ALBUQUERQ...	97	PRL 78 3252	I.F. Albuquerque <i>et al.</i>	(FNAL E761 Collab.)
BAER	97	PR D57 567	H. Baer, M. Brhlik	
BARATE	97K	PL B405 379	R. Barate <i>et al.</i>	(ALEPH Collab.)
BARATE	97L	ZPHY C76 1	R. Barate <i>et al.</i>	(ALEPH Collab.)
BERNABEI	97	ASP 7 73	R. Bernabei <i>et al.</i>	(DAMA Collab.)
CARENA	97	PL B390 234	M. Carena, G.F. Giudice, C.E.M. Wagner	
CSIKOR	97	PRL 78 4335	F. Csikor, Z. Fodor	(EOTV, CERN)
DATTA	97	PL B395 54	A. Datta, M. Guchait, N. Parua	(ICTP, TATA)
DEGOUVEA	97	PL B400 117	A. de Gouvea, H. Murayama	
DERRICK	97	ZPHY C73 613	M. Derrick <i>et al.</i>	(ZEUS Collab.)
EDSJO	97	PR D56 1879	J. Edsjo, P. Gondolo	
ELLIS	97	PL B394 354	J. Ellis, J.L. Lopez, D.V. Nanopoulos	
HEWETT	97	PR D56 5703	J.L. Hewett, T.G. Rizzo, M.A. Doncheski	
KALINOWSKI	97	PL B400 112	J. Kalinowski, P. Zerwas	
TEREKHOV	97	PL B412 86	I. Terekhov	(ALAT)
ABACHI	96	PRL 76 2228	S. Abachi <i>et al.</i>	(D0 Collab.)
ABACHI	96B	PRL 76 2222	S. Abachi <i>et al.</i>	(D0 Collab.)
ABE	96	PRL 77 438	F. Abe <i>et al.</i>	(CDF Collab.)
ABE	96D	PRL 76 2006	F. Abe <i>et al.</i>	(CDF Collab.)
ABE	96K	PRL 76 4307	F. Abe <i>et al.</i>	(CDF Collab.)
AID	96	ZPHY C71 211	S. Aid <i>et al.</i>	(H1 Collab.)
AID	96C	PL B380 461	S. Aid <i>et al.</i>	(H1 Collab.)
ARNOWITT	96	PR D54 2374	R. Arnowitt, P. Nath	
BAER	96	PR D53 597	H. Baer, M. Brhlik	
BERGSTROM	96	ASP 5 263	L. Bergstrom, P. Gondolo	
CHO	96	PL B372 101	G.C. Cho, Y. Kizukuri, N. Oshimo	(TOKAH, OCH)
FARRAR	96	PRL 76 4111	G.R. Farrar	(RUTG)
LEWIN	96	ASP 6 87	J.D. Lewin, P.F. Smith	
TEREKHOV	96	PL B385 139	I. Terkhov, L. Clavelli	(ALAT)
ABACHI	95C	PRL 75 618	S. Abachi <i>et al.</i>	(D0 Collab.)
ABE	95N	PRL 74 3538	F. Abe <i>et al.</i>	(CDF Collab.)
ABE	95T	PRL 75 613	F. Abe <i>et al.</i>	(CDF Collab.)
ACCIARRI	95E	PL B350 109	M. Acciarri <i>et al.</i>	(L3 Collab.)
AKERS	95A	ZPHY C65 367	R. Akers <i>et al.</i>	(OPAL Collab.)
AKERS	95R	ZPHY C67 203	R. Akers <i>et al.</i>	(OPAL Collab.)
BEREZINSKY	95	ASP 5 1	V. Berezinsky <i>et al.</i>	
BUSKULIC	95E	PL B349 238	D. Buskalic <i>et al.</i>	(ALEPH Collab.)
CLAVELLI	95	PR D51 1117	L. Clavelli, P.W. Coulter	(ALAT)
FALK	95	PL B354 99	T. Falk, K.A. Olive, M. Srednicki	(MINN, UCSB)
LOSECCO	95	PL B342 392	J.M. LoSecco	(NDAM)
AKERS	94K	PL B337 207	R. Akers <i>et al.</i>	(OPAL Collab.)
BECK	94	PL B336 141	M. Beck <i>et al.</i>	(MPIH, KIAE, SASSO)
CAKIR	94	PR D50 3268	M.B. Cakir, G.R. Farrar	(RUTG)
FALK	94	PL B339 248	T. Falk, K.A. Olive, M. Srednicki	(UCSB, MINN)
SHIRAI	94	PRL 72 3313	J. Shirai <i>et al.</i>	(VENUS Collab.)
ADRIANI	93M	PRPL 236 1	O. Adriani <i>et al.</i>	(L3 Collab.)
ALITTI	93	NP B400 3	J. Alitti <i>et al.</i>	(UA2 Collab.)
CLAVELLI	93	PR D47 1973	L. Clavelli, P.W. Coulter, K.J. Yuan	(ALAT)
DREES	93	PR D47 376	M. Drees, M.M. Nojiri	(DESY, SLAC)
DREES	93B	PR D48 3483	M. Drees, M.M. Nojiri	
FALK	93	PL B318 354	T. Falk <i>et al.</i>	(UCB, UCSB, MINN)
HEBBEKER	93	ZPHY C60 63	T. Hebbeker	(CERN)
KELLEY	93	PR D47 2461	S. Kelley <i>et al.</i>	(TAMU, ALAH)
LOPEZ	93C	PL B313 241	J.L. Lopez, D.V. Nanopoulos, X. Wang	(TAMU, HARC+)
MIZUTA	93	PL B298 120	S. Mizuta, M. Yamaguchi	(TOHO)
MORI	93	PR D48 5505	M. Mori <i>et al.</i>	(KEK, NIIG, TOKY, TOKA+)
ABE	92L	PRL 69 3439	F. Abe <i>et al.</i>	(CDF Collab.)
BOTTINO	92	MPL A7 733	A. Bottino <i>et al.</i>	(TORI, ZARA)
Also		PL B265 57	A. Bottino <i>et al.</i>	(TORI, INFN)
CLAVELLI	92	PR D46 2112	L. Clavelli	(ALAT)

DECAMP	92	PRPL 216 253	D. Decamp <i>et al.</i>	(ALEPH Collab.)
LOPEZ	92	NP B370 445	J.L. Lopez, D.V. Nanopoulos, K.J. Yuan	(TAMU)
MCDONALD	92	PL B283 80	J. McDonald, K.A. Olive, M. Srednicki	(LISB+)
ROY	92	PL B283 270	D.P. Roy	(CERN)
ABREU	91F	NP B367 511	P. Abreu <i>et al.</i>	(DELPHI Collab.)
AKESSON	91	ZPHY C52 219	T. Akesson <i>et al.</i>	(HELIOS Collab.)
ALEXANDER	91F	ZPHY C52 175	G. Alexander <i>et al.</i>	(OPAL Collab.)
ANTONIADIS	91	PL B262 109	I. Antoniadis, J. Ellis, D.V. Nanopoulos	(EPOL+)
BOTTINO	91	PL B265 57	A. Bottino <i>et al.</i>	(TORI, INFN)
GELMINI	91	NP B351 623	G.B. Gelmini, P. Gondolo, E. Roulet	(UCLA, TRST)
GRIEST	91	PR D43 3191	K. Griest, D. Seckel	
KAMIONKOW...	91	PR D44 3021	M. Kamionkowski	(CHIC, FNAL)
MORI	91B	PL B270 89	M. Mori <i>et al.</i>	(Kamiokande Collab.)
NOJIRI	91	PL B261 76	M.M. Nojiri	(KEK)
OLIVE	91	NP B355 208	K.A. Olive, M. Srednicki	(MINN, UCSB)
ROSZKOWSKI	91	PL B262 59	L. Roszkowski	(CERN)
SATO	91	PR D44 2220	N. Sato <i>et al.</i>	(Kamiokande Collab.)
ADACHI	90C	PL B244 352	I. Adachi <i>et al.</i>	(TOPAZ Collab.)
GRIEST	90	PR D41 3565	K. Griest, M. Kamionkowski, M.S. Turner	(UCB+)
BARBIERI	89C	NP B313 725	R. Barbieri, M. Frigeni, G. Giudice	
NAKAMURA	89	PR D39 1261	T.T. Nakamura <i>et al.</i>	(KYOT, TMTC)
OLIVE	89	PL B230 78	K.A. Olive, M. Srednicki	(MINN, UCSB)
ELLIS	88D	NP B307 883	J. Ellis, R. Flores	
GRIEST	88B	PR D38 2357	K. Griest	
OLIVE	88	PL B205 553	K.A. Olive, M. Srednicki	(MINN, UCSB)
SREDNICKI	88	NP B310 693	M. Srednicki, R. Watkins, K.A. Olive	(MINN, UCSB)
ALBAJAR	87D	PL B198 261	C. Albajar <i>et al.</i>	(UA1 Collab.)
ANSARI	87D	PL B195 613	R. Ansari <i>et al.</i>	(UA2 Collab.)
ARNOLD	87	PL B186 435	R.G. Arnold <i>et al.</i>	(BRUX, DUUC, LOUC+)
NG	87	PL B188 138	K.W. Ng, K.A. Olive, M. Srednicki	(MINN, UCSB)
TUTS	87	PL B186 233	P.M. Tuts <i>et al.</i>	(CUSB Collab.)
ALBRECHT	86C	PL 167B 360	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
BADIER	86	ZPHY C31 21	J. Badier <i>et al.</i>	(NA3 Collab.)
BARNETT	86	NP B267 625	R.M. Barnett, H.E. Haber, G.L. Kane	(LBL, UCSC+)
GAISSER	86	PR D34 2206	T.K. Gaisser, G. Steigman, S. Tilav	(BART, DELA)
VOLOSHIN	86	SJNP 43 495	M.B. Voloshin, L.B. Okun	(ITEP)
COOPER-...	85B	PL 160B 212	Translated from YAF 43 779.	
DAWSON	85	PR D31 1581	A.M. Cooper-Sarkar <i>et al.</i>	(WA66 Collab.)
FARRAR	85	PRL 55 895	S. Dawson, E. Eichten, C. Quigg	(LBL, FNAL)
GOLDMAN	85	Physica 15D 181	G.R. Farrar	(RUTG)
HABER	85	PRPL 117 75	T. Goldman, H.E. Haber	(LANL, UCSC)
BALL	84	PRL 53 1314	H.E. Haber, G.L. Kane	(UCSC, MICH)
BARBER	84B	PL 139B 427	R.C. Ball <i>et al.</i>	(MICH, FIRZ, OSU, FNAL+)
BRICK	84	PR D30 1134	J.S. Barber, R.E. Shrock	(STON)
ELLIS	84	NP B238 453	D.H. Brick <i>et al.</i>	(BROW, CAVE, IIT+)
FARRAR	84	PRL 53 1029	J. Ellis <i>et al.</i>	(CERN)
BERGMA	83C	PL 121B 429	G.R. Farrar	(RUTG)
CHANOWITZ	83	PL 126B 225	F. Bergsma <i>et al.</i>	(CHARM Collab.)
GOLDBERG	83	PRL 50 1419	M.S. Chanowitz, S. Sharpe	(UCB, LBL)
HOFFMAN	83	PR D28 660	H. Goldberg	(NEAS)
KRAUSS	83	NP B227 556	C.M. Hoffman <i>et al.</i>	(LANL, ARZS)
VYSOTSKII	83	SJNP 37 948	L.M. Krauss	(HARV)
KANE	82	PL 112B 227	M.I. Vysotsky	(ITEP)
CABIBBO	81	PL 105B 155	Translated from YAF 37 1597.	
FARRAR	78	PL 76B 575	G.L. Kane, J.P. Leveille	(MICH)
Also		PL 79B 442	N. Cabibbo, G.R. Farrar, L. Maiani	(ROMA, RUTG)
			G.R. Farrar, P. Fayet	(CIT)
			G.R. Farrar, P. Fayet	(CIT)